# Entropy of Liquid Water as Predicted by the Two-Phase Thermodynamic Model and Data-Driven Many-Body Potentials

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#### Abstract

We investigate the entropy of liquid water at ambient conditions using the two-phase thermodynamic (2PT) model, applied to both common pairwise-additive water models and the MB-pol and MB-pol(2023) data-driven many-body potentials. Our simulations demonstrate that the 2PT model yields entropy values in semiquantitative agreement with experimental data when using MB-pol and MB-pol(2023). Additionally, our analyses indicate that the entropy values predicted by pairwise-additive water models may benefit from error compensation between the inherent approximations of the 2PT model and the known limitations of these water models in describing many-body interactions. Despite its approximate nature, the simplicity of the 2PT model makes it a valuable tool for estimating relative entropy changes of liquid water across various environments, especially when combined with water models that provide a consistently robust representation of the structural, thermodynamic, and dynamical properties of liquid water.

# Introduction

Entropy is a central thermodynamic property across various fields, having broad implications not only in chemistry and physics but also in biology, information theory, and social sciences.<sup>1-4</sup> The second law of thermodynamics states that the total entropy of the universe cannot decrease, thus determining the spontaneity of all transformations.<sup>5</sup> For example, entropy changes play a crucial role in determining reaction dynamics,<sup>6-8</sup> understanding protein folding,<sup>9</sup> exploring phase transitions,<sup>10</sup> and designing high-entropy materials.<sup>11</sup> Although Boltzmann's statistical interpretation of entropy laid the groundwork over a century ago,<sup>12</sup> accurately calculating entropy changes for complex molecular systems continues to represent a significant challenge.

The two-phase thermodynamic (2PT) model provides a straightforward approach for estimating the entropy of a liquid by decomposing the total density of states into solidlike and gas-like components.<sup>13–15</sup> Within the 2PT framework, the solid-like component is approximated by normal modes describing an ideal solid, while the gas-like component is approximated by the diffusive motions of a hard-sphere gas. Since the diffusion of a liquid can be fully attributed to the gas-like component, the decomposition of the total density of states can be uniquely determined. It follows that the entropy of a liquid can then be obtained by summing the entropy contributions from the solid-like and gas-like components, which are calculated independently.

Water is arguably the most important liquid on Earth, directly mediating chemical, physical, and biological processes.<sup>16</sup> It is thus not surprising that many molecular models<sup>17–20</sup> have been developed since the first Monte Carlo (MC)<sup>21</sup> and molecular dynamics (MD)<sup>22</sup> simulations of liquid water. The most common models describe the water molecules as rigid objects and represent the underlying interactions as a sum of physics-inspired, pairwise-additive contributions. These models attempt to capture actual many-body interactions through effective pairwise terms that are empirical parameterized to reproduce a subset of experimental properties (e.g., density, freezing point, enthalpy of vaporization, etc.).<sup>18</sup> Examples of empirical pairwise-additive models include SPC,<sup>23</sup> SPC/E,<sup>24</sup> TIP3P,<sup>25</sup> TIP4P,<sup>25</sup> TIP4P,<sup>25</sup> TIP4P,<sup>25</sup> TIP4P,<sup>26</sup> and SPC/Fw<sup>30</sup>) have also been developed. Although empirical pairwise-additive models have been widely used in computer simulations of water and various aqueous solutions they suffer from intrinsic shortcomings that limit their predictive power and transferability across different phases.<sup>20</sup>

Recent advances in correlated electronic structure methods and machine learning algorithms have enabled the development of water potentials rigorously derived from the manybody expansion of the energy (MBE) calculated at the coupled cluster level of theory, including single, double, and perturbative triple excitations, *i.e.*, CCSD(T). CCSD(T) is currently considered the "gold standard" of chemical accuracy for molecular interactions.<sup>31,32</sup> Among these data-driven many-body potentials, MB-pol<sup>33–35</sup> has been shown to accurately reproduce the properties of water from gas-phase clusters to liquid water to ice.<sup>36,37</sup> Notably, MB-pol is the first and, to date, only molecular model capable of accurately predicting the phase diagram of water.<sup>38</sup> In addition, MB-pol has also enabled accurate simulations of the hydration properties of halide and alkali metal ions from small clusters<sup>39–46</sup> to aqueous so-lutions,<sup>47–51</sup> as well as of water adsorbed in metal-organic frameworks at different relative humidity values.<sup>52–60</sup> More recently, an updated version of MB-pol, MB-pol(2023), trained on larger training sets of CCSD(T) many-body energies, has been shown to achieve even higher predictive accuracy for simulations of water in both gas and liquid phases.<sup>61</sup>

