

Estimation of ligand binding free energy using multi-eGO.

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ABSTRACT: The computational study of the binding of a ligand to a target protein provides mechanistic insight into the molecular determinants of this process and can improve the success rate of *in silico* drug design. All-atom molecular dynamics (MD) simulations can be used to evaluate the binding free energy, typically by thermodynamic integration, and to probe binding mechanisms, including the description of protein conformational dynamics. The advantages of MD come at a high computational cost, which limits its use. Such cost could be reduced by using coarse-grained models, but their use is generally associated with an undesirable loss of resolution and accuracy. To address the trade-off between speed and accuracy of MD simulations, we describe the use of the recently introduced multi-eGO atomic model for the estimation of binding free energies. We illustrate this approach in the case of the binding of benzene to lysozyme by both thermodynamic integration and metadynamics, showing multiple binding/unbinding pathways of benzene. We then show equally accurate results for the binding free energy of dasatinib and PP1 to Src kinase by thermodynamic integration. Finally, we show how we can describe the binding of the small molecule 10074-G5 to Aβ42 by single molecule simulations and by explicit titration of the ligand as a function of concentration. These results demonstrate that multi-eGO has the potential to significantly reduce the cost of accurate binding free energy calculations and can be used to develop and benchmark *in silico* ligand binding techniques.

INTRODUCTION

The study of ligand binding plays a pivotal role in the understanding of biochemical pathways and in drug design. Ligands, which can be small molecules, peptides, or other macromolecules, interact with specific targets modulating their conformations, interactions and biological activity¹⁻³.

Characterizing a ligand-receptor interaction is a three-level problem (structure, thermodynamics and kinetics) that involves the characterization of the bound configuration (the pose), of the stability of the interaction (the binding free energy), and of the kinetics of the process (the on and off rate constants) including the detailed binding mechanism. Although a wide range of experimental and computational methods have been developed to address all these aspects⁴⁻⁶, there is still a need for scaling them up to reduce time and cost of the studies, in particular for drug design.

A common computational approach for drug design is based on the use of molecular docking methods⁷⁻⁹, including recent machine learning ones^{10,11}, to search and score the binding pose for large chemical libraries. The binding process of the resulting top-scoring ligands can then be studied

by a variety of methods, which include free energy calculations such as thermodynamic integration (TI) and free energy perturbation (FEP), long time-scale MD simulations, Markov state models, and reaction coordinate based techniques such as umbrella sampling and metadynamics (MetaD)¹²⁻¹⁹. The intensive computational cost of these methods can be reduced by lowering the resolution of the structural representation of the ligand-receptor system (i.e. by coarse graining), but usually at the price of reducing the accuracy of the results^{20,21}. There is therefore a need of developing methods that achieve a more advantageous trade-off between speed and accuracy.

Recently, we introduced multi-eGO to enable the simulation of protein self-assembly processes²². Multi-eGO simplifies the description of protein interactions by deriving effective Lennard-Jones (LJ) parameters by combining one or more training conventional MD simulations of a system describing its fluctuations in one or more free energy minima of interest with prior models representing baseline interactions. We showed that multi-eGO enables to obtain accurate results on protein folding, intrinsically disordered proteins (IDPs) conformational dynamics, and peptide aggregation,

at a fraction of the computational cost while maintaining atomic resolution²³.

Here we extended the applicability of the multi-eGO approach to ligand binding simulations (Figure 1). The training data consist in the simulations of a protein in its apo and holo states. The holo training is performed with additional free ligands in the box to also estimate the probabilities of non-specific interactions (both protein-ligand and ligand-ligand). The apo training intramolecular contact probabilities are reweighted by the previously introduced random coil prior, which represents a protein as a self-avoiding polymer. To reweight the protein-ligand and ligand-ligand contact probabilities, we introduce as prior a simulation performed at the same concentration of ligands of the training but with ligands and protein interacting only by hard-core repulsions (i.e. LJ C_{12} interactions). This prior captures the rotational and translation entropy associated to the free ligand at the training concentration. We expect to have a two-fold advantage: (i) the cost of training simulations, while not negligible, to be much less than that of a TI and MetaD calculation with a conventional force field, and (ii) for TI (and by extension also other endpoint techniques) removing both the explicit solvent and the electrostatic coupling allows us to have a single calculation to perform without the need of multiple thermodynamic cycles. For MetaD the smaller number of particles makes the multi-eGO-based simulations have a negligible cost in comparison to that of the training and to the same calculations at all-atom resolution (cf. Table S1 and S2).

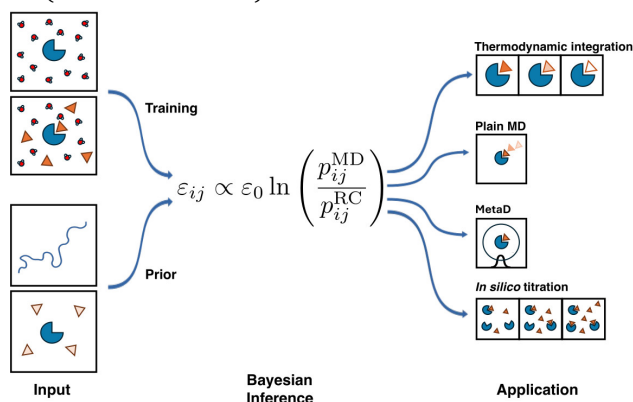


Figure 1. Schematic representation of the multi-eGO approach for ligand bindings. Training simulations for the apo and holo state, including free ligands in the box are weighted, following Bayes, for prior models representing a self-avoiding polymer for the protein and the random rotations, translations and clashes of the ligands with the folded proteins. The resulting model can be employed to efficiently perform multiple simulation techniques including TI, MetaD, plain MD as well as concentration dependent simulations.

To test the applicability of the multi-eGO approach, we applied it to four systems of increasing complexity. The first test case is the binding of benzene (BNZ) to the L99A mutant of the bacteriophage T4 lysozyme (LYZ)^{24,25}, which is

widely used to benchmark computational methods of protein-ligand binding^{26–29}. LYZ is a small enzyme (162 amino acids in the case of T4) that catalyzes the hydrolysis of 1,4- β -linkages of cell-wall peptido-glycans and consists of two domains (N-terminal and C-terminal). The L99A mutation produces a small apolar cavity in the C-terminal domain, which allows the binding of small nonpolar ligands such as BNZ with a binding free energy $\Delta G = 21.7 \pm 0.8$ kJ/mol²⁴. The second and third systems are the binding of dasatinib (DAS) and PP1 to the c-Src kinase (SRC). Dasatinib is a chronic myeloid leukemia drug that targets at nM affinity multiple tyrosine kinases, including Src; PP1 is an Src kinase inhibitor with lower IC_{50} ^{30–34}. Here, the two molecules allow us to explore the effect of complex (in terms of internal degrees of freedom) ligands and strong binding free energies, while also being well characterized *in silico*³⁵. The fourth system is the binding of the small molecule 10074-G5³⁶ to monomeric A β 42, which is an intrinsically disordered protein (IDP) whose aggregation is associated with Alzheimer's diseases, allowing to test a binding process dominated by conformational dynamics^{37,38}.

