Dynamic, remote-controllable electroactive hydrogel waveguide architectures

Oscar Alejandro Herrera Cortes, Kathryn A. Benincasa, Kevin Vaughan, Derek. R. Morim, Natalie Blanchard, Tomas Omasta and Kalaichelvi Saravanamuttu\*

Affiliations :

Department of Chemistry and Chemical Biology, McMaster University, 1280 Main St. West, Hamilton, Ontario L8S 4M1, Canada

\*Correspondence to:

Prof. Kalaichelvi Saravanamuttu

Department of Chemistry and Chemical Biology

McMaster University

1280 Main St. West, L8S 4M1 Ontario, Canada

Email: kalai@mcmaster.ca

## 1 Abstract

2 Remote-controllable waveguide architectures inspired by living organisms with unique flexible, 3 light guiding properties were fabricated using self-trapped beams of incoherent light. Made of 4 electroactive hydrogels, light-guiding structures are generated through a nonlinear, self-inscription 5 process that utilizes visible beams from light-emitting diodes (LEDs). Due to irreversible refractive index 6 changes experienced by photoinduced chemical reactions, these self-trapped beams permanently 7 inscribe cylindrical waveguides along their paths. Taking advantage of this phenomena, we fabricate 8 macro-scale, remote controllable waveguide structures in the form of rectangular prisms and arrays of 9 cylindric waveguides. We also fabricate micro-scale structures for remote actuation, in the form of 10 rectangular prisms embedded with thousands of waveguide units. By applying and varying external 11 electric fields, we dynamically control the bending, angular orientation, and rotation (up to 360°) of these 12 pliant light-guiding structures. This allows precise, remote control of the waveguided light output.

#### 14 **Main**

15 Understanding the relationship between structure and function in biological systems and materials has fascinated researchers for generations.<sup>1</sup> The female angler fish, for example, has a 16 17 symbiotic relationship with a species of luminescent bacteria. This bacterium is housed within a lure that 18 dangles atop the female's head, where it has a safe environment to grow. The angler fish can direct the light output of this bioluminescent species to attract prey or potential mates, by flexing this lure.<sup>2,3</sup> Squid, 19 20 cuttlefish and octopi display remarkable camouflaging due to pigment filled organs, called 21 chromatophores, present in their skin. These organs expand upon contraction of muscle fibers when 22 triggered by an electrical stimulus. As a result, the chromatophores increase their size temporarily.<sup>4</sup> The 23 visual perception by an observer is a changeable pattern in the form of an array of coloured dots that 24 seem to migrate in space, distracting the eye. Inspiration by these types of hybrid systems can lead to the 25 fabrication of stimuli-responsive, flexible, light guiding materials with remote actuation for various 26 applications.

27 Most reported soft actuators reported in the literature comprise of multi-material light guiding 28 architectures. Reminiscent to the angler fish's ability to direct light, there are currently few examples of 29 waveguide actuators reported for remote light guidance. This includes the development of liguid crystal elastomer fibers patterned with gold nanoparticles,<sup>5</sup> shape memory polymer fibers embedded with 30 Al<sub>2</sub>O<sub>3</sub>/ZnO photonic shells of varying thickness and observable structural colour, <sup>6</sup> ferromagnetic soft-31 32 robots embedded with optical fibers, <sup>7</sup> and liquid crystal fiber arrays.<sup>8</sup> Reminiscent to the visual migration 33 of chromatophores seen in squid skin, and how light can be reflected towards and away from a visual 34 observer, thermoactivated liquid crystals elastomers are able to remotely actuate and act as an optical 35 sensor. Liquid crystal elastomers undergo a transition from an opaque domain to a transparent upon 36 heating. As a result, light coupled into the transparent elastomer can then be received by a detector 37 (leading to an optical signal). The bending state of the liquid crystal elastomer translates to the amount of light monitored and collected at the detector.9 38

39 3-D printing techniques that offer a fast route to create mechanically stable objects have emerged
40 in recent years. While some of them rely on sequentially depositing layers similar to the methods used in

traditional additive manufacturing techniques, there also exist methods capable of producing seamless objects.<sup>10–13</sup> Techniques that create objects employing 3D-unit operations have also emerged, such as volumetric methods that rely on tomographic reconstruction<sup>14,15</sup> and those that rely on the propagation of nonlinear optochemical waves known as computed axial lithography (CAL).<sup>16</sup> Objects are created by launching 2D patterns calculated using the same concepts from computed tomography imaging. The combined energy dose from the 2D images solidifies the resin creating the desired object.