In this study, we report the entropy values of liquid water at ambient conditions for the MB-pol and MB-pol(2023) potentials as predicted by the 2PT model and compare them with values calculated with several common pairwise additive water models. Since the 2PT model involves the calculation of the vibrational density of states, which depends sensitively on the water model used in the simulations, we also investigate the ability of empirical pairwise-additive models, as well as the MB-pol and MB-pol(2023) potentials, to consistently describe both the entropy and the structural and dynamical properties of liquid water at ambient conditions. Our analyses indicate that the entropy values calculated with both MB-pol and MB-pol(2023) are about 10% smaller than the corresponding experimental value. Since these data-driven many-body potentials provide remarkable agreement with experimental data across all water's phases, the discrepancy found for the entropy of liquid water might primarily be due to the approximate nature of the 2PT model. Similar discrepancies are also found for empirical pairwise-additive models (e.g., TIP4P/2005 and TIP4P/Ew) that are known to provide a robust description of the properties of water at ambient conditions, supporting the hypothesis that these discrepancies reflect inherent limitations of the 2PT model. This is further corroborated by the good performance in reproducing the experimental entropy value exhibited by pairwise-additive water models (e.g., TIP3P and SPC) that provide a poor description of both structural and dynamical properties of liquid water. This apparent agreement thus hinges on fortuitous error compensation between the approximations inherent to the 2PT model and the known inaccuracies of these pairwise-additive water models in correctly representing many-body interactions. Nevertheless, given its simplicity, the 2PT model can still serve as a valuable tool for estimating relative changes in the entropy of liquid water across various environments, especially when applied to simulations carried out with models that provide a reliable representation of the structural, thermodynamic, and dynamical properties of water across all its phases.

#### Methods

We consider water models belonging to the SPCx (SPC, <sup>23</sup> SPC/E, <sup>24</sup> SPC/Fw<sup>30</sup>) and TIPnP (TIP3P, <sup>25</sup> TIP4P, <sup>25</sup> TIP4P/2005, <sup>27</sup> TIP4P/2005f, <sup>29</sup> TIP4P/Ew, <sup>26</sup> TIP5P<sup>28</sup>) families of pairwise-additive water models along with the MB-pol<sup>33–35</sup> and MB-pol(2023)<sup>61</sup> data-driven many-body potentials. All MD simulations were performed using the Large-scale Atomic/Molecular Massively Parallel Simulator (LAMMPS)<sup>62</sup> for a cubic box of N = 256 water molecules in periodic boundary conditions. In the case of MB-pol and MB-pol(2023), the simulations were enabled by the MBX C++ library<sup>63,64</sup> interfaced with LAMMPS.

The equations of motion were propagated using the velocity-Verlet algorithm with a time step of 0.2 fs in the canonical (NVT: constant number of molecules N, volume V, and temperature T) ensemble,<sup>65</sup> with the temperature controlled using a global Nosé-Hoover chain consisting of four thermostats.<sup>66</sup> The nonbonded interactions were truncated at an atom-atom distance of 9.0 Å. The long-range electrostatic interactions for models belonging to the SPCx and TIPnP families were calculated using the particle-particle particle-mesh solver as implemented in LAMMPS. The long-range electrostatic interactions for MB-pol and MB-pol(2023) were calculated using the particle mesh Ewald method as implemented in MBX.<sup>63,67,68</sup> The entropy values of all water models were calculated using the 2PT model<sup>14</sup> by averaging over 20 independent NVT trajectories, each 50 ps long.

The tetrahedral order parameter,  $q_{tet}$ , was calculated according to<sup>69</sup>

$$q_{tet} = 1 - \frac{3}{8} \sum_{j=1}^{3} \sum_{k=j+1}^{4} \left( \cos \psi_{ijk} + \frac{1}{3} \right)^2, \tag{1}$$

where  $\psi_{ijk}$  is the angle between the oxygen atom of the central water molecule with index *i* and the oxygen atoms of two neighboring water molecules with index *j* and *k* lying within a distance smaller than 3.5 Å from the central molecule.  $q_{tet} = 0$  corresponds to a completely disordered structural arrangement as in an ideal gas, while  $q_{tet} = 1$  corresponds to a perfect tetrahedral arrangement. The orientational correlation function was calculated by averaging over the same 20 trajectories according to the following expression<sup>70</sup>

$$C_2(t) = \langle P_2[\vec{e}(0) \cdot \vec{e}(t)] \rangle.$$
<sup>(2)</sup>

Here,  $\vec{e}$  represents a unit vector along one of the OH bonds of a water molecule,  $P_2$  is the second-order Legendre polynomial, and the bracket indicates an ensemble average over all OH bonds at a given time t.