Our results show that multi-eGO can quantitatively estimate the binding free energies of benzene, dasatinib and PP1, reproduce their hypothesized binding mechanisms and, importantly, allow quantitative and systematic comparison of different simulation techniques. In the case of A β 42:10074-G5, we demonstrate that multi-eGO can account for extremely flexible systems, and we can provide a comparison between single-molecule simulations and concentration-dependent simulations, i.e. models able to closely mimic an *in vitro* experimental setup.

On the basis of these results, we suggest that multi-eGO represents an *in silico* platform for ligand binding studies that may be useful for both method development and benchmarking, as well as for integration into drug screening pipelines.

RESULTS

A multi-eGO model for lysozyme:benzene simulations. To train a multi-eGO model for LYZ:BNZ binding, we performed multiple apo and holo simulations using the DES-Amber³⁹ and the GAFF2 force field (see Methods and Table S1). With respect to the apo state, the simulations reproduced a breathing motion of the two domains over the microsecond time scale that is well characterized experimentally by residual dipolar couplings in nuclear magnetic resonance (NMR) spectroscopy⁴⁰ (Fig. 2A), as shown by the distribution of the root mean square deviation of the atomic positions (RMSD) calculated with respect to the backbone of a reference structure (PDB 1L84). The simulations of the holo state were performed with one BNZ molecule in the binding cavity and four additional molecules in solution. We found that BNZ is more likely to be found in regions surrounding the binding cavity, suggesting a role of surface interactions in directing BNZ to its binding site (Fig. 2B).

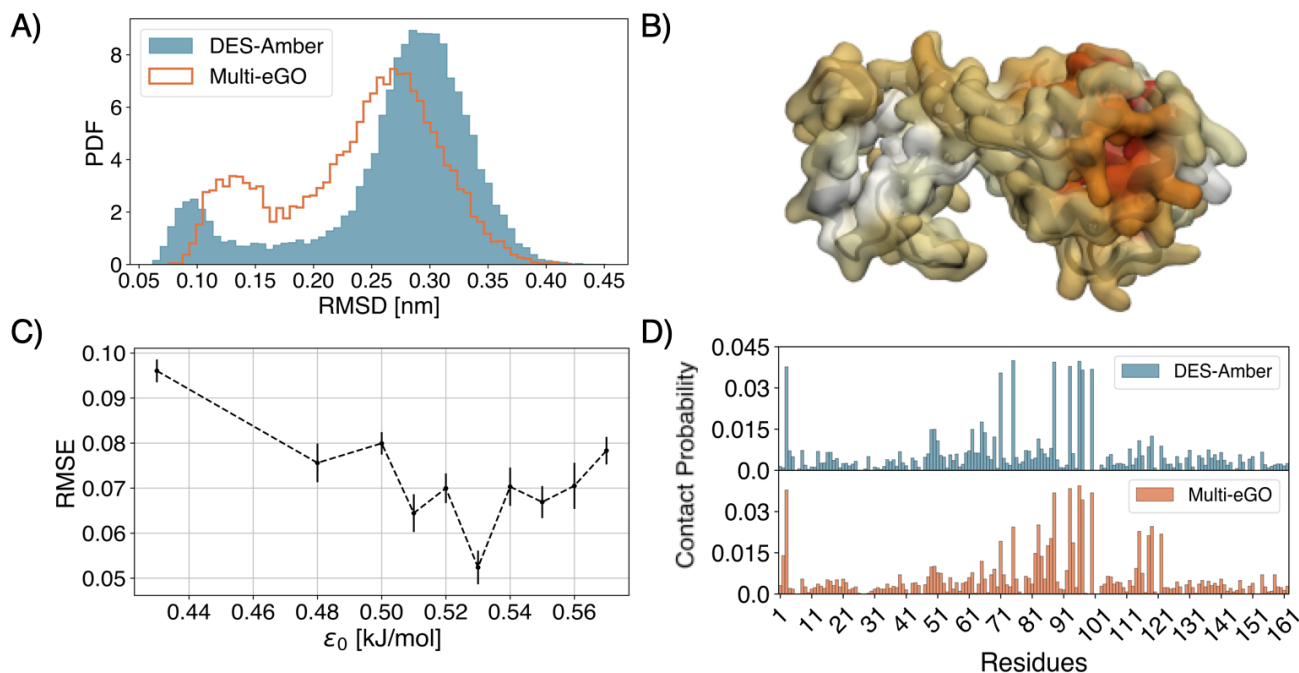


Figure 2. (A) RMSD distributions comparing to a reference structure (PDB 1L84) the ensemble of structures from the LYZ DES-Amber and multi-eGO simulations. These RMSD distributions illustrate that multi-eGO effectively reproduces the two-state distribution observed in LYZ. (B) Representation of the LYZ:BNZ interaction obtained from training simulations. Increased red coloring indicates stronger attraction, highlighting the binding pocket in the C-terminal domain and various surface interactions of the protein, some of which correspond to binding entry pathways. (C) Scoring function used to calibrate the free parameter of the model for intermolecular interactions. The RMSE was computed by comparing LYZ:BNZ contact probabilities between training and multi-eGO simulations. Error bars were derived by averaging the RMSE over different segments of multi-eGO trajectories. (D) Residue-wise intermolecular contact probabilities from training simulations versus multi-eGO simulations, optimized with an $\epsilon_0^{\text{inter-molecular}}$ of 0.53 kJ/mol.

