Stable and Accurate Atomistic Simulations of Flexible Molecules using Conformationally Generalisable Machine Learned Potentials

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Abstract

Computational simulation methods based on machine learned potentials (MLPs) promise to revolutionise shape prediction of flexible molecules in solution, but their widespread adoption has been limited by the way in which training data is generated. Here, we present an approach which allows the key conformational degrees of freedom to be properly represented in reference molecular datasets. MLPs trained on these datasets using a global descriptor scheme are generalisable in conformational space, providing quantum chemical accuracy for all conformers. These MLPs are capable of propagating long, stable molecular dynamics trajectories, an attribute that has remained a challenge for MLPs. We deploy the MLPs in obtaining converged conformational free energy surfaces for flexible molecules via well-tempered metadynamics simulations; this approach provides a hitherto inaccessible route to accurately computing the structural, dynamical and thermodynamical properties of a wide variety of flexible molecular systems.

1. Introduction

Few areas of the physical and biological sciences remain untouched by the valuable contributions of atomistic simulation techniques, either based on ab initio quantum mechanics (QM)¹ or empirical potentials.² Despite the many compelling structural, kinetic and thermodynamic insights obtained from these two well-established approaches, the inherent limitations are well understood. Quantum chemical methods offer accuracy but are bounded by computational inefficiency and their use quickly becomes intractable with increasing system size or time scale. Empirical potentials are highly efficient to compute; however, the use of simple physics-based functional forms to represent the potential energy surface can neglect or misrepresent important quantum mechanical effects, with inaccurate property prediction across diverse fields as a result.³⁻⁹ The emergence of machine learned interatomic potentials (MLPs) offers a possible solution to bridge this accuracy-efficiency gap. 10-14 MLPs are trained to learn the complex mapping between a system's atomic structure, encoded by carefully chosen input features, and its multidimensional potential energy surface, using regression algorithms. Extremely flexible functional forms of the MLP enable much higher accuracy than can be achieved using empirical potentials, providing that a sufficiently high-quality ab initio QM reference dataset is available for training. The capability to perform simulations with forces and energies at quantum chemical accuracy without the prohibitive computational cost is revolutionising the atomistic modelling toolkit for applications spanning materials science, 15, ¹⁶ chemistry¹⁷ and biology. ¹⁸

For flexible molecular systems, which can have complex conformational energy surfaces defined by torsional motions around one or more interdependent rotatable bonds, the accuracy-efficiency gap is particularly problematic.¹⁹ An interplay of many competing effects, such as conjugation, steric repulsion, hydrogen bonding, dispersion, electronic repulsion and solvation can lead to subtle differences in relative conformational energies. Many of these effects are not well represented by conventional empirical potential methods, sometimes resulting in a failure to correctly identify the lowest energy conformer.^{5,7} On the other hand, the timescales involved in torsional motions are often too long to observe conformational changes using *ab initio* molecular dynamics (AIMD) methods due to the requirement to overcome large activation energy barriers. One area in which this trade-off impedes progress is drug design, where it is critical for atomistic simulations to be able to accurately distinguish between the energies of the unbound and protein-bound conformational states of flexible drug molecules, as well as conduct long simulations, to sample a representative ensemble and thus

correctly assess their binding affinities. In this particular field, MLPs may confer significant advantages over conventional techniques.^{20, 21}

The general training protocol for MLPs is well-established but the last few years have seen the algorithmic variety of schemes in the literature rapidly evolve, ^{22, 23} including those based on feed-forward, 24-27 message-passing 28, 29 and equivariant 30-33 neural networks, Gaussian process regression, 34-36 kernel ridge regression 37-39 and atomic cluster expansions. 40-⁴² In spite of this important progress, the widespread adoption of MLPs in molecular simulation requires several remaining challenges to be surmounted.^{39, 43, 44} One of these challenges pertains to trajectories that are unstable over the long timeframes required for sampling the underlying probability distribution to calculate meaningful simulation observables. These instabilities can arise from unphysical configurations in regions of the potential energy surface that were poorly sampled in the reference dataset used for training, 45, 46 where there is no guarantee that MLPs can extrapolate to predict physically reasonable forces. It must therefore be ensured that reference datasets contain structures corresponding to all relevant local minima, as well as those corresponding to transition paths, to prevent unstable trajectories. This challenge is particularly acute for flexible molecules which may possess a distribution of distinct conformational states associated with separate energetic minima. ¹⁹ Developing models that generalise across all regions of conformational space is the necessary next step underpinning the widespread deployment of MLPs for flexible molecules. This naturally brings into focus the task of how to efficiently generate high quality reference datasets for training MLPs.

In contrast to the development of data efficient learning algorithms, comparatively little attention has been paid on how best to generate reference dataset structures. It is most common for structures to be sampled from equilibrium molecular dynamics (MD) simulations.^{37, 47, 48} Although elevated temperatures are frequently employed to extend the range of sampling, there is no guarantee that an arbitrary high temperature will be sufficient to surmount large activation energy barriers and sample all the relevant conformers of a flexible molecule. Another common approach is normal-mode sampling.^{24, 49-51} However, exhaustive normal-mode sampling, which depends on computing the Hessian matrix, may be a prohibitively costly overhead. In addition, the non-linear distortions involved in conformational change may be poorly approximated by linear normal-mode sampling. Other more sophisticated approaches, such as active learning,⁵²⁻⁵⁴ adaptive learning^{55, 56} and query-by-committee,⁵⁷ have been proposed to generate structures in poorly sampled regions. However, in many cases, these approaches are unnecessarily complex and computationally expensive, especially when the rotatable bonds of

interest are essentially already known. The use of enhanced sampling techniques,⁵⁸⁻⁶¹ which are well-established in MD simulations, could be used to improve sampling with little computational overhead instead. For example, Yang et al.⁶² recently showed that enhanced sampling could be used to generate the transition state structures needed to train and employ a reactive MLP to simulate urea decomposition.

In this study, we highlight that comprehensive conformational sampling, i.e. ensuring that all energetically relevant conformers are represented by structures in the reference dataset, is a crucial consideration when training MLPs for flexible molecules. We firstly demonstrate that the most popular reference dataset used for MLP benchmarking, the revised MD17 (rMD17) dataset,^{37, 63} which contains the gas-phase trajectories of several flexible drug molecules, is missing important conformers. Simulations of these flexible molecules using MLPs trained on the rMD17 dataset ultimately fail due to the generation of unphysical structures. We then propose a straightforward alternative scheme for generating robust reference datasets for training MLPs to remedy this shortcoming. Finally, we remark upon the implications of our findings in the training of MLPs for simulating flexible molecules and their stability in MD simulations.