47 Our group has demonstrated that 3D objects can be created in a single step when patterned blue nonlinear beams propagate through a photopolymerizable medium.<sup>17</sup> In traditional stereolithographic 48 49 methods, the thickness of the resulting film is limited to  $< 100 \,\mu$ m, due to the divergence experienced by the incident beam blurring as it travels through the medium, decreasing the overall resolution. This beam 50 51 divergence can be controlled in nonlinear media, such as a photopolymer material, as the divergence of 52 the beam is counteracted by the photo-induced changes in refractive index originating from the 53 photopolymerization reactions. This results in the creation of 2D objects with depths > 100 µm.<sup>18</sup> 54 Additionally, these objects are homogenous along their direction of propagation since they are obtained in 55 a continuous matter without layering artifacts. Therefore, each element obtained through this method is a 56 polymer waveguide. This method can be extended to long-area waveguide lattices with large populations 57 of filaments embedded in a shape, as imposed by a 2D amplitude mask.<sup>19-22</sup> This lattice would be 58 impossible to fabricate with traditional photolithographic methods, including current volumetric 3-D printing techniques discussed earlier. Polymer waveguides have potential applications in light-base 59 60 therapies including photo-activated drug release and the migration of biological cells within a hydrogel (via waveguided light). 23,24 61

Herein, we describe the fabrication and demonstration of different geometries of flexible, light guiding structures as inspired by deep sea creatures. These geometries include a planar rectangular prism; individual and small arrays of cylindrical waveguides; and a rectangular prism embedded with longrange micro-scale waveguide filaments. Using a stimuli-responsive electroactive hydrogel, we have demonstrated that the fabricated waveguides motion and orientation can be precisely and remotely controlled through external electrical fields. Reminiscent of the bioluminescent lure directed by the female angler fish to attract prey, we demonstrate how a flexible electroactive waveguide's light output can be 69 precisely controlled due to the degree of bending experienced by the waveguide, and how the light output 70 is successfully directed to and from a detector. Reminiscent of chromatophores found in various 71 cephalopods and how the arrays of pigment migrate in space to distract predators, we demonstrate how 72 micro-lattices of waveguides display a similar behaviour, guiding their light output as a bright array of dots 73 in space to a detector.

74 Through the vectorial sum of external electric fields, we can dynamically control the bending, angular orientation, and rotation (up to 360°) of these flexible, light-guiding structures. This allows for 75 76 precise, remote control of the waveguided light output without the use of computers or other electronic 77 devices. Flexible fiber-like structures that decrease their brightness while bending can act as sensors in 78 prosthetic hands or artificial skins.<sup>25,26</sup> Additionally, fabricating stimuli-responsive waveguides or actuators 79 with waveguide-circuitry embedded can help to develop highly sophisticated soft robots that can detect 80 their own shape or deformation state.<sup>27,28</sup> They could also become part of devices that can be used for minimally invasive surgery or light-based therapies.<sup>29</sup> 81

## 83 Unidirectional actuation of hydrogel waveguides

84 An optical waveguide is a dielectric structure that can guide electromagnetic waves. Guided waves are dependent on the phenomenon of total internal reflection (TIR), which confines light to travel 85 along the optical axis of a waveguide. The core region of a waveguide is assumed to have a refractive 86 index (n<sub>1</sub>) higher than the surrounding material (n<sub>2</sub>) to achieve TIR (Figure 1a, left).<sup>30</sup> The material 87 88 presented herein is a flexible, electroactive hydrogel. This means that as an electric field is applied the 89 flexible waveguide will bend in the direction of the cathode, continuing to guide light along its optical axis 90 (Figure 1a, center). This light will be guided to a detector (i.e., CCD camera) and will continue to be 91 collected so long as the waveguided light is within the detectors field of view, or light collection range. If 92 the polarity of the electric field is inverted, the waveguide will bend the opposite way continuing to guide 93 light as it moves (Figure 1a, right). The bending of the hydrogel is caused by the flow of ions in solution 94 as they migrate through the porous hydrogel to the cathode. This is a direct result of the material used to 95 fabricate the waveguide structures.