The apo training simulation described above was analyzed to obtain intramolecular interatomic contact probabilities, which were reweighted using two priors. First, a random coil (RC) simulation was run, representing the contact probability of a self-avoiding polymer with the same sequence of LYZ. Then, an interdomain reference was used to decouple the intra-domain and inter-domain multi-eGO energy scales and correctly reweight the inter-domain training probabilities. The inter-domain prior consisted of a multi-eGO simulation in which all the inter-domain interactions between the N-terminal domain (residues 1-71) and the C-terminal domain (residues 72-162) were described as hard-core repulsions. This procedure allowed to obtain an informative prior of the random interdomain collisions (see Methods). The resulting multi-eGO model for LYZ, with an $\epsilon_0^{\text{intradomain}}$ of 0.34 kJ/mol and an $\epsilon_0^{\text{interdomain}}$ of 3 kJ/mol, reproduced the equilibrium between a closed and an open conformation, (see Fig. 2A and Fig. S1). Focusing on the LYZ:BNZ interactions, the intermolecular LYZ:BNZ and BNZ:BNZ contact probabilities were reweighted using an intermolecular reference simulation consisting of intermolecular (hard-core) clashes between LYZ and BNZ molecules with the same BNZ concentration used in the training simulation to account for the roto-translational entropy of ligands. The $\epsilon_0^{\text{intermolecular}}$ was obtained by minimizing the root mean square error (RMSE) between inter-molecular residue contact maps as shown in Fig 2C, resulting in an $\epsilon_0^{\text{intermolecular}}$ of 0.53 kJ/mol. The resulting multi-eGO model reproduced the intermolecular contact probabilities

between LYZ and BNZ (Fig. 2D), demonstrating that multi-eGO can effectively learn heterogeneous interactions.

As a first test, we then performed 40 unbiased binding simulations, starting from random configurations of 4 unbound BNZ molecules at the same concentration as the training simulation (see SI Movie 1 for a representative trajectory). The simulations were stopped when the minimum distance between the BNZ atoms and the C α atom of the A99 residue in the binding pocket fell below 0.4 nm. The extracted binding times were then fitted to a Poisson distribution, yielding a mean binding time τ of 963 ± 20 ns with a p-value of 0.988 obtained from a Kolmogorov-Smirnov (KS) test⁴¹, corresponding to a k_{on} of $7.6 \times 10^7 \text{ M}^{-1} \text{ s}^{-1}$ (see Fig. S2). This rate should be considered as nominal due to the simplified nature of the model, but when compared with the experimental value⁴² of $8 \times 10^5 - 10^6 \text{ M}^{-1} \text{ s}^{-1}$ it gives an indication of the intrinsic speed up of multi-eGO due to degrees of freedom removal.

Thermodynamic integration and volume-based metadynamics on lysozyme:benzene. Given the efficiency of multi-eGO, we wanted to test its accuracy in estimating the LYZ:BNZ binding free energy by also comparing the consistency of alternative methods. We focused on Thermodynamic Integration (TI) and Volume Based Metadynamics (VMetaD), the former being the industry standard for such calculations⁴³ and the latter a potentially more informative but also more computationally expensive alternative²⁸.

TI of LYZ:BNZ was performed in 6 replicates using three alternative restraints (CL1, CL2, and CL3) to keep the ligand

in the binding pose and to speed up the convergence²⁹ (see Methods and Fig. S3). The binding free energies, calculated as the mean and the standard deviation of the mean, are shown in Fig. 3A. The three estimates of 21.41 ± 0.04 kJ/mol, 21.45 ± 0.05 kJ/mol, and 21.46 ± 0.07 kJ/mol from the three restraints, respectively, show a high statistical precision and a remarkable accuracy when compared to the experimental binding free energy of 21.7 ± 0.8 kJ/mol (Fig. 3A).

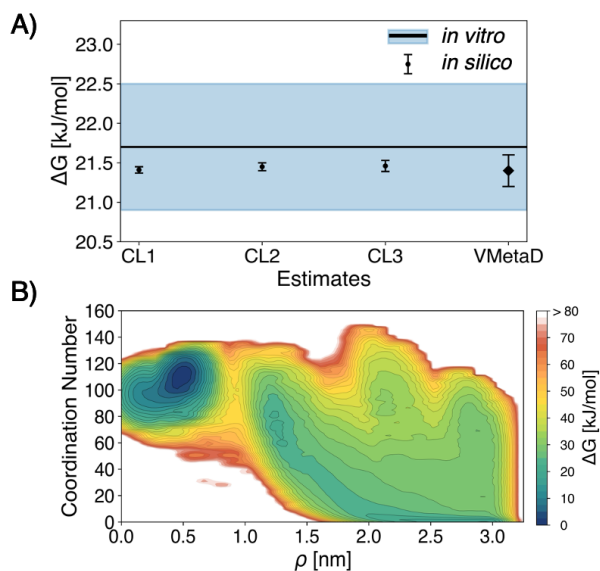


Figure 3. (A) Estimation of the binding ΔG obtained with the multi-eGO force field using different techniques: CL1, CL2, and CL3 refer to TI with three different restraints. The error was calculated as the standard error of the mean of six replicates for each restraint. VMetaD denotes the volume-based metadynamics ΔG estimate, with the error calculated as the standard error of the mean of four replicates. (B) Free energy surface (FES) obtained by reweighting of VMetaD in an appropriate collective variable (CV) space. On the x-axis, ρ represents the distance of the BNZ center of mass from the center of the sphere used to define the sampling volume. On the y-axis, the coordination number of BNZ with the residues of LYZ.

An additional binding free energy estimate was obtained using VMetaD (see Methods) by running four independent replicates of 1.5 μ s each, reaching convergence after approximately 200 ns (Fig. S4). A free energy surface (FES) representing the binding process was obtained by reweighting VMetaD simulations as a function of the distance of the BNZ center of mass from the center of the sphere centered on the binding site, which was used to define the sample volume, and the coordination number of BNZ with the residues of LYZ (see Fig. 3B and Methods section). From the FES, we estimated a ΔG of 21.4 ± 0.2 kJ/mol, obtained as the mean and standard deviation of the mean over the 4 replicates (Fig. 3A), which is within the statistical precision of the TI estimates, demonstrating the reliability of the two different approaches. The use of VMetaD also allowed us to

compare our sampled binding/unbinding pathways with those already reported in the literature. Using a dynamic time-warping clustering algorithm⁴⁴, we were able to observe most, if not all (depending on the classification approach), of the previously observed pathways^{28,44} (see Fig. S5 and S6), further supporting the ability of multi-eGO to provide relevant details about the binding mechanism.

Thermodynamic integration of c-Src kinase:dasatinib and c-Src kinase:PP1 binding. To test the transferability of the protocol developed for LYZ:BNZ, we reapplied it to the binding of DAS and PP1 to SRC. First, we performed several training MD simulations for the apo system as well as for the SRC:DAS and SRC:PP1 holo systems, in both cases also adding 4 ligands in solution (see Methods and Table S1). The training simulations showed different interactions between ligands and the protein in correspondence of known binding sites, as shown in previous work. Besides the ATP binding site, we found important interactions with the N-lobe (β -sheet), the α G helix, the P-loop, binding in the MYR site for both ligands and binding in the PIF site for PP1 (see Fig. S7).

