

Thermodynamically-Favoured Bulk O₂ Formation on Charging in LiCoO₂ and LiNiO₂ Cathode Materials

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The formation of O₂ molecules within the bulk of LiCoO₂ and LiNiO₂ cathode materials during charging has been reported experimentally, but the validity and significance of these observations remain debated. Using density functional theory calculations, we demonstrate that at high states of charge, layered LiNiO₂ and LiCoO₂ are thermodynamically unstable with respect to both O₂ gas release and internal decomposition forming bulk-confined O₂ molecules, indicating that, from a thermodynamic perspective, the oxygen chemistry of these stoichiometric layered oxides parallels that of Li-rich oxides, and may likewise limit their long-term cycling stability.

The useful energy density of transition metal (TM) oxide lithium-ion cathodes over repeated charge-discharge cycles is limited by their oxygen chemistry. At high voltage, TM-oxide cathodes release gas-phase O₂, with this process linked to structural degradation, capacity loss, and gas build-up in cells [1]. It has recently been observed that in some TM-oxide cathodes, O₂ molecules also form within the bulk [2]. This bulk O₂ formation is linked to the formation of voids [3–7], and results in energy-density loss that cannot be prevented through surface engineering only.

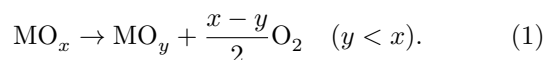
Bulk O₂ formation was first identified in Li-rich (Li_{1+x}M_{1-x}O₂) cathodes (where M = Mn, Co, Ni), from resonant inelastic X-ray scattering (RIXS) data, magnetic probes, and neutron diffraction [2, 8]. The standard explanation for O₂ formation in Li-rich cathodes is that the full charge capacity of these materials cannot be accounted for by formal TM redox alone [2], assuming a maximum accessible TM oxidation state of 4+. The additional capacity then is attributed to partially reversible bulk O₂ formation [2], associated with oxidation of O²⁻ ions with collinear Li–O–Li coordination [9].

Recently, bulk O₂ formation has been reported in Ni-rich Li-stoichiometric (LiMO₂-based) cathode materials, suggesting possible similarities with the oxygen chemistry of Li-rich cathodes and raising questions about their internal degradation processes. However, the interpretation of the characteristic RIXS signature of molecular O₂ in charged LiNiO₂ [4, 10] as evidence for bulk O₂ has been contested [11].

The possibility of bulk O₂ formation in LiMO₂ cathodes could be considered surprising when considered in the context of prevailing explanations for O₂ formation in Li-rich cathodes. In LiMO₂ cathodes, TM oxidation to a formal 4+ oxidation state can fully account for the theoretical capacity at the top of charge, seemingly negating the need for O oxidation to charge-balance Li⁺ extraction. Moreover, in non-defective layered LiMO₂ cathodes, the collinear Li–O–Li coordination proposed as nec-

essary for O²⁻ oxidation [9] is absent.

An alternative view of possible O₂ formation is provided by considering the corresponding chemical reaction from a thermodynamic perspective. For a fully delithiated cathode material MO_x, the thermodynamic driving force for forming O₂ is governed by the reaction



O₂ release from charged LiNiO₂ and LiCoO₂ can be rationalised if this reaction has a negative free energy change when the O₂ product is gas-phase [12]. Consequently, delithiated CoO₂ and NiO₂ are metastable and, if possible, will decompose into molecular O₂ and a more metal-rich oxide, MO_y, with the O₂ released as gas. Their ability to be cycled without decomposing is due to kinetic stability, specifically a combination of self-limiting oxygen release from the surface and kinetically slow TM cation diffusion in the bulk [12].

From a thermodynamic perspective, O₂ release and bulk O₂ formation differ only in the location and degree of confinement of the resulting O₂ [13]. Bulk O₂ forms at internal surfaces rather than at the external surface and is confined to a small volume within the cathode, resulting in a higher molar free energy. In Li-rich cathodes (e.g., Li₂MnO₃-based), reaction (1) has a negative free energy change for both O₂ release [14] and bulk O₂ formation [7], consistent with experimental observations of both phenomena in these systems. This established behaviour in Li-rich systems provides a valuable point of comparison for thinking about the oxygen redox chemistry in stoichiometric LiMO₂ cathodes.