#### **Results and Discussions**

The entropy values calculated with all water models considered in this study are listed in Table 1 along with the experimental value. The entropy values obtained for SPC, SPC/E, SPC/Fw, TIP3P, TIP4P/2005, and TIP4P/Ew agree with those reported in previous studies.<sup>14,15</sup> SPC and TIP3P display the closest agreement with the experimental value within the SPCx and TIPnP families of models, respectively. All other SPC-type and TIPnP-type models, on the other hand, predict similar entropy values that are consistently smaller than the experimental value. In this regard, it should be noted that, despite water being empirical and described by effective two-body potential energy terms, TIP4P and TIP5P provide slightly better agreement with the experimental value than MB-pol and MB-pol(2023), which

Table 1: Entropy of liquid water at 298 K calculated using the 2PT model and MD simulations carried out at 298 K with the MB-pol and MB-pol(2023) data-driven many-body potentials, as well as various pairwise-additive models belonging to the SPCx and TIPnP families.

Model	Entropy (J/mol-K)
Experiment <sup>71</sup>	$69.95\pm0.03$
Data-driven many-body potentials	
MB-pol MB-pol(2023)	$62.9 \pm 0.3$ $62.0 \pm 0.5$
MD-p01(2023)	$02.0 \pm 0.0$
SPCx pairwis	e-additive models
SPC	$66.9\pm0.3$
SPC/E	$61.7\pm0.3$
SPC/Fw	$61.4 \pm 0.4$
TID. D	a addition madala
TIPNP pairw	ise-additive models
TIP3P	$70.6 \pm 0.2$
TIP4P	$64.3 \pm 0.2$
TIP4P/2005	$58.3 \pm 0.2$
TIP4P/2005f	$60.2\pm0.2$
TIP4P/Ew	$59.9\pm0.2$
TIP5P	$63.5\pm0.3$

were developed based on a rigorous many-body formalism and CCSD(T) reference energies.

The differences in entropy values among the water models are related to the differences in structural and dynamical properties of liquid water predicted by the models themselves. To gain insights into the relationship between the structure of water and the calculated entropy values, Figure 1 shows the oxygen-oxygen radial distribution function (RDF) of liquid water calculated with each model at 298 K and experimental density. While the SPC and TIP3P models provide the closest agreement with the experimental data, their RDFs lack prominent long-range structural organization, which leads to an overall poor representation of the overall structure of liquid water. On the other hand, all other models predict similar RDFs, correctly predicting the presence of well-defined first and second solvation shells. In particular, MB-pol and MB-pol(2023), which predict similar entropy values, provide the



Figure 1: Oxygen-oxygen radial distribution function,  $g(R_{OO})$ , calculated from MD simulations of liquid water carried out at 298 K with the MB-pol and MB-pol(2023) data-driven many-body potentials, as well as various pairwise-additive models belonging to the SPCx (A) and TIPnP (B) families. The experimental oxygen-oxygen radial distribution function from Ref. 72 is also shown for comparison.

closest agreement with the experimental oxygen-oxygen RDF.

It should be noted that the trend in entropy values predicted by the various water models qualitatively follows the structural ordering inferred from the analysis of the corresponding oxygen-oxygen RDFs. In particular, the high entropy values predicted by the SPC and TIP3P models are reflected in the absence of a second solvation shell in their oxygenoxygen RDFs. As shown in Table 1, the SPC/E, SPC/Fw, TIP4P/2005, TIP4P/2005f, and TIP4P/Ew models predict smaller entropy values than MB-pol and MB-pol(2023). This difference correlates with the first hydration shell in the oxygen-oxygen RDFs calculated with these models being overstructured compared to that predicted by MB-pol and MB-pol(2023). Overall, the trend in entropy values displayed by the various water models appears to be inversely correlated with their ability to predict the correct structure of liquid water. This analysis suggests that the better agreement of the SPC and TIP3P models with the experimental entropy value, compared to all other models, is likely due to error compensation between the approximations adopted by the 2PT model and the intrinsic limitations of these water models in correctly representing many-body interactions in liquid water. Additionally, based on the consistent agreement of MB-pol and MB-pol(2023) with experimental data for several structural, thermodynamic, and dynamical properties of water, it seems reasonable to assume that, due to its inherent approximate nature, the 2PT model likely underestimates the entropy of liquid water at ambient conditions predicted by these two models by approximately 10%.