96 Hydrogel resin used to fabricate planar waveguides was developed using two previously described formulations.<sup>31,32</sup> The precursor, denoted as H1 (see Methods and Materials 5.1 for resin 97 formulation), consists of acrylic acid, N,N'-methylenebis(acrylamide) crosslinker, Irgacure ® 819 ( $\lambda_{max}$  = 98 99 468 nm) photoinitiator and sodium phosphate buffer solution (pH = 7.4). Although there is buffer present 100 in the resulting hydrogel, only  $\approx 2$  % of monomers are deprotonated after fabrication. When the hydrogel 101 is immersed in buffer solution, most acrylic acid monomers lose their acidic proton producing an anionic 102 network (Figure 1b, left). When an electric field is applied, the counterions within the hydrogel bound to a 103 carboxylate group remain within the network to preserve the electroneutrality inside the gel. Counterions 104 that are not bound migrate freely in the solution and within the gel towards the cathode. The migration of 105 negative ions within the gel is hindered since the network is negatively charged. As a result, there is a 106 higher concentration of ions on the cathode side than in the anode side (Figure 1b, center). This 107 concentration gradient generates a higher osmotic pressure on the anode side ( $\pi_1$ ; swelling) than on the cathode side ( $\pi_2$ ; shrinkage).<sup>33,34</sup> The pressure difference,  $\pi_1 > \pi_2$ , causes bending of the gel (**Figure 1b**, 108 109 right) toward the cathode.



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Figure 1 | Electroactive waveguide working principles. a) Scheme depicting a rectangular waveguide 111 (refractive index, n1) guiding an optical mode to the entrance face of the waveguide (represented by the 112 yellow circle) through TIR. In the presence of an electric field, the waveguide bends moving to a new 113 position while still guiding light. b) Scheme of bending mechanism in electroactive hydrogels. Changes in 114 115 osmotic pressure ( $\pi_1 > \pi_2$ ) result in bending. c) Scheme depicting the self-inscription process of rectangular waveguides obtained by launching a blue LED beam through a rectangular shape mask. The 116 resultant waveguide generated through photopolymerization of H1 precursor is shown on the right. (scale 117 118 bar = 2 mm) 119 Rectangular hydrogel waveguides were obtained by launching a blue ( $\lambda = 445$  nm) LED beam

- 120 through a rectangular shape mask (Figure 1c, left) pasted to a glass coverslip. Free-radical
- 121 polymerization of acrylic acid monomers were initiated launching these patterned beams to cells
- 122 containing the **H1** hydrogel precursor. The refractive index change initiated by the photopolymerization
- 123 reaction counteracts the natural diffraction of light, allowing the patterned beam to travel without diverging

(Figure 1c, middle) while inscribing a polymeric structure (Figure 1c, right). These rectangular prisms
 remained on the coverslip (with shape mask removed), acting as the entrance face of the cell in which
 they were fabricated.

127 To investigate the remote actuation of these hydrogel structures, a rectangular waveguide 128 adhered to a glass coverslip was placed between two carbon electrodes (placed 7 mm apart) and partially 129 immersed in a phosphate buffer solution (pH = 7.4). A blue ( $\lambda$  = 445 nm) LED beam was collimated and 130 launched through the part of the waveguide that was not immersed in the buffer solution. The beam 131 output of the waveguide was monitored with a CCD camera while the position of the waveguide was 132 tracked using a cellphone camera (Figure 2a). With an electric field applied, the waveguide bends 133 towards the cathode and the beam output shifts in the same direction. This demonstrates how light 134 confined within the prism can be spatially manipulated. (Figure 2b).

135 The polarity of the electric field was inverted when the monitored intensity output for the planar 136 waveguide reached a steady, nonchanging value (roughly 60 s). When the polarity of the electric field 137 was inverted, the bending direction reverses as a result (Figure 2b and video S1). The maximum 138 bending angle for rectangular waveguides 10 mm long is calculated to be -12° (see Supplementary 139 Information S1 and Figure S1 for the calculation of bending angle). This was achieved at 410 s, after 140 which the electric field had been inverted twice (Figure 2c). As expected, there is a noticeable drop in the 141 intensity when bending angles begin to approach a local maximum. This is due to confined light within the 142 waveguide being directed away from the detector (i.e., CCD camera). The net effect every time the 143 waveguide reaches the maximum bending angle, or similarly large angle, is a bright sheet (with  $\approx 2$  mm 144 thickness) that seems to turn "off" from the point of the observer or detector. Quantitatively, this directly 145 translates to the intensity output of the waveguide decreasing by  $\approx 90\%$  on average (Figure 2d). The 146 intensity of the beam is recovered when the waveguide is perpendicular to the detector or reaches a 147 bending angle of  $\theta \approx 0^\circ$ . For the waveguides 6 and 9 mm long the maximum bending angle is  $\approx -2^\circ$  and  $\approx$ -11°, respectively. When these bending angles are reached, the intensity of the beam output decreases to 148  $\approx$  22 and  $\approx$  75% of its original value, respectively. Graphs showing the time evolution of bending angle 149 150 and beam output intensity of these waveguides can be found in Figure S2.