To assess the thermodynamic drive for LiMO₂ (M = Co and Ni) cathodes at high states of charge to form trapped O₂, we computed formation energies of structures along the O₂–MO₂–MO tie-line for each system using density functional theory. These calculations employed the r²SCAN+rVV10 functional due to its accuracy describing O–O bonding energies (Methods). For the O₂ end-member, we considered both gas-phase O₂

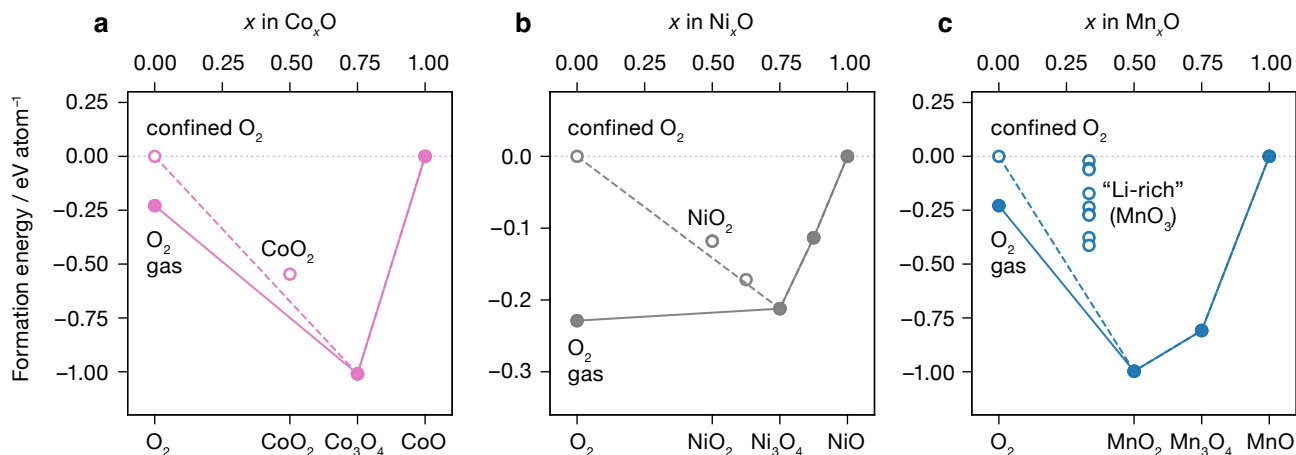


FIG. 1. Formation energy versus composition plots showing the thermodynamic stability of delithiated LiMO_2 cathode materials. The x -axis represents the composition of the structures along the O_2 – MO_2 – MO tie-line, while the y -axis represents their formation energy. Tie-lines are constructed for $M =$ (a) Co, (b) Ni, (c) Mn with the O_2 end-member calculated either as gas-phase O_2 (solid lines) or O_2 confined within the bulk (dashed lines). Filled markers denote ground-state structures, while empty markers represent metastable structures.

and O_2 confined in the cathode bulk. For comparison, we include equivalent data for the O_2 – MnO_2 – MnO system. Our calculated formation energies are presented in Fig. 1 as convex hull plots: points above each convex hull correspond to phases that are unstable with respect to decomposition into competing phases with the same average composition; this condition corresponds to a negative free energy change for reaction (1), while points on the convex hull are thermodynamically stable.

The free energy of bulk-confined O_2 is higher than that of gas-phase O_2 by ~ 0.23 eV/atom due to its high density (~ 1.7 g cm^{-3}) [7]. Despite this, our results predict that layered CoO_2 and NiO_2 are both unstable with respect to decomposition into their respective metal-rich MO_x phases ($x < 2$) and molecular O_2 , even if the O_2 is then confined in the bulk (Fig. 1(a,b)). In contrast, MnO_2 is thermodynamically stable (Fig. 1(c)). As a point of comparison, we also show points for delithiated Li-rich Mn-oxides, MnO_x ($x > 2$), which are predicted to be unstable with respect to forming gas-phase or confined O_2 , in agreement with previous studies [7, 14].

Our analysis indicates that formation of bulk-trapped O_2 is thermodynamically favourable in Li-stoichiometric layered LiCoO_2 and LiNiO_2 on charging, just as in Li-rich $\text{Li}_{1+x}\text{M}_{1-x}\text{O}_2$ cathode materials. This result explains experimental reports of the RIXS signature of molecular O_2 in LiNiO_2 -based stoichiometric cathodes, [4, 10, 11] and observed void formation [4, 6, 10]. These phenomena are linked: voids correspond to regions containing only O_2 and are formed by TM migration [7], driven by phase separation of delithiated MO_2 into O_2 and MO_x ($x < 2$).

Recent DFT calculations have shown that in charged LiNiO_2 , O_2 formation is energetically favourable at twin

boundaries [15]. Our results extend this understanding, indicating that the formation of bulk O_2 is thermodynamically favoured even in the perfect crystal. However, the formation of O_2 within the bulk of stoichiometric cathode materials depends not only on the thermodynamics but also on the kinetics of this process. O – O dimerisation requires localisation of holes on specific O^{2-} ions, and this is likely facilitated by point and extended defects that break crystalline symmetry. Any consequent O^{2-} formation will further disrupt the local crystal structure, enabling subsequent hole trapping and associated O^{2-} formation and, potentially, void growth.

Our results suggest that the long-term cycling stability in LiMO_2 -based cathode materials depends on kinetic stabilisation not only of their surfaces but also within their bulk. As such, one potential strategy to enhance stability is to minimise defects that may act as nucleation sites for bulk O_2 formation.

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