The relationship between the structure and the entropy of water is further supported by the analysis of the tetrahedral order parameter  $(q_{tet})$  for each water model, as shown in Figure 2. Except for SPC and TIP3P, all water models exhibit a similar bimodal distribution, characterized by a shallow peak at  $q_{tet} \approx 0.5$  and a pronounced peak at  $q_{tet} \approx 0.8$ . The  $P(q_{tet})$  distributions calculated with the SPC and TIP3P models instead display a single peak at  $q_{tet} \approx 0.5$ . The analyses of the RDFs (Figure 1) and  $q_{tet}$  distributions (Figure 2) demonstrate that the SPC and TIP3P models are unable to accurately predict the structure of both the first and second solvation shells, despite providing the closest agreement with the experimental entropy values. Moreover, the larger entropy values predicted by the SPC and TIP3P models compared to other models qualitatively align with their  $P(q_{tet})$  distributions, since both models provide more pronounced distributions at smaller  $q_{tet}$  values, directly correlating with more disordered structural arrangements. In addition, the slightly higher



Figure 2: Probability distribution of the tetrahedral order parameter,  $P(q_{tet})$ , calculated from MD simulations of liquid water carried out at 298 K with the MB-pol and MB-pol(2023) data-driven many-body potentials, as well as various pairwise-additive models belonging to the SPCx (A) and TIPnP (B) families.

entropy value predicted by TIP4P compared to MB-pol and MB-pol(2023) can be understood in terms of the differences in the corresponding  $P(q_{tet})$  distributions. Specifically, the distribution calculated with TIP4P is noticeably higher at  $q_{tet} \approx 0.5$  and lower at  $q_{tet} \approx 0.8$ compared to the distributions calculated with MB-pol and MB-pol(2023), indicating slightly more disordered arrangements of water molecules residing within the first solvation shell.

Insights into the relationship between entropy values and dynamical properties predicted by the different water models can be gained by analyzing the calculated orientational correlation functions, as shown in Figure 3. Following the same trends displayed by the entropy values, RDFs, and  $P(q_{tet})$  distributions, all water models, except for SPC, TIP3P, and TIP4P, predict a similar decay of the orientational correlation functions as a function of time. The orientational relaxation times for all water models,  $\tau_2$ , which were determined by fitting the long-time decay of the corresponding orientational correlation functions to a single exponential function, are listed in Table 2. Among all models considered in this study, MB-pol and SPC/E provide the closest agreement with the experimental value, followed by MB-pol(2023), SPC/Fw, TIP4P/2005f, TIP4P/Ew, and TIP5P. Interestingly, while MB-pol and MB-pol(2023) effectively display the same ability to predict the structural properties



Figure 3: Orientational correlation function,  $C_2(t)$ , calculated from MD simulations of liquid water carried out at 298 K with the MB-pol and MB-pol(2023) data-driven many-body potentials, as well as various pairwise-additive models belonging to the SPC (A) and TIPnP (B) families.

Table 2: Orientational relaxation time,  $\tau_2$ , calculated from MD simulations of liquid water carried out at 298 K with the MB-pol and MB-pol(2023) data-driven many-body potentials, as well as various pairwise-additive models belonging to the SPC (A) and TIPnP (B) families. For of each water mode,  $\tau_2$  was calculated by fitting the long-time decay of the corresponding orientational correlation function to an exponential function:  $C_2(t) = A \exp(-t/\tau_2)$ .

Model	$ au_2  ext{ (ps)}$	
Experiment <sup>73</sup>	2.5	
<b>Data-driven ma</b> MB-pol MB-pol(2023)	$\begin{array}{l} \textbf{my-body potentials}\\ 2.6 \pm 0.1\\ 2.8 \pm 0.1 \end{array}$	
SPCx pairwise-additive models		
SPC	$1.5 \pm 0.1$	
SPC/E	$2.6 \pm 0.1$	
SPC/Fw	$3.0 \pm 0.1$	
TIPnP pairwise-additive models		
TIP3P	$1.1 \pm 0.1$	
TIP4P	$1.8 \pm 0.1$	
TIP4P/2005	$3.1 \pm 0.1$	
TIP4P/2005f	$2.7\pm0.1$	
TIP4P/Ew	$2.7\pm0.1$	
TIP5P	$2.3 \pm 0.1$	

of liquid water, as shown in Figures 1 and 2, they provide slightly different orientational correlation functions. In addition, it is noteworthy that models with smaller  $\tau_2$  values generally exhibit larger entropy values, as shown in Table 1. Smaller  $\tau_2$  values, corresponding to rapid decays of the corresponding correlation functions, indicate that water molecules more quickly lose memory of their initial orientations. Consequently, information about the water structure is more rapidly lost over time. In this context, the trends in  $\tau_2$  and entropy values among all water models suggest that, despite being inherently approximate, the 2PT model qualitatively captures the essential features of entropy, directly correlating the entropy value with the extent of structural order/disorder and information content.