2. Methods

2.1 Training scheme.

MLPs were trained according to a robust variant of the PairFE-Net global descriptor scheme,⁶⁴ which is based on the original PairF-Net formalism.²⁷ For certain applications, the use of global descriptors to encode atomic structure confers advantages over local descriptors. In avoiding the use of cut-offs or symmetry functions inherent to local descriptor models, all quantum chemical and long-range effects are included. The inability of local descriptor models to account for long-range interactions may have a significantly deleterious impact on the predictive ability of an MLP,⁶⁵⁻⁶⁸ potentially hindering its ability to distinguish between conformers that subtly differ in stability. Furthermore, although local descriptor models can be parallelised and designed with molecular generalisability and size-extensivity in mind,^{14, 24, 26, 35, 40} these attributes are not always necessary or desirable depending on the application at hand.²² Indeed, in drug design, Lipinski's "rule of 5" mnemonic⁶⁹ suggests that candidate ligands should have a molecular weight of less than 500 g mol⁻¹ in order to permit their penetration through cell membranes. This size regime is within the capacity of global descriptor models such as PairFE-Net. Furthermore, where quantum chemical accuracy is only required for small regions of a system, MLPs can be combined or embedded with much cheaper

empirical potentials, e.g. to model solvent or a protein binding pocket, in a manner analogous to hybrid quantum mechanical/molecular mechanical approaches, as implemented in our earlier work⁶⁴ and elsewhere.^{70, 71}

In PairFE-Net, an atomic structure is encoded using pairwise nuclear repulsion forces,

$$F_{ij}^{NR} = \frac{Z_i Z_j}{r_{ij}^2} \tag{1}$$

where Z_i is the nuclear charge and r_{ij} is the interatomic distance. This internal coordinate system guarantees rotational and translation invariance of the trained MLP. The use of the full set of pairwise nuclear repulsion forces as input features without cut-offs ensures that all long-range interactions are included. To train an MLP, the absolute energy of a structure, E, is first converted to an energy scaled over the range of forces, E^* , where

$$E^* = \frac{F_{i,k}^{max}(E - E_{min})}{(E_{max} - E_{min})}$$
 (2)

 $F_{i,k}^{max}$ is the maximum absolute force across all atoms i and Cartesian dimensions, k, and E_{max} and E_{min} are the maximum and minimum energies, across the combined training and validation sets. In our previous study,⁶⁴ forces and scaled energies were simultaneously decomposed into a set of pairwise interatomic coefficients using a transformation matrix. Here we employed a simplified version of the method, in which only the scaled energy is decomposed into a set of interatomic energies for the i and j atom pair, ε_{ij} , according to

$$E^* = \sum_{ij}^{N_{pair}} X_{ij} \varepsilon_{ij} \tag{3}$$

where N_{pair} is the number of distinct atom pairs in the system, given by

$$N_{pair} = \frac{N(N-1)}{2} \tag{4}$$

where N is the number of atoms in the system. X_{ij} is the pair energy bias, which depends only on their reciprocal interatomic distance,

$$X_{ij} = \frac{M}{r_{ij}} \tag{5}$$

and M is a normalization constant

$$M = \left(\sum_{ij}^{N_{pair}} \frac{1}{r_{ij}^2}\right)^{-\frac{1}{2}} \tag{6}$$

During training according to the PairFE-Net scheme, artificial neural networks directly predict the pairwise interatomic energies, $\hat{\varepsilon}_{ij}$, which are then recombined to predict the scaled energy, \hat{E}^* , according to Equation 2. The predicted energy is then unscaled using

$$\hat{E} = \frac{\hat{E}^* (E_{max} - E_{min})}{F_{i,k}^{max}} + E_{min}$$
 (7)

from which the predicted conservative atomic forces for each atom i and in each Cartesian direction k, $\hat{F}_{i,k}$, can be calculated from

$$\hat{F}_{i,k} = -\vec{\nabla}_i \hat{E} \tag{8}$$

ensuring that energy is strictly conserved.

2.2 Neural network training.

Reference datasets containing 10,000 structures were sub-divided into separate training, validation and test sets containing 8000, 1000 and 1000 structures, respectively. Feed-forward artificial neural networks were trained, using a batch size of 32 structures, to predict forces and energies by minimising the custom mean-squared error loss function, L,

$$L = \lambda_E \|E - \hat{E}\|^2 + \lambda_F \frac{1}{3N} \sum_{i=1}^{N} \sum_{j=1}^{k} \|F_{i,k} - \hat{F}_{i,k}\|^2$$
(9)

for the training set, where λ_E and λ_F are the (unscaled) energy and force weights, given by

$$\lambda_F = \frac{3N}{3N+1} \qquad \lambda_E = 1 - \lambda_F \tag{10}$$

Compared to our previous study,⁶⁴ the pairwise interatomic coefficients do not feature in the loss function, enabling improved force and energy prediction accuracy. The learning rate, initially set to 5×10^{-4} , was chosen on the basis of balancing the need to minimise training time and avoid exploding gradients. As well as being used to monitor for overfitting, the validation set was used to determine the convergence criteria of the training process. Specifically, if there was no improvement in the validation loss over the preceding 2000 epochs, the learning rate was reduced by a factor of 0.5, and training was stopped once the learning rate had decreased to 1×10^{-7} . All networks were comprised of an input layer containing N_{pair} nodes, three hidden layers each containing 360 nodes with sigmoid linear unit (SiLU) activation functions and a linear output layer containing N_{pair} nodes. The use of multiple hidden layers greatly enhances the ability of the trained networks to fit complex and highly non-linear functional dependences, improving their predictive capability compared to our previous study.⁶⁴ Network hyperparameters were chosen on the basis of systematic evaluation and optimisation of test set predictions using an independent molecule in the rMD17 training set (malonaldehyde).

Separate neural networks were trained for each molecule and dataset. All neural networks were constructed and trained using Tensorflow, version 2.12.⁷²

2.3 Dataset generation.

To sample structures in the reference datasets, molecular dynamics simulations were performed using the GAFF potential.⁷³ The simulations were initiated using a single molecule centred in a simulation cell with side lengths of 2.5 nm. Simulations were performed for 11 ns in the canonical ensemble, numerically integrating the equations of motion using the velocity Verlet algorithm and a 1 fs timestep. Target temperatures were maintained using the Nosé-Hoover chain thermostat⁷⁴ with a collision frequency of 50 ps⁻¹, a chain with 10 beads, 5 thermostat iterations and 5 Yoshida-Suzuki integration parameters. For each molecule, partial atomic charges were derived from the restrained electrostatic potential (RESP) approach⁷⁵ at the HF/6-31G* level of theory. Non-bonded interactions were calculated without cut-offs. In enhanced sampling simulations, the metadynamics⁵⁸ bias potential was constructed as the sum of penalty Gaussian functions with a fixed height, h, and width of 0.24 kcal mol⁻¹ and 0.35 radians, deposited every 0.5 ps. In all sampling simulations, structures were saved every 1 ps in the final 10 ns of simulation time. All MD simulations were carried out using the OpenMM library, ⁷⁶ version 8.0 and, in the case of enhanced sampling with metadynamics, using the OpenMM Plumed plugin.⁷⁷ The proportions of the (φ, ψ) -surface populated with structures shown in Table 3 were estimated using a 32×32 grid.