To parameterize the multi-eGO models, after performing an RC simulation of SRC, we set the $\epsilon_0^{\text{intramolecular}}$ parameter for the protein to match the radius of gyration (R_g), the residue-wise root-mean square fluctuations of the atomic positions (RMSF), and RMSD distributions, finding an optimal value of 0.33 kJ/mol (Fig. S8). To describe intermolecular interactions, we first performed an intermolecular reference simulation at the same ligand concentration as the training simulation to account for random collisions between ligand molecules. We then optimized $\epsilon_0^{\text{ligand-ligand}}$ and found a value of 0.43 kJ/mol and 0.6 kJ/mol for DAS:DAS and PP1:PP1, respectively, which could reproduce their training contact probabilities. A second intermolecular reference simulation was then performed in which the ligands were allowed to interact with each other, while the protein and ligands only interacted with hard-core repulsion. Finally, by minimizing the RMSE between the training and multi-eGO residue-wise intermolecular contact probabilities, we found an optimal $\epsilon_0^{\text{intermolecular}}$ of 0.32 kJ/mol and 0.35 kJ/mol for SRC:DAS and SRC:PP1, respectively (Fig. S9).

The resulting multi-eGO models were used to perform 40 binding simulations for both PP1 and DAS, starting from a random conformation of four unbound ligands (at the same concentration as in training). To establish the binding, we chose three reference atoms of the protein, namely CD:GLU310, N:MET341, CA:THR338 for SRC:DAS and N:ILE294, N:MET341, N:THR338 for SRC:PP1, and we stopped each simulation when the three minimum distances between one of the ligands and the three reference atoms were below 0.5 nm, 0.4 nm, 0.5 nm for SRC:DAS and 0.5 nm, 0.35 nm, 0.5 nm for SRC:PP1.

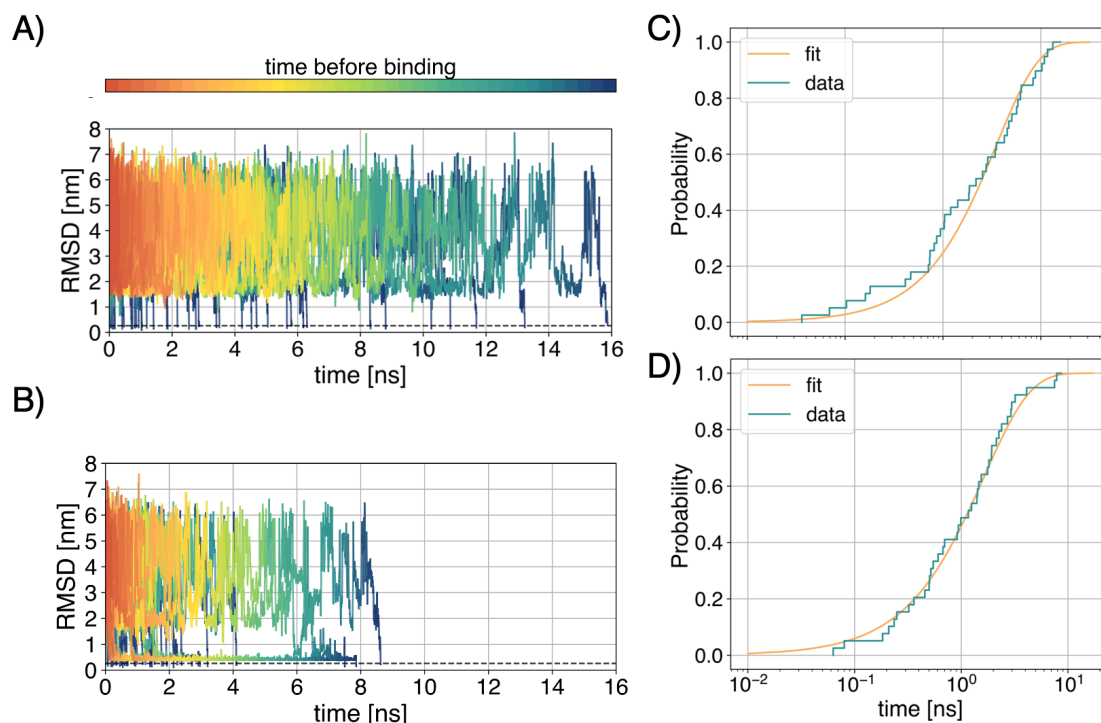


Figure 4. (A) RMSD of DAS with respect to its experimental binding pose over time. The ligand is considered bound when the RMSD is less than 0.26 nm. (B) Cumulative DAS binding times and Poisson fit with mean binding times of 3.5 ± 0.1 ns, $p=0.92$ from a Kolmogorov-Smirnov test. (C) RMSD of PP1 with respect to its experimental binding pose over time. The ligand is considered as bound when the RMSD is less than 0.26 nm. (D) Cumulative PP1 binding times and Poisson fit with mean binding times of 1.64 ± 0.02 ns, $p=0.99992$ from a Kolmogorov-Smirnov test.

The calculated RMSD of the ligand to its bound conformation after alignment of the protein is shown in Fig. 4A,C (see also SI Movie S2 and S3). Binding times (RMSD less than 0.26 nm) were fit to a Poisson distribution, yielding mean binding times of 3.5 ± 0.1 ns and 1.64 ± 0.02 ns for SRC:DAS and SRC:PP1, respectively, with corresponding p -values of 0.92 and 0.99992 obtained with KS-test (Fig. 4B,D). The corresponding k_{on} for SRC:DAS and SRC:PP1 are $2.51 \times 10^{10} \text{ M}^{-1}\text{s}^{-1}$ and $5.39 \times 10^{10} \text{ M}^{-1}\text{s}^{-1}$, respectively. Despite the fact that, as for LYZ:BNZ, the time scales are nominal (i.e., the *in vitro* k_{on} for DAS is $\sim 5 \times 10^6 \text{ M}^{-1}\text{s}^{-1}$), the relative rates we found are consistent with those obtained by previously published MD simulations³⁵, with both ligands reported to have k_{on} of the same order of magnitude, with a slightly higher k_{on} for PP1.

Having shown that the model can simulate the correct binding of the two ligands, we estimated their binding free energy by TI, using three alternative restraints as before (see Methods and Fig. S10 and S11). For each restraint we performed 4 TI replicates. For SRC:DAS we estimated a binding free energy of 55.0 ± 0.5 kJ/mol, 55.5 ± 0.3 kJ/mol, 55.6 ± 0.4 kJ/mol for the three restraints, respectively (Figure 5A), while for SRC:PP1 we estimated a binding free energy of 38.34 ± 0.06 kJ/mol, 38.4 ± 0.1 kJ/mol, 38.3 ± 0.1 kJ/mol (Figure 5B). As for LYZ:BNZ all estimates are compatible and show a high statistical precision.