In the experiment shown in **Figure 2c**, the rectangular waveguide is able to reach maximum bending angles faster after the electric field has been applied in two actuations (i.e., after 300 s), than compared to the beginning of the experiment. This is most likely caused by an increase of fixed charge density in the network (mole fraction of charged species) as the hydrogel swells and more carboxylic acid groups get deprotonated. It has been previously shown, in hydrogels with acrylic acid as the monomeric unit, that the rate of bending increases with the fixed charge density if the total charge density is below 70%.<sup>32,35</sup>



Figure 2 | Remote actuation of a planar, hydrogel waveguide. a) Optical setup used for the electroactuation of rectangular hydrogel waveguides. Light from a blue LED was collimated (L1, f.l. = 250 mm) and launched to the entrance face of the waveguide. Imaging of the optical profile at the exit face of the waveguide was achieved using a pair of a planoconvex lenses (L2, L3, f.l. = 250 mm) and a highresolution charge-coupled device (CCD) camera. A phone camera captured the top-view actuation of the gel. b) Top-view waveguide actuation (scale bar = 4 mm) with complimentary 2D beam intensity profiles (scale bar = 1 mm) of a waveguide 10 mm long under the effect of an electric field (1300 V/m) when

irradiated with a blue LED (see **video S1**). **c)** Time evolution of the bending angle as the electric field is applied. The maximum bending angle (-12°) is indicated by a red line on the graph. Electric field polarity inversions are indicated above the graph. **d)** Normalized intensity of the beam output of the waveguide as the electric field is applied. As the waveguide bends away from the optical set up, the intensity on the detector decreases  $\approx$  90% when the maximum bending angle is reached (as indicated by the red line).

171 The method of hydrogel fabrication is widely versatile to different shapes and architectures of light 172 guiding geometries. To demonstrate this versatility, cylindrical hydrogel waveguides were fabricated using 173 a blue LED ( $\lambda$  = 445 nm) launched through a circular shape mask (Figure 3a) placed at the entrance face 174 of a cell containing hydrogel precursor H1. The fabricated cylindrical waveguide was kept vertical for the 175 following unidirectional remote actuation investigations (Figure 3b). A scheme of the setup used for these 176 experiments is shown in Figure S3, which consists of a power supply connected to rectangular carbon 177 electrodes inserted on the sides of a custom 3D printed cage. In the middle of the cage is a hole in which 178 a 10 mm cylindrical hydrogel waveguide can be placed. Sodium phosphate buffer (pH = 7.4) solution 179 covered around 8 mm ( $\approx$  80%) of the cylindrical waveguide, with the remaining top portion exposed to air. 180 While an electric field was applied to actuate the cylinder, an attenuated beam from a compact laser 181 module ( $\lambda$  = 520 nm) was launched through the bottom. The spatial evolution of the cylinder and the 182 intensity of its beam output was recorded using a Moticam 3+ CMOS camera (Figure 3c).

183 To observe the unidirectional motion of the hydrogel cylinder, the polarity of the applied electric field was inverted every 180 s. A maximum bending angle of -12° was achieved at 180 s, at the end of the 184 185 first actuation. After applying the electric field for roughly 400 s (i.e., two electric field actuations), the 186 intensity of the outputted beam starts to decrease over time. This is a result of three factors: fatigue 187 damage; the fracture plane created near the base of the hydrogel when the cylinder bends and/or swells; 188 and droplets of water containing carbon particles (from the carbon electrodes) built up on the waveguide 189 tip over time, consequently blocking the beam output. As a result of these factors the intensity we can 190 visually see the beam fully recover at  $\approx$  260 s, as the waveguide points directly at the camera again (i.e., 191  $\theta \approx 0^{\circ}$ ) (see **video S2** for the entire recorded experiment).

Unidirectional actuation of square arrays of four (2 x 2) and nine (3 x 3) cylindric waveguides were also demonstrated. To actuate cylindrical arrays of waveguides vertically, a similar setup was used as described in **Figure S3a**. The only difference presented here is the presence of additional holes in the base of the cage to accommodate for the additional cylindrical waveguides (**Figure 3d**). Arrays of 196 waveguides were illuminated and actuated in a unidirectional fashion. A general trend can be observed 197 for both rectangular and cylindric waveguide unidirectional remote actuation, with respect to light 198 guidance and control. When the electric field is applied, the waveguide in question bends towards the 199 cathode and the intensity of the beam output decreases as the light output is guided away from the 190 detector. This trend is also observed when the polarity is reversed and the waveguide bends in the 191 opposite direction. The light intensity can be fully regained once the waveguide is aligned again with the 192 detector, thus the light output is guided towards it and is recovered.