The 10,000 structures sampled from the empirical potential simulations were evaluated using single point calculations at the B3LYP/6-31G* level⁷⁸⁻⁸¹ to obtain reference datasets containing quantum chemical Cartesian atomic forces and energies. Torsional energy profiles for ϕ or ψ were obtained by constrained geometry optimisations. For each value of ϕ , a conformational search with respect to ψ was carried out to find the minimum energy, and vice versa. The torsion energy profiles were calculated at the B3LYP/6-31G* level⁷⁸⁻⁸¹ for consistency with the newly generated datasets and, separately, at the PBE/def2-SVP level for consistency with the rMD17 reference datasets, respectively. All DFT calculations employed Gaussian 09, revision D.01,⁸² using a pruned 99,590 point (ultrafine) integration grid and tight SCF convergence criteria of <10⁻⁸ RMS change in the density matrix, which corresponds to a change in energy of approximately 10^{-7} kcal mol⁻¹.

Due to the use of input descriptors based on an internal coordinate system which imposes an arbitrary order on input features, the PairFE-Net scheme is not invariant with respect to atomic permutation. When combined with inherently uneven MD sampling, this can result in significantly different energy and force predictions for symmetrically equivalent

structures and an inability of trained MLPs to correctly reproduce torsional energy profiles. One solution to this problem is augmenting the dataset with structural copies of all symmetrically equivalent structures, but this is unfeasible for all but the smallest molecules due to a combinatorial explosion in the number of permutations. Instead, after the DFT evaluation and prior to training MLPs, random swap moves were applied to all permutationally equivalent atoms or groups of atoms in the reference dataset. For each structure, 10 swap moves were applied to every symmetry-related functional group. This "permutational shuffling" of the reference dataset essentially eliminates the problem associated with a lack of permutational invariance combined with uneven MD sampling, enabling accurate prediction of torsional energy profiles.

2.4 MLP simulations.

In the MLP simulations, atomic forces were predicted from trained PairFE-Net neural networks. Constant temperatures were maintained using the Langevin thermostat with a coupling constant of 1 ps⁻¹. All other simulation parameters were as described above for the empirical potential simulations. Although the time that instability arises is obvious by visual analysis of MD trajectories, unstable structures were defined as those in which either (i) any bond deviates by \pm 0.25 Å relative to the initial (optimised) structure; or (ii) any interatomic distance is less than 0.75 Å. Stability was first assessed using a short 10 ps simulation, saving structures every 1 fs. On successful completion of this simulation, longer 25 ns simulations were performed, saving structures every 1 ps, while again checking for stability using the above criteria. Converged conformational free energy surfaces were obtained using 25 ns well-tempered metadynamics simulations. As opposed to the conventional metadynamics technique, a bias factor of 6.0 was used to scale the h, from an initial value of 0.24 kcal mol⁻¹.

3. Results

3.1 Benchmarking.

A crucial first step in the development of new MLPs is benchmarking with reference datasets. The rMD17 dataset^{37, 63} has been the subject of extensive testing and benchmarking and is by far the most popular benchmark dataset for MLPs. It contains 10 small organic molecules with structures sampled from short *ab initio* MD simulations in vacuum at 500 K. To assess the performance of PairFE-Net for each molecule, the first 10000 structures were taken from the rMD17 dataset. Of these 10000 structures, the first 8000 were used for training, the subsequent 1000 for cross-validation purposes and the final 1000 for testing. The 1000 test structures

drawn from the reference dataset will herein be referred to as the sub-sample test sets, to distinguish them from independently generated test sets.

Benchmarking with 8000 training structures demonstrates the exceptional performance of PairFE-Net in predicting Cartesian atomic forces and energies for each molecule in rMD17. The mean absolute errors (MAEs) for these sub-sample test sets are in the ranges of 0.020 – $0.221 \text{ kcal mol}^{-1} \text{ Å}^{-1}$ and $0.004 - 0.055 \text{ kcal mol}^{-1}$ for forces and energies, respectively (Table 1, Fig. 1a). 100% of energies and 99.6% of forces were predicted within 1 kcal mol⁻¹ (Å⁻¹) by PairFE-Net. Sorting the absolute errors and plotting them against their percentile produces a sigmoidal "S-curve", enabling visualisation of the distribution of errors.⁸⁴ The S-curves for aspirin demonstrate the very high fidelity of PairFENet-trained MLPs with respect to the underlying DFT dataset (Fig. 1c), as well as highlighting the long tail in the force and energy error distributions. Due to these long tails, it is perhaps more instructive to evaluate MLPs using maximum forces and energy errors, which are particularly important when it comes to simulation stability and prediction of physically relevant structures using trained MLPs: maximum absolute errors are typically an order of magnitude larger than the MAEs (Table 2). The maximum force error across all 468,000 Cartesian force components in the 10 molecule test sets was 3.729 kcal mol⁻¹ Å⁻¹ and the maximum energy error across all 10000 test set structures (1000 per molecule) was 0.603 kcal mol⁻¹. In addition to this exceptional prediction accuracy, the PairFE-Net scheme is computationally efficient, with each MLP taking no longer than 1-2 days to train using a single GPU node. The learning curve for aspirin trained on the rMD17 dataset is shown in Fig. 1b.

To enable a like-for-like comparison with other MLP schemes the performance of PairFE-Net trained networks was also benchmarked using a smaller training set comprised of the first 1000 structures, which is the typical dataset size used for benchmarking with rMD17 in the literature. The accuracy of MLPs trained on the smaller dataset is diminished compared to the 8000 structure training set, with the force and energy MAEs increasing to 0.187 – 1.020 kcal mol⁻¹ Å⁻¹ and 0.023 – 0.327 kcal mol⁻¹, depending on the molecule (Fig. S1). For aspirin, separate MLPs were trained using the training splits suggested by Christensen and von Lilienfeld⁶³ but no significant differences were observed in resulting test set MAEs compared to training MLPs using the first 1000 structures.

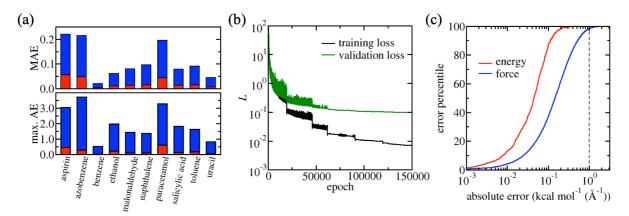


Fig. 1 Performance of PairFE-Net MLPs, trained on 8000 structures from the rMD17 benchmark dataset. (a) Force (kcal mol⁻¹ Å⁻¹) and energy (kcal mol⁻¹) test set prediction errors using the rMD17 dataset shown in blue and red, respectively. The top panel shows the mean, and the bottom panel shows the maximum, absolute errors for each 1000 structure test set. (b) Learning curves for the training (black) and validation (green) sets for aspirin. The loss function, L, which is minimised during training, is defined in Methods. Sudden drops in L correspond to a reduction in the learning rate. (c) Force (blue) and energy (red) S-curve plots of error distribution. The black dashed line marks the threshold for quantum chemical accuracy (1 kcal mol⁻¹).