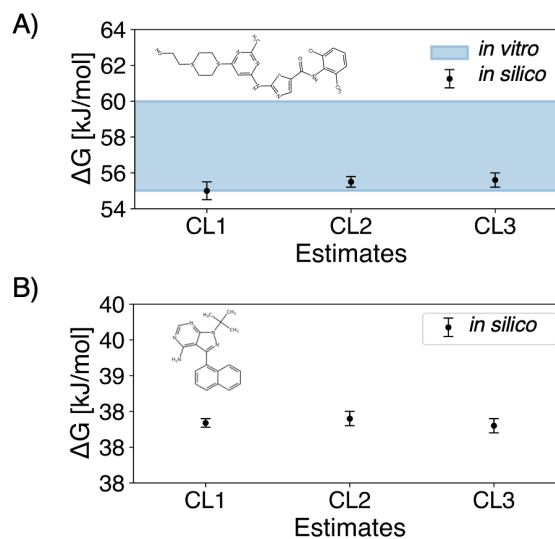


Figure 5. (A) TI estimates of the binding ΔG of c-Src kinase and dasatinib using multi-eGO. The three estimates (CL1, CL2, CL3) correspond to different restraints used to keep the ligand in the binding pose during decoupling. Errors were calculated as the standard error of the mean of four replicates for each restraint. The blue band represents the interval of *in vitro* measurements (REF). (B) TI estimates of the binding ΔG of c-Src kinase and PP1 using multi-eGO. CL1, CL2, CL3 corresponds to different restraints. Errors were calculated as the standard error of the mean of four replicas for each constraint.

Importantly, the ΔG obtained for SRC:DAS is compatible with values reported in the literature for *in vitro* measurements, ranging from 55.17 kJ/mol to 60.23 kJ/mol (Figure 5A). For SRC:PP1, we compared the experimental ratio of the IC_{50} of SRC:DAS (0.5 nM)⁴⁵ and SRC:PP1 (170 nM)³⁴ with the ratio of the $K_D = \exp\left(-\frac{\Delta G}{k_B T}\right)$ estimates obtained with multi-eGO. We found that the ratio of the multi-eGO K_D values is in qualitative agreement with the IC_{50} ratio, namely $\frac{K_D^{PP1}}{K_D^{DAS}} = 740 \pm 114$ and $\frac{IC_{50}^{PP1}}{IC_{50}^{DAS}} = 340$.

Single-molecule and titration simulations of A β 42:10074-G5. The above binding processes represent the most common scenario where a pose is known or hypothesized under the lock-and-key mechanism where the ligand minimally perturbs the structure of the receptor. At the other end of the spectrum is the case of ligand binding to IDPs, which do not entail a well-defined binding site due to the absence of a stable tertiary structure. To test multi-eGO in this scenario, we used the large publicly available dataset for A β 42 simulated in its apo and 10074-G5 holo states^{37,46}. After parametrizing the A β 42 monomer as previously published with a $\epsilon_0^{intramolecular}$ of 0.33 kJ/mol, we parametrized 10074-G5 starting from its GROMOS ATB-derived model⁴⁷ and optimized its bonded parameters to match those of the training simulation (see Methods). We then performed a reference simulation in which mutual interactions between A β 42 and 10074-G5 are represented only by steric clashes in the same training condition. After this simulation we set the intermolecular interaction between A β 42 and 10074-G5 by minimizing the RMSE of the intermolecular residue-wise contact maps (see Fig. S12), obtaining a value of $\epsilon_0^{intermolecular}$ of 0.385 kJ/mol. The multi-eGO holo-simulation showed the same R_g behavior observed in the training simulation (Fig. S13), qualitatively reproducing the same peak shift and narrowing of the distribution. Note that this effect is only due to the ligand binding, since the intramolecular interactions are only learned from the apo simulation.

With the parameters found above, we ran a 2 μ s multi-eGO MD simulation and calculated the probability of binding vs. unbinding to estimate the K_D . In the case of single molecule

simulations, direct estimation of K_D as $\frac{[L][P]}{[LP]}$ is affected by the finite size of the simulation box⁴⁸. To correctly calculate the dissociation constant, considering the box size, we followed Lopez et al. and calculated K_D as

$$K_D = \frac{1}{V N_A p_B(V)} \frac{1 - p_v(V)}{1 - \frac{v}{V}},$$

where V is the box volume, N_A is Avogadro's number, $p_B(V)$ is the bound probability, v is the interacting volume (the volume where protein and ligand interact but the ligand is not in the bound pose), and p_v is the probability of finding the ligand within the volume V . In the case of ligand binding with an IDP, the difference between the interacting and bound volumes is subtle (with no precise binding pose).

By defining the bound state considering the configurations where the minimum distance between A β 42 and 10074-G5 is less than 0.33 nm and as interacting the state where the minimum distance is less than 0.4 nm, we obtained a value of the dissociation constant of $683 \pm 7 \mu$ M for multi-eGO (with the error calculated as standard deviation of the mean of 4 replicates of 2 μ s each) (see Figure 6A), to be compared with $381 \pm 5 \mu$ M obtained by analyzing the training simulation (where the error is calculated using a bootstrapping method with 95th percentile). We note that the difference between the training and multi-eGO dissociation constants corresponds to 1.5 kJ/mol in free energy, less than the thermal fluctuation energy (~ 2.5 kJ/mol at 300 K). However, these values are not in quantitative agreement with the experimental value of 7-40 μ M, most likely because of limitations in the force-field accuracy. In this case, a substantial advantage of using multi-eGO is represented by the possibility to update the model to reproduce the experimental affinity by changing a single parameter, $\epsilon_0^{intermolecular}$, which represents the energy scale for the A β 42:10074-G5 interaction, without the need of repeating the computationally-intensive all-atom training simulation. A $\epsilon_0^{intermolecular}$ of 0.46 kJ/mol corresponded to a dissociation constant of $18 \pm 1 \mu$ M (Figure 6A), before and after the volume correction, where the data represent the mean and the standard deviation of the mean obtained from 5 replicates of 2 μ s each.