### Conclusions

In this study, we have investigated the entropy of liquid water as predicted by the 2PT model used in combination with various models, including the MB-pol and MB-pol(2023) data-driven many-body potentials and models belonging to the SPCx and TIPnP families. We have demonstrated that the entropy values calculated using the 2PT model qualitatively correlate with the extent of structural ordering and fast dynamics predicted by a given water model. However, our analyses indicate that the 2PT model does not consistently align the entropy value calculated with a given water model with the ability of the same model to reproduce other structural, thermodynamic, and dynamical properties of liquid water. This discrepancy can be attributed to the interplay between the approximations adopted by the 2PT model and the ability of a water model to accurately describe many-body interactions in water. This is evidenced by the ability of the SPC and TIP3P models to predict entropy values in remarkable agreement with the experimental value while providing a poor description of other structural and dynamical properties. On the other hand, the MB-pol and MB-pol(2023) data-driven many-body potentials consistently provide excellent agreement with experimental data for water across all phases but, when used in 2PT calculations, they underestimate the entropy value by approximately 10%. Overall, our study suggests that, despite its inherent approximate nature, the 2PT model can serve as an efficient tool to estimate relative entropy changes of water in different environments when it is used in combination with water models that provide a robust description of both the structure and dynamics of liquid water.

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#### Data availability

Any data generated and analyzed for this study that are not included in this article are available from the authors upon request.

#### Code availability

The molecular models used in the MD simulations are publicly available on GitHub (https://github.com/paesanilab/Data\_Repository/tree/main/water\_entropy) in the format for LAMMPS or the MBX<sup>64</sup> interface with LAMMPS.<sup>62</sup> All computer codes used in the analysis presented in this study are available from the authors upon request.

## References

- (1) Wehrl, A. General Properties of Entropy. Rev. Mod. Phys. 1978, 50, 221.
- (2) Shannon, C. E. A Mathematical Theory of Communication. Bell Syst. Tech. J. 1948, 27, 379–423.
- (3) Georgescu-Roegen, N. The Entropy Law and the Economic Process; Harvard university press, 1971.
- (4) Brooks, D. R.; Wiley, E. O.; Brooks, D. Evolution as Entropy; University of Chicago Press Chicago, 1988.

- (5) Fermi, E. *Thermodynamics*; Courier Corporation, 2012.
- (6) Bernstein, R.; Levine, R. Entropy and Chemical Change. I. Characterization of Product (and Reactant) Energy Distributions in Reactive Molecular Collisions: Information and Entropy Deficiency. J. Chem. Phys. 1972, 57, 434–449.
- (7) Ben-Shaul, A.; Levine, R.; Bernstein, R. Entropy and Chemical Change. II. Analysis of Product Energy Distributions: Temperature and Entropy Deficiency. J. Chem. Phys. 1972, 57, 5427–5447.
- (8) Alhassid, Y.; Levine, R. Entropy and Chemical Change. III. The Maximal Entropy (Subject to Constraints) Procedure as a Dynamical Theory. J. Chem. Phys. 1977, 67, 4321–4339.
- (9) Dill, K. A.; MacCallum, J. L. The Protein-Folding Problem, 50 Years On. Science 2012, 338, 1042–1046.
- (10) Stanley, H. E. Phase Transitions and Critical Phenomena; Clarendon Press, Oxford, 1971; Vol. 7.
- (11) George, E. P.; Raabe, D.; Ritchie, R. O. High-Entropy Alloys. Nat. Rev. Mater. 2019, 4, 515–534.
- (12) Lebowitz, J. L. Boltzmann's Entropy and Time's Arrow. *Phys. Today* **1993**, *46*, 32–38.
- (13) Lin, S.-T.; Blanco, M.; Goddard III, W. A. The Two-Phase Model for Calculating Thermodynamic Properties of Liquids from Molecular Dynamics: Validation for the Phase Diagram of Lennard-Jones Fluids. J. Chem. Phys. 2003, 119, 11792–11805.
- (14) Lin, S.-T.; Maiti, P. K.; Goddard III, W. A. Two-phase Thermodynamic Model for Efficient and Accurate Absolute Entropy of Water from Molecular Dynamics Simulations. J. Phys. Chem. B 2010, 114, 8191–8198.