3.2 Simulation stability and MLP generalisability.

Despite excellent test error performance, we observed that MD simulations for three molecules from rMD17 with conformational flexibility, aspirin, paracetamol and salicylic acid, frequently display simulation instabilities and unphysical structures when using MLPs trained on the rMD17 benchmark dataset. This can be attributed to incomplete conformational sampling in the reference dataset, i.e. the absence of structures corresponding to certain conformers and an inability of the trained MLP to predict forces for these conformers. Extrapolation to these unsampled regions leads to the generation of unphysical structures in subsequent timesteps of the simulation. Force predictions are poor for these newly generated unphysical structures, also not present in the training set, which results in further extrapolation and ultimately failure of the simulation trajectory either immediately or after a period of stable simulation. This issue undermines the ability to use the MLP to calculate any properties of interest from the MD simulation.

Table 1. Force and energy mean absolute errors of flexible drug molecule MLPs trained on the rMD17, MD-300K, MD-500K and Meta-300K datasets. After training, each MLP was separately tested on three independent test sets: the reference dataset sub-sample of 1000 structures, as well as torsion scans for rotation about ϕ and ψ .

molecule	test set	force (kcal mol ⁻¹ Å ⁻¹)			energy (kcal mol ⁻¹)				
		rMD17	MD-300K	MD-500K	Meta-300K	rMD17	MD-300K	MD-500K	Meta-300K
aspirin	sub-sample	0.22	0.15	0.31	0.27	0.06	0.03	0.10	0.16
	φ-scan	7.93	6.51	1.27	0.19	5.42	1.63	2.04	0.12
	ψ-scan	4.76	0.98	0.20	0.15	0.51	0.38	0.06	0.15
	sub-sample	0.20	0.19	0.29	0.31	0.04	0.03	0.06	0.06
paracetamol	φ-scan	15.61	21.50	22.75	0.12	3.60	5.74	8.00	0.02
	ψ-scan	3.78	0.06	0.08	0.09	0.50	0.02	0.02	0.02
salicylic acid	sub-sample	0.08	0.06	0.14	0.13	0.01	0.01	0.03	0.03
	φ-scan	136.60	20.28	0.19	0.25	175.24	11.52	0.05	0.16
	ψ-scan	5.09	2.99	0.74	0.20	0.83	1.89	1.37	0.09

Table 2. Force and energy maximum absolute errors of flexible drug molecule MLPs trained on the rMD17, MD-300K, MD-500K and Meta-300K datasets. Each trained MLP was separately tested on three independent test sets: the reference dataset sub-sample of 1000 structures, as well as torsion scans for rotation about ϕ and ψ .

molecule	test set	force (kcal mol ⁻¹ Å ⁻¹)			energy (kcal mol ⁻¹)				
		rMD17	MD-300K	MD-500K	Meta-300K	rMD17	MD-300K	MD-500K	Meta-300K
aspirin	sub-sample	3.04	4.49	9.85	30.94	0.46	0.31	0.88	2.81
	φ-scan	91.18	104.77	15.85	1.11	13.97	9.60	6.06	0.22
	ψ-scan	32.0	7.59	2.62	0.76	3.03	1.43	0.15	0.23
	sub-sample	3.17	2.58	4.11	3.50	0.42	0.16	0.28	0.43
paracetamol	φ-scan	172.50	170.41	137.27	0.92	17.43	31.60	44.59	0.10
	ψ-scan	29.81	0.44	0.55	0.63	4.10	0.05	0.04	0.08
	sub-sample	1.82	1.09	6.20	11.97	0.12	0.06	0.50	1.40
salicylic acid	φ-scan	3763.97	484.94	2.82	1.90	849.23	44.44	0.37	0.55
	ψ-scan	33.28	52.39	10.23	1.42	6.30	8.23	6.20	0.18

The principal torsional degrees of freedom defining the conformational landscapes of aspirin, paracetamol and salicylic acid are highlighted in Fig. 2a, 3a and 4a. The landscapes are dominated by rotation around the φ torsion angle, for which *trans* and *cis* conformers are defined as $\varphi = 180^\circ$ and $\varphi = 0^\circ$, respectively. For each molecule, the trans conformer is well sampled in the rMD17 dataset, the *cis* conformer is entirely absent (Fig. 2b, 3b and 4b). As a result, forces and energies for the *cis* conformer are poorly predicted and MD simulations initiated from these conformers become unstable within just a few steps. In addition, for salicylic acid, sampling with respect to ψ in the rMD17 dataset is also confined to a narrow region of $-65^\circ < \psi < 80^\circ$ (Fig. 4b). Indeed, in the entire rMD17 dataset of 10^5 structures, only 23.9%, 24.1% and 10.4% of the (φ, ψ) –surface of aspirin, paracetamol and salicylic acid is sampled, respectively (Table 3). Force and energy test MAEs are commonly treated as a proxy for the quality of a trained MLP. However, the failure of simulations using rMD17-trained MLPs despite their excellent test MAEs, demonstrates that MAEs do not necessarily provide a reliable approximation of the true generalisation error for the entire potential energy surface.

Table 3. Proportion (%) of (ϕ, ψ) —surface sampled in the rMD17 dataset and the new datasets generated in this work.

dataset	aspirin	paracetamol	salicylic acid
rMD17	23.9	24.1	10.4
MD-300K	16.0	18.9	7.6
MD-500K	26.8	24.5	29.4
Meta-300K	100.0	100.0	99.9

In addition to testing the aspirin, paracetamol and salicylic acid MLPs using these subsample test sets, they were separately tested on independently generated datasets representing full torsional scans around φ and ψ . Specifically, for each torsion angle, new datasets containing 72 structures were obtained from adiabatic scans via the density functional theory (DFT) at the B3LYP/6-31G* level. The computed DFT torsional profiles predict that, in the gas phase, the *trans* conformers have lower potential energies than the *cis* conformers by only 4.4, 2.0 and 3.9 kcal mol⁻¹ for aspirin, paracetamol and salicylic acid, respectively, with transition barriers of 9.7, 13.5 and 16.1 kcal mol⁻¹. Although these activation energies are unlikely to be crossed in a gas phase MD simulation of an isolated molecule at 300 K, it is probable that effects due to entropy, solvation or intermolecular interactions with other

molecules or surfaces will lower energy barriers or alter the conformational preference of the molecule. In any case, it is clearly an undesirable trait for a simulation initiated from any state other than the global energy minimum to be unstable. It is therefore critically important that trained MLPs can yield low force and energy test errors for the full torsion scan test sets as well as the sub-sample test sets.