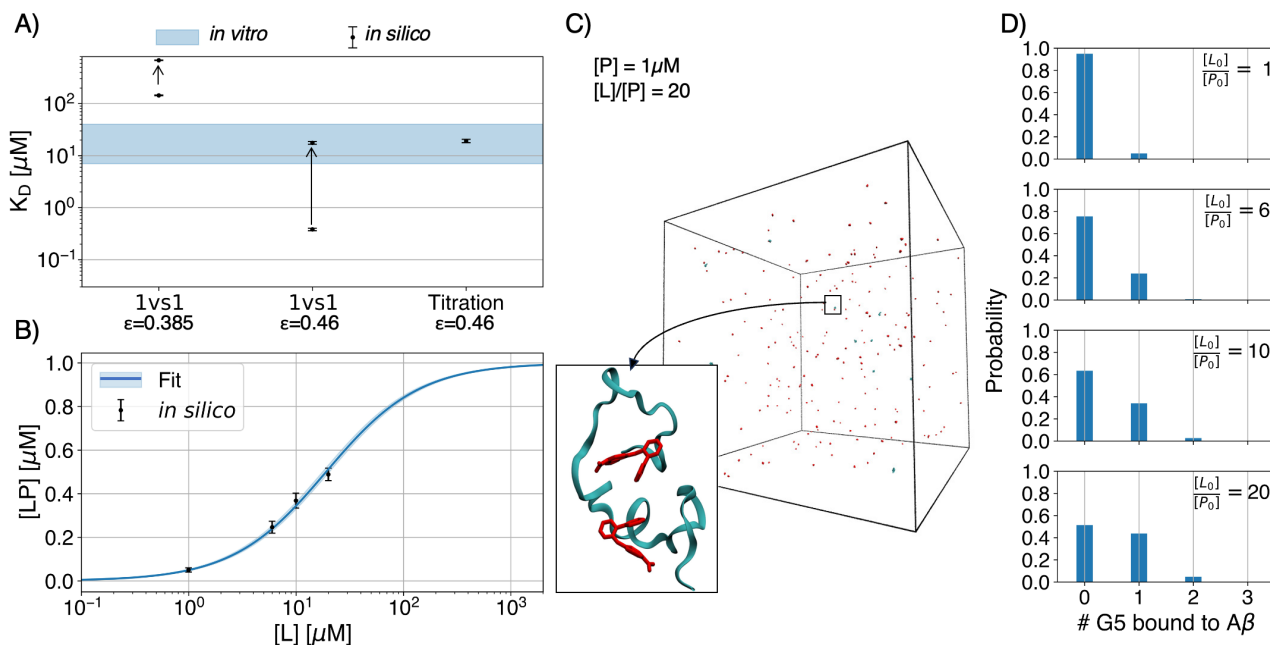


Figure 6. (A) Comparison of dissociation constants obtained from single molecule and titration multi-eGO simulations. The arrows indicate the correction for the finite size effects. (B) Aβ42:10074-G5 bound state population as a function of 10074-G5 concentration as obtained from multi-eGO 10074-G5 titration simulations, the line is the fit using the Hill equation. (C) Display of a simulation box for 10 Aβ42 molecules at 1 μM concentration mixed with 200 10074-G5 molecules. The inset shows an Aβ42 configuration bound to two 10074-G5 molecules. (D) Probability of observing 0 to 3 10074-G5 molecules simultaneously bound to an Aβ42 molecule as a function of 10074-G5 concentration.

The updated model was then used to study Aβ42:10074-G5 binding in the presence of an increasing number of 10074-G5 molecules. Following *in vitro* experiments, we performed an *in silico* titration experiment mixing 1 μM Aβ42 concentration (consisting of 10 monomers in a 255 nm cubic box, cf. Fig. 6C) with increasing concentrations of 10074-G5, namely 1:1, 1:6, 1:10, and 1:20 Aβ42:10074-G5 ratios. For each of these, we ran 5 μs long multi-eGO simulations (see Fig. S14 for binding probability as a function of time and block averaging), and by fitting the concentration of bound Aβ42:10074-G5 as a function of ligand concentration to the equation $[LP] = \frac{[L][P]}{[L]+K_D}$, we found a K_D of 19.2 ± 1.4 μM (as shown in Figure 6B with error estimates resulting from a 1 μs block averaging), which is consistent with that obtained with the volume-corrected single molecule approach. This case further emphasizes the usefulness of multi-eGO as a platform for method comparison and development, in this case providing a direct proof of the validity of the volume correction of Lopez *et al.*⁴⁸ when estimating dissociation constants from single molecule equilibrium simulations.

In Figure 6D, we report the probability of having multiple concurrent bindings at the different concentrations, showing the presence of 2 or 3 ligands bound to one Aβ42 monomer (see also the inset of Figure 6C for a representative configuration). This analysis suggests that the effect of simultaneous binding is negligible at the concentrations of the experiments. The lack of cooperativity is consistent with the entropic expansion model with minimal structural perturbations⁴⁹. Of note, we anticipate that our simulations could

be extended to account for additional Aβ42 interactions, such as those that could be derived from a fibril structure, paving the way for the study of protein aggregation itself and its inhibition mechanisms.

DISCUSSION

Structure-based models have been introduced to the study of protein folding based on the hypothesis that the native structure of a protein should capture the most relevant interactions across the whole free energy landscape of a protein. They have proven useful for studying the otherwise inaccessible folding mechanism at high resolution⁵⁰. Recently, we have introduced multi-eGO, an ensemble-based model rooted in Bayesian statistics, where one or more conformational ensembles of a protein, representing the fluctuations of relevant free energy minima, are weighted with prior information^{22,23}. This approach enables to obtain a model capable of simultaneously describing folded and disordered proteins, as well as self-assembly processes such as peptide amyloidogenic aggregation.

Here, we have extended the multi-eGO approach to include small molecules. Our initial results on four case studies allow us on the one hand to propose multi-eGO as a model for benchmarking and developing *in silico* techniques for ligand binding studies, and on the other hand to suggest that multi-eGO can be used for quantitative estimation of ligand affinities. This is because multi-eGO describes the system at atomic resolution, with binding free energies compatible with those measured *in vitro*, but with a dramatic acceleration of the kinetics (that we can tentatively quantify between 2 and 4 orders of magnitude). We exploited the LYZ:BNZ system to apply orthogonal free energy calculation

techniques (TI with multiple restraining conditions and VMetaD) showing their applicability and robustness on the model. Note that all simulations are replicated multiple times to obtain statistically meaningful estimates, with a minimal computational cost. While A β 42:10074-G5 allowed us to show a use case inaccessible to the current conventional approaches, namely explicit titration as a function of ligand concentration, with a direct comparison with concentrations used *in vitro*. This, on the one hand, demonstrated the efficiency of multi-eGO and, on the other hand, underlined the importance of correcting single-molecule simulations for the finite-size effect, following procedures such as those described in Lopez et al.