204 Figure 3 | Vertical remote actuation of cylindrical hydrogel waveguide. a) Fabrication scheme for the cylindrical hydrogel waveguides. Light from a blue LED is launched through a circular shape mask, to the 205 H1 precursor. b) Image of the cylindric waveguide when illuminated with a green laser beam (scale bar = 206 207 5 mm). c) Time evolution images of a cylindric waveguide under the effect of an electric field (520 V/m) 208 (see video S2). Electric field polarity switches are as indicated. d) Scheme of cages used to house cylindric waveguide arrays, highlighting the distance away from the cathode each cylinder was based on 209 the cage and array implemented. Cylindric waveguides are irradiated with an expanded green beam ( $\lambda$  = 210 211 520 nm) and actuated with an electric field (520 V/cm) (see video S3).

### 213 Multidirectional actuation of arrays of cylindric waveguides

214 Most actuation of electroactive hydrogel structures involves unidirectional bending and locomotion with a single pair of electrodes.<sup>32,35</sup> Upward propulsion can also be achieved with a single pair 215 216 of electrodes by applying alternating electric fields, which was demonstrated by Zhao et al.<sup>36</sup> Here the 217 authors designed an electroactive hydrogel, which had four "tentacles", inspired by jellyfish movements. 218 Bending the tentacles upwards using a voltage of 7 V followed by a downward realignment by applying a 219 higher voltage (15 V) in the opposite direction produced an upward propulsion. Using arrays of four 220 electrodes has been proven to be an efficient way to achieve complex motions in electroactive hydrogel 221 structures. This approach has also been applied to ionic polymer metal composites by placing electrodes on orthogonal positions of the polymer surface,<sup>37,38</sup> where bending is limited to 90° and 45°. 222

223 We have demonstrated arrays of cylindric waveguides can be unidirectionally actuated with the 224 application of a single electric field. By combining two electric fields through vectorial analysis (Figure S3b and Figure S3c for vectoral analysis used for the following investigations), arrays of cylindric 225 226 waveguides can be remotely bent, precisely, in multiple directions. To accomplish this, the combined 227 electric field was kept constant (520 V/m) to achieve targets angles of: 0°, 60°, 120°, 180°, 240°, 300° and 228 360°. Bending angles for the following experiments were calculated only for the individual cylindric 229 waveguide placed in the center of the square array (Figure 4a). We used the first bending direction of the 230 cylinder as a reference point (i.e.,  $\theta = 0^{\circ}$ ). For the initial bending direction, the electric field was applied for 231 240 s, after which it was applied for 180 s for each direction after. This process was repeated until 232 cylindrical waveguides completed a full cycle (see video S4 for single and arrays of waveguides rotating). Experimental bending angles for the individual waveguide were calculated to be: 0°, 46°, 84°, 173°, 263°, 233 234 302° and 338° (as indicated in Figure 4a). While the calculated experimental angles show some deviation 235 from the target angles, this approach allows precise control over the waveguide's direction, thereby 236 allowing for multidirectional rotation of the waveguide. The intensity of the waveguide output remains 237 stead as it bends, initially. Over time the intensity output begins to decrease, which can be attributed to damage experienced by the hydrogel. The same combined electric field was also applied for each 238 239 configuration of cylindric waveguide arrays;  $(2 \times 2)$  (Figure 4b) and  $(3 \times 3)$  (Figure 4c).



Figure 4 | Multidirectional bending of cylindric waveguide arrays. a) Configurations used to direct the cylindric waveguide. Target angles are displayed along with the electric field vectors used to orient the waveguide. Experimental images of each output displayed below, with calculated bending angles (red) displayed alongside each image. b) 2x2 and c) 3x3 cylindrical waveguide array experimental output. Target angles (white) displayed alongside images.

### 247 Unidirectional remote actuation of a micro-waveguide array

We have previously shown that it is possible to fabricate micro-waveguide arrays in an organosiloxane photopolymer.<sup>39</sup> We hypothesize that creating these arrays using an electroactive hydrogel would allow the formation of a structure with a bright array of dots that would seem to migrate in space from the point of view of an observer, reminiscent of the chromatophores seen in various types of cephalopods which present as dots migrating in space. Like the behaviour observed in the planar and cylindric hydrogel waveguides, light traveling through these waveguide micro-channels will be confined and guided towards the direction in which the waveguide is bending.