Although rMD17-trained MLPs were able to reproduce the DFT torsion scans in wellsampled regions of conformation space, they were unable to reproduce it in poorly sampled regions. For the φ-scan, they do not predict physically sensible energies or forces for the cis conformers or near the energy barrier (red lines, Fig. 2f, 3f and 4f). Using these test sets, the inability to predict forces and energies around certain conformational minima and maxima results in very large MAEs, in some cases exceeding 100 kcal mol⁻¹ (Å⁻¹). For aspirin and paracetamol, rMD17-trained MLPs were able to accurately reproduce the ψ-scan reasonably well (red lines, Fig. 2g and 3g). However, for salicylic acid, rMD17-trained MLPs were also not able to reproduce the ψ-scan (red line, Fig. 4g). The observation that the MAEs for torsion scan test sets are much higher than for the sub-sample test sets, suggests that the latter may be a simple consequence of limited conformational sampling and significant structural redundancy in the rMD17 reference dataset. This observation also suggests that MAEs, when evaluated over a narrow conformational distribution such as in the sub-sample test sets, are a poor proxy for MLP quality and a weak predictor of stability upon deployment in MD simulations. The described failures of MLPs trained on the rMD17 dataset provide the dual motivations of the present study, namely to (i) enhance current understanding of the relationship between conformational sampling in the reference dataset and the stability MD simulations using MLPs, and to (ii) develop a systematic new methodology for training accurate and stable MLPs for flexible molecules.

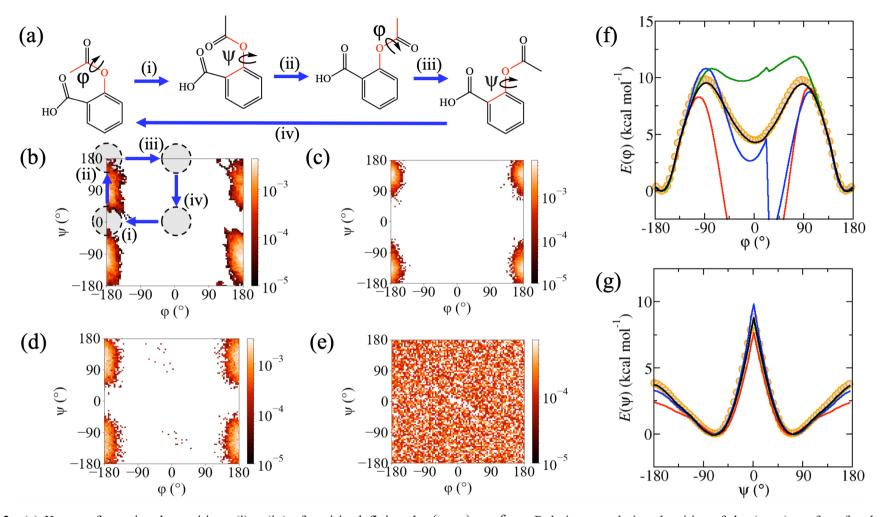


Fig. 2 (a) Key conformational transitions (i) – (iv) of aspirin defining the (φ, ψ)–surface. Relative population densities of the (φ, ψ)–surface for the (b) rMD17, (c) MD-300K, (d) MD-500K and (e) Meta-300K datasets with conformational transitions (i) – (iv) highlighted in (b). Energy profiles with respect to torsional rotation around (f) φ, and (g) ψ, obtained from B3LYP/6-31G* calculations (orange circles) and predicted from MLPs trained on rMD17 (red lines), MD-300K (blue lines), MD-500K (green lines) and Meta-300K (black lines) datasets.

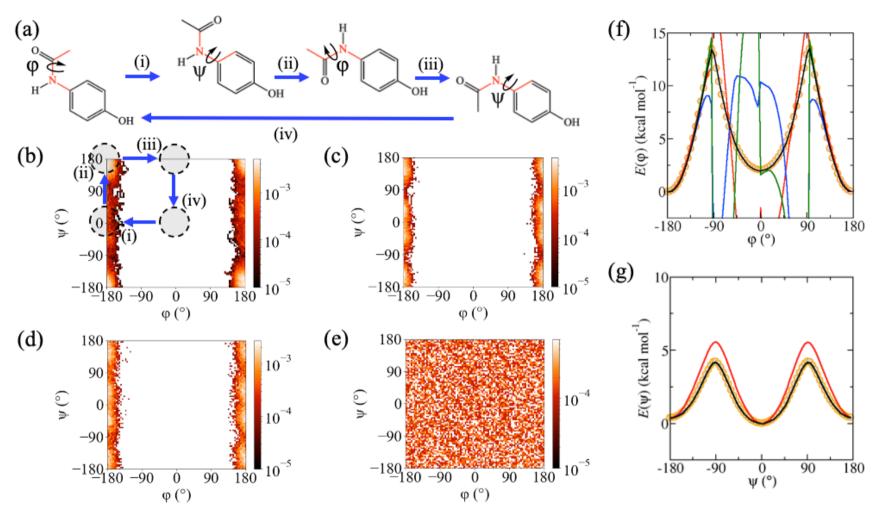


Fig. 3 (a) Key conformational transitions (i) – (iv) of paracetamol defining the (φ, ψ) –surface. Relative population densities of the (φ, ψ) –surface for the (b) rMD17, (c) MD-300K, (d) MD-500K and (e) Meta-300K datasets with conformational transitions (i) – (iv) highlighted in (b). Energy profiles with respect to torsional rotation around (f) φ , and (g) ψ , obtained from B3LYP/6-31G* calculations (orange circles) and predicted from MLPs trained on rMD17 (red lines), MD-300K (blue lines), MD-500K (green lines) and Meta-300K (black lines) datasets.

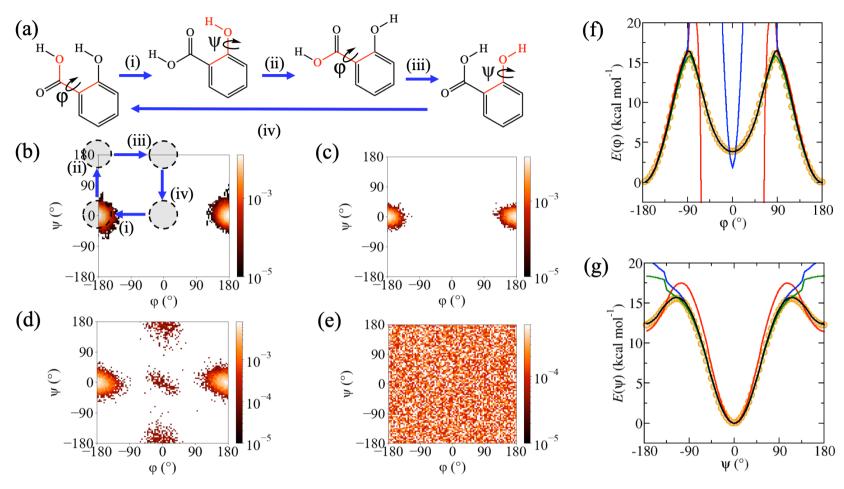


Fig. 4 (a) Key conformational transitions (i) – (iv) of salicylic acid defining the (φ, ψ)–surface. Relative population densities of the (φ, ψ)–surface for the (b) rMD17, (c) MD-300K, (d) MD-500K and (e) Meta-300K datasets with conformational transitions (i) – (iv) highlighted in (b). Energy profiles with respect to torsional rotation around (f) φ, and (g) ψ, obtained from B3LYP/6-31G* calculations (orange circles) and predicted from MLPs trained on rMD17 (red lines), MD-300K (blue lines), MD-500K (green lines) and Meta-300K (black lines) datasets.