The use of multi-eGO requires the training of the model, which consists of running training simulations and finding a set of free parameters. Here, we have introduced a possible strategy to set our energy scale parameters excluding the intervention of any possible user bias. In practice, we minimized the RMSE with the protein fluctuation as well as the RMSE for the intermolecular residue-wise contact map, comparing the multi-eGO and the training simulations. This approach is first successfully tested on the LYZ:BNZ system and then successfully replicated for the SRC:DAS case. Successfully, because in both cases the knowledge of the experimental binding free energy was not used in the input but was accurately obtained by the simulations. In both cases, our training simulations include a ligand in its binding pose as well as several ligands free to sample the surface of the protein. The latter data are used to set the energy scale of the system. In the case of A β 42:10074-G5, we used a previously generated large dataset of simulations designed to represent the equilibrium between the apo and holo states, but here our approach of learning equilibrium data about the entire binding process inherited the limitation of the training simulation, namely an overestimated dissociation constant. This is where the advantage of multi-eGO becomes apparent, as it is sufficient to tune a single parameter to make the model match the experimental value. This allows us to speculate that multi-eGO may be particularly suitable for studying ligand binding processes such as those represented by the LYZ:BNZ, SRC:DAS, and SRC:PP1 cases, where the training data can be generated with very limited computational cost and should represent only the weak (i.e., faster to sample) non-specific interactions of a ligand with the protein surface, given knowledge or a hypothesis about the binding pose. We anticipate that further work will be needed to better understand the strengths and limitations of multi-eGO in ligand binding studies, including cases with more complex binding mechanisms such as induce-fit or conformational selection, as well as to streamline and possibly automatize the simulation setup protocol.

CONCLUSIONS

Ligand binding studies are among the most important applications of MD simulations since this computational technique can capture the conformational dynamics of both the ligand and the receptor. It is thus possible to obtain a relatively accurate description of their physicochemical interactions, and both kinetic and thermodynamic information about the binding process. These results can be obtained provided that sufficient computational resources are

available. While several methods have been developed to reduce the computational cost of obtaining accurate kinetic or thermodynamic data, MD simulations are still affected by the problem of the timescale for sampling the relevant regions of the conformational space to obtain reliable estimates of thermodynamic properties. Our results highlight the potential of multi-eGO both as a model to benchmark and develop *in silico* free energy calculation techniques and as an accurate and efficient framework for ligand binding studies, potentially extending the current capabilities, in terms of time and number of particles, of *in silico* molecular studies.

METHODS

Conventional molecular dynamics simulations. All-atom, explicit solvent, training simulations of LYZ:BNZ, SRC:DAS and SRC:PP1, were performed using the DES-Amber force field³⁹ in TIP4P-D water⁵¹, with ligands parameterized using GAFF2. To match the scaled electrostatic interactions of DES-Amber force field, we rescaled the AM1-BCC-derived charges for the ligands by a factor of 0.9.

All systems were prepared at pH 7.4 and adding Na⁺ and Cl⁻ ions to maintain physiological salinity (150mM) and to neutralize the total charge of the system. All systems were subjected to energy minimization using the steepest descent algorithm until the maximum force converges to a value < 1000 kJ/(mol nm), followed by a conjugate-gradient minimization until the maximum force converges to a value < 100 kJ/(mol nm) and a LBFGS algorithm until the maximum force converges to a value < 10 kJ/(mol nm). Subsequently, the minimized configuration was relaxed for 2 ns at a constant pressure of 1 bar and constant temperature, keeping the protein atoms restrained to the position of the minimum energy configuration. The simulations used the leapfrog algorithm with a time step of 2 fs and LINCS restraints⁵² for hydrogen atoms. Non-bonded interactions were cut off at 1 nm using PME for long-range electrostatics. Temperature and pressure were controlled by stochastic velocity rescaling⁵³ and cell rescaling⁵⁴ algorithms, respectively. Number of replicates and simulation times are reported in Table S1. All simulations were run using GROMACS 2022⁵⁵.

LYZ was modelled starting from the 1L84 PDB structure⁵⁶ and simulated at 300 K. The training of the LYZ:BNZ complex was performed inserting one BNZ into and 4 BNZ outside the binding pocket in order to explore the interaction with the surface of the protein.

SRC was modelled from the 1Y57 PDB structure⁵⁷ and simulated at 310 K. To train the holo complex of SRC:DAS/PP1 we inserted 1 molecule (either DAS or PP1) in the binding pose and 4 ligands in the solvent to explore the interaction with the surface of the protein. To train our multi-eGO model we removed from the holo training trajectories the non-equilibrium portions consisting in irreversible aggregation of three or more ligands with each other, which was also previously observed³⁵.

MD simulations of apo and holo A β 42:10074-G5 were available from previous studies^{37,46}. Briefly, the apo simulation was carried out at 278 K using the CHARMM22* forcefield⁵⁸ and the TIP3P water model⁵⁹. The holo simulation performed under the same condition parameterizing the ligand 10074-G5 using the Force Field Toolkit⁶⁰.

Multi-eGO Simulations. Multi-eGO is a multi-ensemble hybrid transferable/non-transferable force field²². The non-transferable (structure-based) part of the force field is obtained from state-of-the-art simulation(s) of the system of interest sampling the fluctuations of specific states (i.e. monomer, apo, fibril, etc.), and is then weighted with a prior (or reference) simulation obtained with a simplified description of the system meant to account for trivial results. From the training and the reference simulations we extract pair-wise distances and contact probabilities P_{ij}^{MD} and P_{ij}^{RC} respectively. Then through a Bayesian reweighting we obtain an estimation of the atom-pairs non-bonded interaction energy as

$$\varepsilon_{ij} = -\frac{\varepsilon_0}{\ln P_{threshold}^{RC}} \cdot \ln \frac{P_{ij}^{MD}}{\max(P_{ij}^{RC}, P_{threshold}^{RC})},$$

where i, j are atom indices, ε_0 is the energy scale to be set as a parameter, and $P_{threshold}^{RC}$ is a minimum probability used for regularization. A detailed description of the model can be found in ref.²³ and the associated code and parameters are available on GitHub.