255 To explore this, waveguide lattices were inscribed within hydrogel prisms (L = 6 mm, W = 2.5256 mm). The hydrogel precursor used to fabricate waveguide lattices was similar to precursor H1, with the 257 added use of camphorquinone (CQ) co-initiator. This adapted precursor is hence denoted as H2. The 258 inclusion of CQ as a co-initiator in the photoinitiated free radical reaction will produce a lower 259 concentration of free radicals, thereby resulting in a smaller degree of polymerization.<sup>43</sup> This will produce 260 the refractive index contrast needed for the formation of prisms with waveguides embedded. To fabricate 261 hydrogel micro-lattices, the process to create rectangular waveguides was repeated, with the addition of a 262 periodic amplitude mask ( $\Lambda$  = 80 µm) used to order the emerging periodic microstructure (**Figure 5a**). 263 When examined through transmission optical microscopy, an array of waveguides can be seen in the 264 longitudinal axis and along the axis of light propagation (Figure 5b).

265 When an electric field is applied to the hydrogel, the micro-lattice will bend according to the 266 polarity of the applied field, therefore directing the light guided by the embedded waveguides (Figure 5c). This is seen experimentally, when a blue LED ( $\lambda$  = 445 nm) launched into one end of the hydrogel (at low 267 268 intensities) shows the presence of an array of bright micro-beams with non-uniform intensity and with the 269 same periodicity of the amplitude mask used to create the pattern (Figure 5d, middle). Under the effect of 270 an electric field the waveguide bends towards the cathode while the array of micro-beams also moves in 271 the same direction (Figure 5d, left) (see video S5). Changing the polarity of the electric field after 100 s 272 produces the same effect, now in the opposite direction (Figure 5d, right). The ability of the structure to 273 direct light in the same direction of the bending is evident when looking at the zoomed areas at the top of

- each of the CCD images. This is further confirmed qualitatively when the light pattern coming out of the lattice is projected onto a black surface, in conjunction with a compact laser module light source ( $\lambda = 635$
- 276 nm) (see video S6).



Figure 5 | Remote actuation of the hydrogel micro-lattice. a) Fabrication of the hydrogel micro-lattice using a periodic amplitude mask. b) Transmission optical micrographs of hydrogel micro-lattice (H x L x W = 10 mm x 6 mm x 2.5 mm), with coloured regions highlighting areas of optical imaging. c) Bending scheme of hydrogel prism demonstrating light output based on electric field polarity (as indicated). d) Time evolution of the hydrogel waveguide lattice under the effect of an electric field (350 V/cm). The electric field was applied for 100 s, before inverting to reverse the direction of the micro-lattice. The square light pattern generated by the lattice moves towards the cathode (see video S5).

Positions of waveguides within the hydrogel lattice were analyzed using Trackpy,<sup>40</sup> a Python 285 286 particle tracking package (see Supplementary Information S4 for details regarding this program). Data 287 and image sets were collected by applying an electric field for 100 s in one direction (Figure 6a), before 288 inverting the electric field polarity (Figure 6b). Waveguide light output was detected using CCD camera 289 images, where most of the waveguide trajectories ("tracks") can be quantified as measured 290 displacements. In the first 100 s (before the electric field polarity is reversed), the whole pattern is initially 291 displaced 538 ± 23 µm (Figure 6a). After that, the electric field is turned off for 60 s to allow the hydrogel 292 micro-lattice to stabilize, and consequently allow the light output reaching the detector to steady. This 293 pattern of applying the electric field for 100 s before allowing the gel to rest for 60 s is repeated for the 294 rest of the experiment.

295 The largest overall displacement experienced by the gel is  $1031 \pm 52 \mu m$ . This is explained by 296 qualitative swelling experienced by the gel as the experiment progresses, resulting in a more flexible 297 hydrogel. When comparing the movement of all the filaments together (Figure 6c), a concerted pattern 298 emerges. The intensity output of the waveguides in the micro-lattice demonstrates a completely different 299 behavior compared to the intensity output of the rectangular and cylindrical waveguides. Unlike these 300 waveguides, which demonstrate intensity decreases and recovery as light output is directed away and 301 towards a detector (i.e., CCD camera), the intensity output of the micro-lattice shows random fluctuations 302 which indicate complex dynamics taking place within the lattice (Figure S4d).