3.3 Reference datasets from enhanced conformational sampling.

Due to the limitations of the rMD17 dataset, we pursued a two-step approach to generate new reference datasets for aspirin, paracetamol and salicylic acid. In the first step, structures were sampled from a 10 ns MD simulation using the GAFF empirical potential.⁷³ In the second step, structures sampled every 1 ps were evaluated via single point calculations at the B3LYP/6-31G* level,⁷⁸⁻⁸¹ in order to obtain reference forces and energies. This two-step approach has the advantage that the completeness of the dataset with respect to conformational sampling can be readily checked prior to more computationally expensive evaluation using a DFT functional in the second step. Sampling from longer MD trajectories than those used to generate the rMD17 dataset (100 – 200 ps) also reduces the time correlation of structures. For each molecule, three separate sampling strategies were employed: (i) MD simulation at 300 K, (ii) MD simulation at 500 K and (iii) metadynamics⁵⁸ simulation at 300 K. These datasets will be referred to as MD-300K, MD-500K and Meta-300K, respectively.

By sequentially adding Gaussian functions to a history-dependent biasing potential, metadynamics penalises previously well-sampled regions in collective variable space, enabling a system to escape local energy minima. This process can be exploited to enforce sampling of previously unexplored regions of the free energy surface, including barrier regions. In biasing the system away from previously visited structures with respect to the collective variables, metadynamics is ideally suited to the task of sampling structures for MLP reference datasets because it inherently reduces structural redundancy. Biasing potentials were used here to enhance sampling with respect to two collective variables, defined as the two torsion angles ϕ and ψ (Fig. 2a, 3a and 4a). Since the goal is to generate structures over the entire $\phi\psi$ surface rather than obtaining converged free energy surfaces, the original formulation of metadynamics was employed instead of the well-tempered variant. In completely saturating the conformational energy surface with structures in this manner, any problems pertaining to the underlying quality of the empirical potential (i.e. sampling the wrong underlying conformational distribution) are ameliorated.

Irrespective of the sampling methodology, all sub-sample test sets produce very low MAEs (<< 1 kcal mol⁻¹ (Å⁻¹) and maximum absolute errors (Tables 1 and 2). Meta-300K or MD-500K trained MLPs result in slightly larger force and energy MAEs than MLPs trained using the rMD17 or MD-300K datasets. At face value this suggests that the latter MLPs are of higher quality. However, it should be emphasized that the reported MAEs are associated only with structures in the region of conformation space that is sampled in their given test set. The MLPs can more accurately predict the energies and forces of structures that were sampled in

the same manner as the training set, but this metric is not necessarily representative of the generalisation error expected for the entire potential energy surface.

Similar to rMD17, the MD-300K trained MLPs fail to predict forces and energies in some regions of the torsion scans, particularly around the *cis* conformers (blue lines, Fig. 2f, 3f and 4f). This results in MAEs that are significantly higher than those obtained using the subsample test sets, in some cases exceeding 10 kcal mol⁻¹ (Å⁻¹). Indeed, the percentage of the (φ , ψ)–surface sampled in MD-300K is 16.0%, 18.9% and 7.6% for aspirin, paracetamol and salicylic acid, respectively (Table 3); this is even less than for rMD17 and can be attributed to a lower simulation temperature. Once again, simulations employing MD-300K trained MLPs initiated from the *cis* conformers fail to extrapolate reasonable forces, predict unphysical structures and destabilise within just a few timesteps. Although simulations initiated from the *trans* conformer can be stable for long periods, as demonstrated by successful completion of 25 ns MD simulations, when metadynamics is employed to promote a conformational change to the *cis* conformer, the trajectories become unstable and fail. In addition, simulations of salicylic acid initiated from the *trans* conformer with $\psi = 180^{\circ}$ are also unstable. These described pathological behaviours can once again be explained by the poor conformational sampling of the MD-300K datasets (Fig. 2c, 3c and 4c).

When MD-500K trained MLPs are evaluated using the torsion scan test sets, the force and energy MAEs are in general significantly improved compared to rMD17 and MD-300K (Table 1). This can be attributed to improved conformational sampling due to a higher temperature (compared to MD-300K) and longer simulation time (compared to rMD17). The coverage of the (φ, ψ) -surface in MD-500K increases to 26.8%, 24.5% and 29.4% for aspirin, paracetamol and salicylic acid, respectively (Table 3). For salicylic acid, force MAEs for the φ-scan test set are reduced from 20.28 to 0.19 kcal mol⁻¹ Å⁻¹ and energy MAEs from 11.50 to 0.05 kcal mol⁻¹ compared to MD-300K (Table 1). This improvement is because the higher temperature enables the system to escape the trans conformational minimum and numerous structures representing the cis conformer are generated in the reference dataset (Fig. 4d). Consequently, the φ-scan is well reproduced (green line, Fig. 4f) and stable MD simulations of this conformer can be performed. For aspirin, a small number of structures representative of the cis conformer are generated and, although the MD simulation does not immediately fail when initiated from this conformer, they become unstable after just a few picoseconds. In addition, although the energy profile around the cis conformer is improved relative to rMD17 and MD-300K, the relative energy of this conformer is still overestimated by ~6 kcal mol⁻¹ (green line, Fig. 2f). For paracetamol, the higher temperature used to sample structures in the

MD-500K dataset does not improve conformational sampling (Fig. 3d). Hence, like rMD17 and MD-300K, the φ-scan cannot be reproduced in the region of the *cis* conformer (green line, Fig. 3f) and simulations initiated from this conformer fail immediately. Overall, although increasing the simulation temperature from 300 K to 500 K does improve conformational sampling of flexible molecules in some instances, it does not guarantee complete conformational sampling or reliably produce simulation-ready MLPs.