- (15) Pascal, T. A.; Schärf, D.; Jung, Y.; Kühne, T. D. On the Aabsolute Thermodynamics of Water from Computer Simulations: A Comparison of First-Principles Molecular Dynamics, Reactive and Empirical Force Fields. J. Chem. Phys. 2012, 137, 244507.
- (16) Bagchi, B. Water in Biological and Chemical Processes: From Structure and Dynamics to Function; Cambridge University Press, 2013.
- (17) Guillot, B. A Reappraisal of What We Have Learnt During Three Decades of Computer Simulations on Water. J. Mol. Liq. 2002, 101, 219–260.
- (18) Vega, C.; Abascal, J. L. Simulating Water with Rigid Non-Polarizable Models: A General Perspective. *Phys. Chem. Chem. Phys.* **2011**, *13*, 19663–19688.
- (19) Shvab, I.; Sadus, R. J. Atomistic Water Models: Aqueous Thermodynamic Properties from Ambient to Supercritical Conditions. *Fluid Phase Equilib.* **2016**, 407, 7–30.
- (20) Cisneros, G. A.; Wikfeldt, K. T.; Ojamäe, L.; Lu, J.; Xu, Y.; Torabifard, H.; Bartók, A. P.; Csányi, G.; Molinero, V.; Paesani, F. Modeling Molecular Interactions in Water: From Pairwise to Many-Body Potential Energy Functions. *Chem. Rev.* 2016, 116, 7501–7528.
- (21) Barker, J.; Watts, R. Structure of Water; A Monte Carlo Calculation. Chem. Phys. Lett. 1969, 3, 144–145.
- (22) Rahman, A.; Stillinger, F. H. Molecular Dynamics Study of Liquid Water. J. Chem. Phys. 1971, 55, 3336–3359.
- (23) Berendsen, H. J.; Postma, J. P.; van Gunsteren, W. F.; Hermans, J. Intermolecular Forces; Springer, 1981; pp 331–342.
- (24) Berendsen, H.; Grigera, J.; Straatsma, T. The Missing Term in Effective Pair Potentials.
   J. Phys. Chem. 1987, 91, 6269–6271.

- (25) Jorgensen, W. L.; Chandrasekhar, J.; Madura, J. D.; Impey, R. W.; Klein, M. L. Comparison of Simple Potential Functions for Simulating Liquid Water. J. Chem. Phys. 1983, 79, 926–935.
- (26) Horn, H. W.; Swope, W. C.; Pitera, J. W.; Madura, J. D.; Dick, T. J.; Hura, G. L.; Head-Gordon, T. Development of an Improved Four-site Water Model for Biomolecular Simulations: TIP4P-Ew. J. Chem. Phys. 2004, 120, 9665–9678.
- (27) Abascal, J. L. F.; Vega, C. A General Purpose Model for the Condensed Phases of Water: TIP4P/2005. J. Chem. Phys. 2005, 123, 234505.
- (28) Mahoney, M. W.; Jorgensen, W. L. A Five-Site Model for Liquid Water and the Reproduction of the Density Anomaly by Rigid, Nonpolarizable Potential Functions. J. Chem. Phys. 2000, 112, 8910–8922.
- (29) González, M. A.; Abascal, J. L. A Flexible Model for Water Based on TIP4P/2005. J. Chem. Phys. 2011, 135, 224516.
- (30) Wu, Y.; Tepper, H. L.; Voth, G. A. Flexible Simple Point-Charge Water Model with Improved Liquid-State Properties. J. Chem. Phys. 2006, 124, 024503.
- (31) Stanton, J. F. Why CCSD(T) Works: A Different Perspective. Chem. Phys. Lett. 1997, 281, 130–134.
- (32) Rezac, J.; Hobza, P. Benchmark Calculations of Interaction Energies in Noncovalent Complexes and Their Applications. *Chem. Rev.* 2016, 116, 5038–5071.
- (33) Babin, V.; Leforestier, C.; Paesani, F. Development of a "First Principles" Water Potential with Flexible Monomers: Dimer Potential Energy Surface, VRT Spectrum, and Second Virial Coefficient. J. Chem. Theory Comput. 2013, 9, 5395–5403.
- (34) Babin, V.; Medders, G. R.; Paesani, F. Development of a "First Principles" Water

Potential with Flexible Monomers. II: Trimer Potential Energy Surface, Third Virial Coefficient, and Small Clusters. J. Chem. Theory Comput. **2014**, 10, 1599–1607.

- (35) Medders, G. R.; Babin, V.; Paesani, F. Development of a "First Principles" Water Potential with Flexible Monomers. III. Liquid Phase Properties. J. Chem. Theory Comput. 2014, 10, 2906–2910.
- (36) Paesani, F. Getting the Right Answers for the Right Reasons: Toward Predictive Molecular Simulations of Water with Many-Body Potential Energy Functions. Acc. Chem. Res. 2016, 49, 1844–1851.
- (37) Reddy, S. K.; Straight, S. C.; Bajaj, P.; Pham, C. H.; Riera, M.; Moberg, D. R.; Morales, M. A.; Knight, C.; Götz, A. W.; Paesani, F. On the Accuracy of the MB-pol Many-Body Potential for Water: Interaction Energies, Vibrational Frequencies, and Classical Thermodynamic and Dynamical Properties from Clusters to Liquid Water and Ice. J. Chem. Phys. 2016, 145, 194504.
- (38) Bore, S. L.; Paesani, F. Realistic Phase Diagram of Water from "First Principles" Data-driven Quantum Simulations. *Nat. Commun.* 2023, 14, 3349.
- (39) Bajaj, P.; Götz, A. W.; Paesani, F. Toward Chemical Accuracy in the Description of Ion–Water Interactions through Many-Body Representations. I. Halide–Water Dimer Potential Energy Surfaces. J. Chem. Theory Comput. 2016, 12, 2698–2705.
- (40) Riera, M.; Mardirossian, N.; Bajaj, P.; Götz, A. W.; Paesani, F. Toward Chemical Accuracy in the Description of Ion–Water Interactions through Many-Body Representations. Alkali-Water Dimer Potential Energy Surfaces. J. Chem. Phys. 2017, 147, 161715.
- (41) Riera, M.; Brown, S. E.; Paesani, F. Isomeric Equilibria, Nuclear Quantum Effects, and Vibrational Spectra of M<sup>+</sup>(H<sub>2</sub>O)<sub>n=1-3</sub> Clusters, with M = Li, Na, K, Rb, and Cs, through Many-Body Representations. J. Phys. Chem. A 2018, 122, 5811–5821.