Figure 6 | Quantification of individual waveguide displacement within the hydrogel micro-lattice. a) CCD images demonstrating the hydrogel micro-lattice moving toward the cathode. The average xdisplacement in this initial actuation (i.e., t = 0-100 s) is 538 ± 23 µm. b) After reversing the polarity of the electric field, the hydrogel micro-lattice moves in the opposite direction resulting in a x-displacement of 683 ± 35 µm (t = 100 - 200 s). c) Graph overlaying all trackable waveguide filaments x-displacements. The average x-displacement is indicated when the electric field polarity is reversed.

## 310 Conclusions

- 311 Taking inspiration from the properties of living organisms such as the bioluminescent lure of the
- 312 female angler fish and chromatophores of squid skin we have fabricated remote controllable
- 313 electroactive hydrogel waveguides for precise light control and guidance. We have demonstrated that
- their orientation, motion and thereby direction of waveguided light can be remotely controlled using
- 315 electric fields. As the method of fabrication is not restricted by the photopolymer used, this results in a

- 316 large range of versatility in the light guiding structures produced. Therefore, this work also pathways to
- 317 fabricate not only electroactive waveguide actuators in a single step but also thermal<sup>41</sup> or
- 318 photoresponsive waveguides<sup>42</sup> that can respond to more than one stimulus.

#### 319 Acknowledgements

We thank the Moran-Mirabel research group and Lyons Media Center at McMaster University for access to instrumentation. We also thank Marcus Rose for his guidance in designing the initial stages of image processing for the Trackpy program. We thank Ciba-Geigy for the generous donation of the photoinitiator, Irgacure ® 819. Funding from the Natural Sciences and Engineering Research Council,

324 Canadian Foundation for Innovation and McMaster University is gratefully acknowledged.

# 325 Methods

#### 326 Hydrogel resin formulation

The hydrogel precursor H1 was used for planar and cylindrical waveguides. H1 was obtained by mixing 67.61 mmol of acrylic acid (Sigma Aldrich, Canada), 6.94 mmol of N,N'-methylenebis(acrylamide) (Sigma Aldrich, Canada) as crosslinker, 4.48 x  $10^{-2}$  mmol of Irgacure ® 819 (Ciba Specialty Chemicals) as photoinitiator and 5 mL of sodium phosphate buffer solution 0.4 M (pH = 7.4). The hydrogel precursor, H2, used to fabricate hydrogel prisms patterned with waveguide channels was a slightly modified version of H1 precursor. In this case, Irgacure ® 819 (2.24 x  $10^{-2}$  mmol) was used in combination with  $1.20 \times 10^{-2}$ mmol camphorquinone (Sigma Aldrich, Canada).

### 334 Fabricating hydrogel waveguides

Waveguides were fabricated by launching a blue LED ( $\lambda$  = 445 nm, Thorlabs, I ≈ 0.94 mW mm<sup>-2</sup>) beam through a rectangular shape mask placed at the entrance face of cells containing hydrogel precursor H1 (Figure 1c). The beam propagated through cells with transparent coverslips as exit and entrance face with pathlengths (L) of 6, 9, and 10 mm (W x H = 2 mm x 10 mm), for 14 - 20 s, while triggering localized free radical polymerization. The beam was blocked when the polymeric waveguide reached the exit face of the cell. Once the process was completed, the coverslip at the entrance face was removed with the polymer structure still attached. This procedure was followed to fabricate cylindric waveguides with pathlength of 10 mm (Figure 3c), with the change of a circular shape mask used in
place of a rectangular one.

To obtain rectangular prisms embedded with micro-waveguides, a blue LED ( $\lambda$  = 445 nm, Thorlabs, I ≈ 0.94 mW mm<sup>-2</sup>) beam was collimated using a planoconvex lens (f.l. = 250 mm). The collimated beam was passed through a 2D periodic amplitude mask ( $\Lambda$  = 80 µm, Photomask PORTAL) before being launched through a rectangular shape mask that defined the shape of the prism (H x L x W = 10 mm x 6 mm x 2.5 mm), containing hydrogel precursor H2 (Figure 6a). All micro-patterned prisms were irradiated for 12 s.

#### 350 Unidirectional remote actuation of rectangular hydrogel waveguides

Fabricated rectangular prism waveguides were used immediately without any further treatment as they began to lose optical transparency upon swelling, becoming more brittle as their swelling ratio increased. Waveguides (W x H = 2 mm x 10 mm) 6, 9 and 10 mm long were glued to coverslips between two graphite electrodes (placed 7 mm apart). Once the cylinder was adhered to a cover slip, this unit was glued in place within a 3D printed rectangular cage that had sides sealed with glass coverslips.