However, for all three molecules, Meta-300K trained MLPs can predict with high accuracy the torsional energy profiles for both ϕ (black lines, Fig. 2f, 3f, 4f) and ψ (black lines, Fig. 2g, 3g, 4g). These MLPs yield very low force and energy MAEs (Table 1) on the torsion scan test sets, unlike those trained on rMD17, MD-300K and MD-500K. The excellent prediction accuracy of Meta-300K trained MLPs using independent datasets is due to comprehensive conformational representation in the reference dataset. The coverage of the $\phi\psi$ surface is now 100% due to the metadynamics-enhanced sampling (Table 3). When tested using the sub-sample test sets, Meta-300K trained MLPs generally perform less well in terms of MAEs (Table 1) and maximum errors (Table 2). These maximum errors can be attributed to a handful of unphysical structures in the training and test sets, generated due to instabilities in the empirical potential metadynamics simulations. The torsion profiles are correctly reproduced in spite of these unphysical structures in the training dataset and off-equilibrium structures such as these may even improve the robustness of the trained MLPs in MD simulations. ⁴⁶

Crucially, Meta-300K trained MLPs are sufficiently robust that 25 ns equilibrium MD simulations initiated from *any* conformer can be performed successfully without generating any unphysical structures; this includes commencing trajectories from cis ($\varphi = 0^{\circ}$) conformers, which were unstable using MLPs trained on the other datasets. In order to demonstrate the power of this new strategy, Meta-300K trained MLPs were employed in well-tempered metadynamics simulations to compute the conformational free energy surfaces with respect to the (φ, ψ) -surface, $F(\varphi, \psi)$, for each molecule (Fig. 5b, 6b and 7b). For purposes of comparison, $F(\varphi, \psi)$ was also calculated using the empirical GAFF potential. Obtaining reliable conformational free energy surfaces of flexible molecules necessitates simulations that are not only accurate, but also stable with respect to conformational change. To obtain converged free energy surfaces (Fig. 5b, 6b and 7b), 25 ns of stable simulation was required, far exceeding the timescales accessible to *ab initio* MD. Successful completion of these simulations is an excellent demonstration of the stability of the Meta-300K trained MLPs even in the regions around activation energy barriers, and the first examples of their kind using MLPs in the

literature. $F(\varphi, \psi)$ plots could not be calculated using the rMD17, MD-300K and MD-500K trained MLPs due to their instability for certain conformers and near energy barriers. As well as the equilibrium and two-dimensional metadynamics simulations, separate simulations were performed with only one collective variable, corresponding to rotation around φ or ψ . The one-dimensional free energy profiles for each molecule are shown in Supporting Information (Fig. S2). Meta-300K trained MLPs were stable in all simulations and no unphysical structures were observed. In total, these simulations amount to more than 100 ns of stable simulation time per molecule, sampling the entire conformational free energy surface of these small organic molecules *in vacuo*, including various conformational minima and free energy barriers.

Significant differences are revealed between $F(\varphi, \psi)$ computed via the empirical GAFF potential and the PairFE-Net based MLPs (Fig. 5, 6 and 7). For aspirin, although the empirical potential predicts the relative energies of the trans ($\varphi = 180^{\circ}$) and cis ($\varphi = 0^{\circ}$) conformers correctly, the free energy barrier for the trans-cis conformational transformation is underestimated (Fig. 5). For paracetamol, the MLP simulation shows that the free energy of the cis ($\phi = 0^{\circ}$) conformer is lower by 1.1 kcal mol⁻¹ (Fig. 6b). This is in contrast to the potential energy, for which the trans ($\varphi = 180^{\circ}$) conformer is 2.0 kcal mol⁻¹ lower than the corresponding cis conformer (Fig. 3f) in agreement with the literature. 85 This contrast suggests that entropy plays a key role in determining the conformational preferences of paracetamol. This important entropic contribution arises from the significant interdependency of the two torsional motions in paracetamol. Specifically, in the *trans* conformer, rotation with respect to ψ is limited where the heavy atoms are coplanar. On the other hand, in the cis conformer, ψ -rotation is relatively unhindered (Fig. S3). This example clearly demonstrates the importance of performing MLPbased MD simulations that combine (i) quantum chemical accuracy, to distinguish between energetically similar conformers; and (ii) sufficient stability to perform the long simulations required to sample adequately the torsional space of the flexible molecule, enabling free energies to be accurately computed. In contrast to the MLP, the empirical potential predicts that the free energy of the *trans* conformer of paracetamol is lower than *cis* by 1.2 kcal mol⁻¹. For salicylic acid, the free energy profile $F(\varphi, \psi)$ predicted by the empirical potential is qualitatively different from the MLP-derived surface (Fig. 7), incorrectly predicting a significant conformational minimum at $\psi = 180^{\circ}$, more stable than $\psi = 0^{\circ}$, which is not seen in the MLP free energy surface. The trans ($\varphi = 180^{\circ}$) conformer is stabilised by a hydrogen bond between the hydroxyl and carboxylic acid functional groups via the carbonyl oxygen.

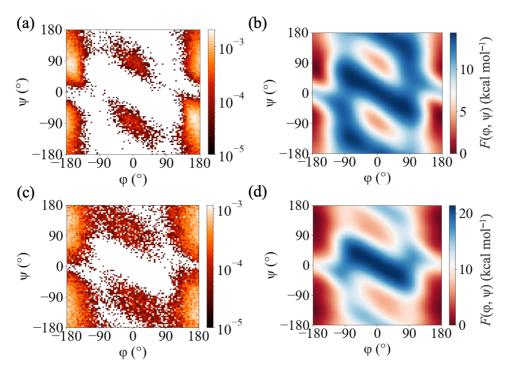


Fig. 5 For aspirin, (a) relative population densities for the (ϕ, ψ) -surface from the well-tempered metadynamics simulation and (b) conformational free energy surface calculated using the Meta-300K trained MLP and (c) relative population density map and (d) conformational free energy surface calculated using GAFF.

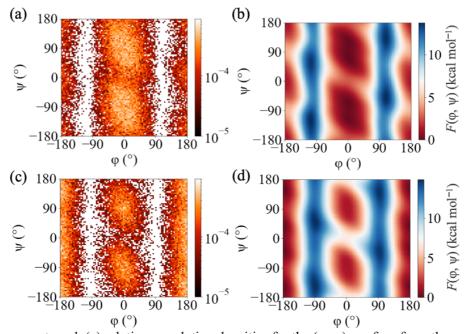


Fig. 6 For paracetamol, (a) relative population densities for the (ϕ, ψ) -surface from the well-tempered metadynamics simulation and (b) conformational free energy surface calculated using the Meta-300K trained MLP and (c) relative population density map and (d) conformational free energy surface calculated using GAFF.

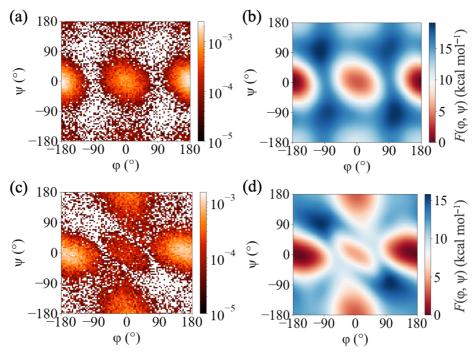


Fig. 7 For salicylic acid, (a) relative population densities for the (ϕ, ψ) -surface the well-tempered metadynamics simulation and (b) conformational free energy surface calculated using the Meta-300K trained MLP and (c) relative population density map and (d) conformational free energy surface calculated using GAFF.