- (42) Bajaj, P.; Wang, X.-G.; Carrington Jr, T.; Paesani, F. Vibrational Spectra of Halide– Water Dimers: Insights on Ion Hydration from Full-Dimensional Quantum Calculations on Many-Body Potential Energy Surfaces. J. Chem. Phys. 2018, 148, 102321.
- (43) Bajaj, P.; Richardson, J. O.; Paesani, F. Ion-Mediated Hydrogen-Bond Rearrangement through Tunnelling in the Iodide–Dihydrate Complex. *Nat. Chem.* 2019, 11, 367–374.
- (44) Bajaj, P.; Zhuang, D.; Paesani, F. Specific Ion Effects on Hydrogen-Bond Rearrangements in the Halide–Dihydrate Complexes. J. Phys. Chem. Lett. 2019, 10, 2823–2828.
- (45) Bajaj, P.; Riera, M.; Lin, J. K.; Mendoza Montijo, Y. E.; Gazca, J.; Paesani, F. Halide Ion Microhydration: Structure, Energetics, and Spectroscopy of Small Halide–Water Clusters. J. Phys. Chem. A 2019, 123, 2843–2852.
- (46) Riera, M.; Talbot, J. J.; Steele, R. P.; Paesani, F. Infrared Signatures of Isomer Selectivity and Symmetry Breaking in the Cs<sup>+</sup>(H<sub>2</sub>O)<sub>3</sub> Complex Using Many-Body Potential Energy Functions. J. Chem. Phys. **2020**, 153, 044306.
- (47) Caruso, A.; Paesani, F. Data-Driven Many-Body Models Enable a Quantitative Description of Chloride Hydration from Clusters to Bulk. J. Chem. Phys. 2021, 155, 064502.
- (48) Caruso, A.; Zhu, X.; Fulton, J. L.; Paesani, F. Accurate Modeling of Bromide and Iodide Hydration with Data-Driven Many-Body Potentials. J. Phys. Chem. B 2022, 126, 8266–8278.
- (49) Zhuang, D.; Riera, M.; Schenter, G. K.; Fulton, J. L.; Paesani, F. Many-Body Effects Determine the Local Hydration Structure of Cs<sup>+</sup> in Solution. J. Phys. Chem. Lett. 2019, 10, 406–412.
- (50) Zhuang, D.; Riera, M.; Zhou, R.; Deary, A.; Paesani, F. Hydration Structure of Na<sup>+</sup>

and K<sup>+</sup> Ions in Solution Predicted by Data-Driven Many-Body Potentials. J. Phys. Chem. B **2022**, 126, 9349–9360.

- (51) Savoj, R.; Agnew, H.; Zhou, R.; Paesani, F. Molecular Insights into the Influence of Ions on the Water Structure. I. Alkali Metal Ions in Solution. J. Phys. Chem. B 2024, 128, 1953–1962.
- (52) Rieth, A. J.; Hunter, K. M.; Dincă, M.; Paesani, F. Hydrogen Bonding Structure of Confined Water Templated by a Metal-Organic Framework With Open Metal Sites. *Nat. Commun.* 2019, 10, 1–7.
- (53) Hunter, K. M.; Wagner, J. C.; Kalaj, M.; Cohen, S. M.; Xiong, W.; Paesani, F. Simulation Meets Experiment: Unraveling the Properties of Water in Metal–Organic Frameworks through Vibrational Spectroscopy. J. Phys. Chem. C 2021, 125, 12451–12460.
- (54) Wagner, J. C.; Hunter, K. M.; Paesani, F.; Xiong, W. Water Capture Mechanisms at Zeolitic Imidazolate Framework Interfaces. J. Am. Chem. Soc. 2021, 143, 21189–21194.
- (55) Ho, C.-H.; Paesani, F. Elucidating the Competitive Adsorption of H<sub>2</sub>O and CO<sub>2</sub> in CALF-20: New Insights for Enhanced Carbon Capture Metal–Organic Frameworks. ACS Appl. Mater. Interfaces **2023**, 15, 48287–48295.
- (56) Ho, C.-H.; Valentine, M. L.; Chen, Z.; Xie, H.; Farha, O.; Xiong, W.; Paesani, F. Structure and Thermodynamics of Water Adsorption in NU-1500-Cr. Commun. Chem. 2023, 6, 70.
- (57) Zhang, J.; Paesani, F.; Lessio, M. Computational Insights into the Interaction of Water with the UiO-66 Metal–Organic Framework and its Functionalized Derivatives. J. Mater. Chem. C 2023, 11, 10247–10258.
- (58) Frank, H. O.; Paesani, F. Molecular Driving Forces for Water Adsorption in MOF-808:
   A Comparative Analysis with UiO-66. J. Chem. Phys. 2024, 160, 094703.