356 Waveguides were partially immersed in a phosphate buffer solution (pH = 7.4, 0.4 M) that acted 357 as an electrolyte solution. A blue LED ( $\lambda$  = 447 nm, Thorlabs) beam was first collimated (L1; f.l. = 250 358 mm) before being launched through the top part of the waveguide not immersed in the buffer solution. 359 Imaging of the optical profile at the exit face of the gel was achieved with a planoconvex lens pair (L2, L3; 360 f.l. = 250 mm) and a high-resolution charge-coupled device (CCD) camera (1200 (H) x 1024 (V) pixels; pixel size = 3.2 µm horizontal x 3.2 µm). Integrated intensities were obtained by adding the intensity of all 361 362 the pixels in each image, which were then normalized to the maximum value. The bending of the hydrogel was monitored with a cellphone camera (Samsung Galaxy 9) placed at the top of the electrochemical cell. 363 364 A DC power supply (BK1735A-ND, B&K Precision) was used to generate the electric field. The polarity of 365 the electric field was inverted when the monitored intensity output for the planar waveguide reached a 366 steady, nonchanging value. This process was repeated until fatigue damaged in the hydrogel was 367 observed.

To actuate micro-patterned prisms the same assembly was used, with the carbon electrodes were placed 43 mm apart. The field strength used in these experiments was 350 V/m, with the polarity of the electric field switching every 100 s. To allow for as many waveguides to be captured on the CCD camera as possible, the electric field was turned off for 60 s after an actuation was complete. This allowed the lattice to stay within the field of view of the camera, allowing the beam output to be captured consistently over the course of the experiment. The intensity and position of each waveguide were tracked using Trackpy (see **Supplementary Information S4**).

#### 375 Unidirectional remote actuation of cylindrical hydrogel waveguides

376 We 3D printed a square prism cage (44 x 44 x 25 mm) with 4 electrode holders on each edge. 377 We placed 2 rectangular glassy carbon electrodes (25 x 25 x 3 mm) parallel to each other (Figure S3a). The cylindric hydrogel (diameter = 2 mm, L= 12 mm) was inserted through circular holes in the bottom of 378 379 the cage. The cage was sealed with a coverslip glued to the bottom. Waveguides were immersed in 380 buffer solution to a depth of 8 mm. These experiments were performed using individual cylinders, 2x2 381 arrays (4 waveguides) and 3x3 arrays (9 waveguides) of waveguides. The cylinders were illuminated with 382 a compact laser module (Thorlabs, USB Connector,  $\lambda$  = 520 nm, 0.9 mW) in conjunction with a beam 383 expander, placed underneath the cage housing the cylindrical waveguides. The time evolution of the waveguide position was followed with a Moticam 3+ camera (1600 pixels (H) × 1200 pixels (V), Motic, 384 Hong Kong). The polarity of the electric field was switched every 180 s. The electric field was turned off 385 386 for 60 s, after which the electric field was inverted. This allows cylinders light intensity output to reach a 387 steady, nonchanging value. The position of the end of the waveguide as well as the intensity of the beam output were analyzed using ImageJ.<sup>43</sup> The normalized intensity of individual waveguides was compared 388 389 to the calculated bending angles.

#### 390 Multidirectional bending of arrays of cylindrical waveguides

To achieve multidirectional bending of individual cylindrical waveguides, we combined electric fields generated from two power supplies connected to four glassy carbon plates (**Figure S3a**). Imaginary lines traced along the direction of the electric field generated by each power supply were treated like an axis. An example of the vectoral analysis carried out to combine the electric fields is shown in **Figure S3**. 395 The voltages employed to achieve each target angle as well as the polarity configurations on the carbon

396 plates are shown in **Figure 5a**. For the initial direction, the electric field was applied for 240 s, after which

the electric field was applied for 180 s. This process was repeated until the waveguide completed a full

398 cycle. It was assumed that the length of the cylinders remained 10 mm during the whole experiment. This

- 399 same procedure for multidirectional bending was used for the 2x2 and 3x3 arrays.
- 400 Structural characterization of prisms with micro-channel waveguides
- 401 A transmission microscope (Olympus BH2-UMA, Upright) coupled with a Moticam 3+ camera
- 402 was used to acquire micrographs of the channel waveguides at the entrance and exit face of the prisms.
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