Structures in the Meta-300K dataset were sampled using a biasing potential constructed from Gaussian functions with heights, h, of 0.24 kcal mol⁻¹. To further illustrate the importance of adequate conformational sampling in training datasets on simulation stability, additional datasets for aspirin were prepared by sampling structures using metadynamics biasing potentials with h = 0.06, 0.12 and 0.18 kcal mol⁻¹. These datasets have (φ, ψ) -surface coverages of 97.8%, 99.6% and 99.9%, respectively. For h = 0.06 kcal mol⁻¹ although the coverage using metadynamics simulations is still vastly improved compared to the rMD17, MD-300K or MD-500K datasets, a well-tempered metadynamics simulation using trained MLPs became unstable after less than 1 ns and as a result the conformational energy surface could not be reliably calculated. This instability arises due to one small poorly sampled region, accounting for just 2.2% of the (φ, ψ) -surface (Fig. S4). The other additional datasets (with h = 0.12 and 0.18 kcal mol⁻¹) did not generate unphysical structures over 25 ns, demonstrating that coverage needs to be approximately 100% in torsional space to guarantee stability. Although simulation stability appears to be strongly dependent on the biasing potential used to generate the dataset, the force and energy MAEs appear insensitive to h (Table S1), once again demonstrating that low subsample test set errors alone are insufficient to demonstrate the quality of an MLP with respect to stability.

4. Conclusions

We have shown that, when the sub-sample test set has a narrow conformational distribution, MAEs can present a misleading representation of the true generalisation error for flexible molecule MLPs. When conformational sampling is improved, prediction accuracy using the independently generated torsion scan test sets improves but prediction accuracy using the subsample test sets reduces. This observation highlights an important trade-off when training flexible molecule MLPs. If MLPs are trained with the sole objective of obtaining low subsample MAEs, this may hinder the ability to perform stable simulations for all conformers, because this objective is most easily achieved using datasets with limited conformational sampling and high structural redundancy. This was observed with the rMD17 and MD-300K datasets, which proved to be the least stable when used in MD simulations when assessed over the full conformational distribution. By contrast, MLPs trained on reference datasets with the most complete conformational sampling (Meta-300K) generally have larger sub-sample test set errors but are more stable in MD simulations. Our observations complement those of Fu et al.,45 who demonstrated that achieving low test set MAEs was not a sufficient condition for obtaining MLPs capable of producing stable trajectories or reproducing simulation-based metrics.

This study highlights the importance of training MLPs on datasets containing all relevant conformers, such as the Meta-300K datasets published in this work; or at least testing MLPs on independently prepared external datasets, because sub-samples test sets will inherit the same sampling deficiencies present in the training set. Finally, also we encourage the reporting of maximum errors when evaluating MLP performance, as these will ultimately determine the stability of simulations. This suggestion resonates with the recent work of Vita et al., ⁸⁶ who proposed evaluation of the loss landscape, to establish the difference in extrapolation abilities of MLPs with similar test errors.

In this work, we have demonstrated that an essential consideration for performing MD simulations of flexible molecules using MLPs is training set selection. Enhanced sampling is key to preparing reference datasets with adequate conformational representation. MLPs trained on Meta-300K datasets have the appealing twin characteristics of enabling stable long timescale MD simulations such that conformational free energy surfaces for example can be calculated; while also inheriting the quantum chemical accuracy of the reference level of theory. The computational effort to obtain these insights directly from numerically very intensive AIMD simulations would be orders of magnitude greater and intractable without

access to large-scale supercomputing resource. The fact that the Meta-300K trained MLPs can accurately reproduce the torsional energy scans and be used in long and stable MD simulations is a *post hoc* justification for our approach to generating reference datasets. It also suggests that a global scheme based on nuclear repulsion force input features can correctly resolve states along the transition paths defined by torsional rotations.

The use of a computationally inexpensive empirical potential was found to be a viable approach for sampling reference dataset structures and the quality of the trained MLP is somewhat independent of the level of theory used to sample conformation space. Even if the chosen empirical potential did not correctly sample the underlying equilibrium probability distribution, an approach based on metadynamics ensures that, in saturating the conformational free energy surface with deposited Gaussian functions, all conformers are well-represented in the reference datasets (Fig. 2e, 3e and 4e). For example, for salicylic acid, an empirical potential predicts a significant free energy minimum not present in the MLP free energy surface (Fig. 7), despite the MLP being trained on structures generated using the empirical potential. The differences between MLP and empirical potential free energy surfaces seen here, which could lead to the incorrect identification of the most stable conformers of drug molecules, can be attributed to much larger errors associated with the empirical potential: this point is demonstrated by evaluating the empirical potential force errors relative to the DFT reference. The force MAE for the GAFF potential from an equilibrium simulation of the aspirin trans conformer was 10.2 kcal mol⁻¹ Å⁻¹, increasing to 12.2 kcal mol⁻¹ Å⁻¹ in a 300 K metadynamics simulation; this is accompanied by a significant increase in the number of large force outliers (Fig. S5). In this respect, the Meta-300K trained MLPs result in at least an order of magnitude improvement in force predictions. The trade-off for this improvement in accuracy relative to the empirical potential is just a two-fold increase in total MD simulation time.

Our key findings have far-reaching implications for the future development of MLPs to compute structural, dynamical and thermodynamical properties of flexible molecules from molecular simulation. When accurate and stable small molecule MLPs are deployed in the condensed phase,^{64, 70} they have the potential to guide innovation in design across a wide variety of applications, from therapeutics and polymers to ionic liquids and environmental contaminants. Related future work should focus on strategies to remove structural redundancies from reference datasets and the use of generalisable collective variables (e.g. those based on the distance matrix⁸⁷), which may pave the way to a fully automated pipeline for generating robust reference datasets for flexible molecules.

Data Availability

A supporting information document contains supporting analysis for this paper. Newly generated reference datasets and trained models are provided and the datasets have also been uploaded to Figshare (https://doi.org/10.6084/m9.figshare.25211540).

Code Availability

All artificial neural network training and MD simulations were conducted within the PairNetOps package. A version of this package is provided in order to reproduce the key results in this paper.

Author Contributions

Christopher D. Williams: Writing, investigation, conceptualisation, methodology, software, data curation, formal analysis.

Jas Kalayan: Methodology, software.

Neil A. Burton: Supervision, conceptualisation, project administration, funding acquisition, writing – review and editing.

Richard A. Bryce: Supervision, conceptualisation, project administration, funding acquisition, writing – review and editing.

Competing interests

There are no conflicts to declare.

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