- (59) Hunter, K. M.; Paesani, F. Monitoring Water Harvesting in Metal–Organic Frameworks, One Water Molecule at a Time. *Chem. Sci.* 2024, 15, 5303–5310.
- (60) Oppenheim, J. J.; Ho, C.-H.; Alezi, D.; Andrews, J. L.; Chen, T.; Dinakar, B.; Paesani, F.; Dincă, M. Cooperative Interactions with Water Drive Hysteresis in a Hydrophilic Metal–Organic Framework. *Chem. Mater.* **2024**, *36*, 3395–3404.
- (61) Zhu, X.; Riera, M.; Bull-Vulpe, E. F.; Paesani, F. MB-pol (2023): Sub-chemical Accuracy for Water Simulations from the Gas to the Liquid Phase. J. Chem. Theory Comput. 2023, 19, 3551–3566.
- (62) Thompson, A. P.; Aktulga, H. M.; Berger, R.; Bolintineanu, D. S.; Brown, W. M.; Crozier, P. S.; in 't Veld, P. J.; Kohlmeyer, A.; Moore, S. G.; Nguyen, T. D.; Shan, R.; Stevens, M. J.; Tranchida, J.; Trott, C.; Plimpton, S. J. LAMMPS-a Flexible Simulation Tool for Particle-based Materials Modeling at the Atomic, Meso, and Continuum Scales. *Comput. Phys. Commun.* **2022**, *271*, 108171.
- (63) Riera, M.; Knight, C.; Bull-Vulpe, E. F.; Zhu, X.; Agnew, H.; Smith, D. G. A.; Simmonett, A. C.; Paesani, F. MBX: A Many-Body Energy and Force Calculator for Data-Driven Many-Body Simulations. J. Chem. Phys. 2023, 159, 054802.
- (64) MBX v1.0. http://paesanigroup.ucsd.edu/software/mbx.html, 2023.
- (65) Shinoda, W.; Shiga, M.; Mikami, M. Rapid Estimation of Elastic Constants by Molecular Dynamics Simulation under Constant Stress. *Phys. Rev. B* 2004, 69, 134103.
- (66) Martyna, G. J.; Klein, M. L.; Tuckerman, M. Nosé–Hoover Chains: The Canonical Ensemble Via Continuous Dynamics. J. Chem. Phys. 1992, 97, 2635–2643.
- (67) Simmonett, A. C.; Brooks, B. R. Analytical Hessians for Ewald and Particle Mesh Ewald Electrostatics. J. Chem. Phys. 2021, 154, 104101.

- (68) Simmonett, A. C.; Brooks, B. R. A Compression Strategy for Particle Mesh Ewald Theory. J. Chem. Phys. 2021, 154, 054112.
- (69) Errington, J. R.; Debenedetti, P. G. Relationship between Structural Order and the Anomalies of Liquid Water. *Nature* 2001, 409, 318–321.
- (70) Bakker, H.; Skinner, J. Vibrational Spectroscopy as a Probe of Structure and Dynamics in Liquid Water. *Chem. Rev.* 2010, *110*, 1498–1517.
- (71) NIST. Chemistry WebBook, Reference Database Number 69 ed.; National Institute of Standards and Technology. 2020.
- (72) Skinner, L. B.; Huang, C.; Schlesinger, D.; Pettersson, L. G.; Nilsson, A.; Benmore, C. J. Benchmark Oxygen-Oxygen Pair-Distribution Function of Ambient Water from X-ray Diffraction Measurements with a Wide Q-Range. J. Chem. Phys. 2013, 138, 074506.
- (73) Rezus, Y.; Bakker, H. On the Orientational Relaxation of HDO in Liquid Water. J. Chem. Phys. 2005, 123, 114502.



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