For each apo system we performed a reference RC simulation consisting in a self-avoiding chain obtained with bonded interactions and C_{12} repulsion multi-eGO transferable potential. For multi-domain protein like LYZ, a second layer of reweighting was introduced. This consists in a multi-eGO simulation in which we turn off the inter-domain attractions, namely between residues 1-71 and 72-162, while keeping the intra-domain ones and the local interaction in the alpha-helix connecting the two domains to maintain it intact. We then use the RC probabilities obtained from the self-avoiding chain for the intra-molecular interactions and the ones obtained from the inter-domain reference. This approach allows us to correctly decouple the inter from the intra domain interactions, assign to each of them a different global energy scale through different ε_0 and reweighting the training atom-pair contact probabilities with the correct prior distribution. This inter-domain decoupling does not affect the binding, since all relevant interaction are only within the C-terminal domain. For intermolecular interaction the reference simulation corresponds to protein:ligand complex at same concentration of the training simulation (considering only the unbound ligands) with only hard-core repulsion between protein and ligands. This allows to account for the concentration-dependent rotational entropy of the ligands removing the effect of the box size in the contact probability estimation. The ligands were parametrized starting from the training topology, removing hydrogens and optimizing the bonded parameters to reproduce the local geometries of the training simulations.

All multi-eGO MD simulation were performed using stochastic dynamics integration with a timestep of 5 fs and a relaxation time of 25 ps. The cutoff for the LJ interactions was set to $2.5\sigma_{max}$, corresponding to 1.44 nm. A 10% larger radius was used for the neighbor lists, which were updated every 20 steps. Different values of ε_0 were tested to maximize the overall agreement between training and multi-eGO simulation. All simulations were run using GROMACS 2022⁵⁵.

Thermodynamic Integration. For the multi-eGO binding free energy estimation, we used TI and the Bennett's accepting ratio (BAR) algorithm^{61,62}. For each system a set three cross linked restraints, namely CL1, CL2, and CL3, to keep the ligand in the binding pose during the decoupling. The effect of this restraining potential can be removed analytically a posteriori²⁹. The thermodynamic cycle was computed through a set of λ values scaling the restraint and the non-bonded interactions (Table S3 to S5). For LYZ:BNZ for each λ we performed a steepest descent energy minimization, followed by a 500 ps relaxation and a 1 ns run at 300K in the NVT ensemble. For SRC:DAS and SRC:PP1 for each λ we performed a steepest descent energy minimization, followed by a 2 ns relaxation and a 1 ns run at 310K in the NVT ensemble. All simulations were run using GROMACS 2022⁵⁵.

Volume-based Metadynamics. VMetaD²⁸ simulations LYZ:BNZ were ran using well-tempered metadynamics⁶³ considering 3 CVs, namely the relative position in spherical coordinates (ρ, θ, φ) of the benzene center of mass with respect to the center of mass of lysozyme binding domain (residues 72-162). We set an initial height of the Gaussians of 1.2 kJ/mol, the widths at 1 Å, $\pi/16$ rad, and $\pi/8$ rad for ρ, θ , and φ , respectively and a bias factor of 20, depositing a Gaussian every 10 ps. We run 4 different 1.5 μ s-long replicas of the simulations, employing the same MD parameters used in the unbiased multi-eGO simulations.

To avoid an extensive sampling time of the unbound state, we set a repulsive spherical potential at a distance of ρ_{sph} (here set at 30 Å) from the center of mass of LYZ:

$$U_{sph}(\rho) = \begin{cases} 0 & \text{if } \rho \leq \rho_{sph} \\ \frac{1}{2}k(\rho - \rho_{sph})^2 & \text{if } \rho > \rho_{sph} \end{cases}$$

with k set to 10000 kJ/mol/Å². To consider the loss of entropy caused by the imposition of such restraint and calculate the binding free energy difference in standard conditions, we apply the following entropic correction

$$\Delta G^0 = \Delta G_{MetaD} - RT \log \left(\frac{V^0}{\frac{4}{3}\pi\rho_{sph}^3 - V_{prot}} \right),$$

where ΔG_{MetaD} is the binding free energy of VMetaD, R is the gas constant, T is the temperature of the system, V^0 is the standard volume (1660 Å³), and V_{prot} is the volume of the protein included in the sphere restraint. Protein volume inside the sphere was calculated using the double cubic lattice method⁶⁴ available in GROMACS. To compute the ΔG_{MetaD} we reweighted the computed free energy surface using the Tiwary-Parrinello algorithm⁶⁵, removing the initial 200 ns of trajectory on CVs that allow the precise definition of the bound and the unbound states. Following ref.²⁸, we choose the distance ρ and the coordination number c , between the set of the ligand atoms A and the set of the protein atoms B , as

$$c = \sum_{i \in A} \sum_{j \in B} \frac{1 - \left(\frac{r_{ij}}{r_0}\right)^6}{1 - \left(\frac{r_{ij}}{r_0}\right)^{12}}.$$

The error estimation has been performed using the standard deviation of the mean of the four replicas. All simulations were run using GROMACS 2022⁵⁵ and PLUMED2^{66,67}.

In Silico Titration Experiment. Four 5 μ s long A β 42:10074-G5 MD multi-eGO simulations were performed by placing 10 monomers of A β 42 in a 255 nm cubic box, corresponding to a concentration of 1 μ M A β 42, and adding 10, 60, 100 or 200 10074-G5 at random, respectively. Simulations were run using GROMACS 2022⁵⁵.

ASSOCIATED CONTENT

Supporting Information

The supporting information file contains comparison between the training simulations and multi-eGO ones for all the showed systems, extended convergence plots for all the TI and VMetaD calculations performed, the representation of binding and unbinding pathways for LYZ:BNZ, the ϵ_0 selection details for SRC:DAS, SRC:PP1, and A β 42:10074-G5; supporting tables with the length of training simulations, extensive details about performances of all-atom and multi-eGO potentials, and the parameters for all the TI performed on the four systems. Finally, we added 4 movies showing the binding events from multi-eGO simulations in the four complexes.

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Author Contributions

BS performed and analyzed all the simulations. ES, with the support of TL, performed proof of principle simulations. BS, ES, FBT, RC, and CC developed the model. RC and CC supervised the work with the support of SG and MV. BS, RC and CC wrote the manuscript with contributions of all authors.

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Notes

Simulations data are publicly available via Zenodo with records DOI: 10.5281/zenodo.12800706, 10.5281/zenodo.12800760, 10.5281/zenodo.12800768; the multi-eGO code and parameters are publicly available on GitHub at <https://github.com/multi-ego/multi-eGO>, use the beta3 tag for a snapshot of the repository associated to this paper.

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ABBREVIATIONS

MD, molecular dynamics. TI, Thermodynamic Integration; VMetaD, Volume-based Metadynamics. RMSD, root mean square deviation. RMSF, root mean square fluctuations, RMSE, root mean square error. Rg, radius of gyration. LYZ, lysozyme. BNZ, benzene. KS, Kolmogorov-Smirnov. SRC, c-Src kinase. DAS, dasatinib. G5, 10074-G5. FES, free energy surface. RC, random coil. IDP, intrinsically disordered protein.

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