

Single-particle photoluminescence measures a heterogeneous distribution of differential circular absorbance of gold nanoparticle aggregates near constricted thioflavin T molecules

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Abstract:

The chirality of biomacromolecules is critical for their function, but the optical signal of this chirality is small in the visible range. Plasmonic nanoparticles are antennas that can couple to this chiral signal. Here, we examine the molecular-scale mechanism behind the induced circular dichroism of gold nanorods (AuNRs) in a solution with insulin fibrils and the fibril-intercalating dye thioflavin T (ThT) with polarization-resolved single-molecule fluorescence and single-particle photoluminescence (PL) imaging. We compared the PL upon excitation by left- and right-handed circularly polarized light to calculate the differential absorbance of AuNRs near insulin fibrils with and without ThT. Overall, our results indicate that AuNRs do not act as chiral absorbers near constricted ThT molecules. Instead, we hypothesize that fibrils promote AuNR aggregation, and this templating is mediated by subtle changes in the solution conditions; under the right conditions, a few significantly chiral aggregates contribute to a large net circular dichroism.

Main

Amyloid fibrils are protein assemblies that often cannot be superimposed onto their mirror image¹. This signature handedness, or chirality, correlates with their biological function². If fibril formation is disrupted or the fibrils are damaged, the fibril chirality is affected^{3–6}. Thus, sensitive detection of a fibril's chirality could improve the early diagnosis of plaque-related diseases^{7–9}. Fibril chirality can be measured with circular dichroism¹⁰. This measurement typically requires high concentrations ($\sim 100 \mu\text{M}$) and large enough volumes ($\sim 1 \text{ mL}$) for enough signal. A recent study has also shown that plasmonic nanoparticles that match the fibril pitch form ordered structures around the fibril, which can be used to detect fibril chirality at low concentrations¹¹. Plasmonic nanoparticles can also amplify the resonant chiroptical signal by plasmon-coupled circular dichroism¹². Here, we probe if $80 \text{ nm} \times 40 \text{ nm}$ gold nanorods (AuNRs) that are strong enough scatterers for detection in single-particle photoluminescence (PL) microscopy and dark-field scattering spectroscopy, but much larger than the fibril pitch, can also be used as a sensitive detector for fibril chirality upon addition of the fibril-binding dye thioflavin T (ThT). We used insulin fibrils as our model fibril system and selected ThT because it acts as a fluorescent probe and a chiral molecule upon binding to the fibrils^{13,14}. The $80 \text{ nm} \times 40 \text{ nm}$ AuNRs have a strong plasmon resonance near 640 nm (Figure 1b), which is spectrally separate from the UV-range absorption peaks of fibrils and ThT^{15,16}.

We investigated the molecular-scale mechanism behind the non-resonant changes in the chiroptical signals that are detected on the ensemble level (Figure 1a, SI Figure S1) and asked if these chiral signals can be attributed to either: (1) plasmon-coupled enhancement of differential absorption by the constricted ThT molecules near the AuNRs, or (2) templated aggregation of the AuNRs by the fibrils and ThT. We designed single-particle fluorescence polarimetry experiments to distinguish between these mechanisms by measuring if there is a significant chiroptical response when a single AuNR is coupled to fibrils and proximal to fibril-bound ThT molecules¹⁰. The experiments were also designed to investigate the heterogeneity in circular differential absorption of single particles, a distribution that is hidden in the ensemble measurement (Figure 1a). We find monomeric AuNRs have no detectable induced CD signal and thus, the net CD signal cannot be attributed to the sum of the induced CD from monomeric AuNRs. Furthermore, we find that small AuNR aggregates are not preferentially handed such

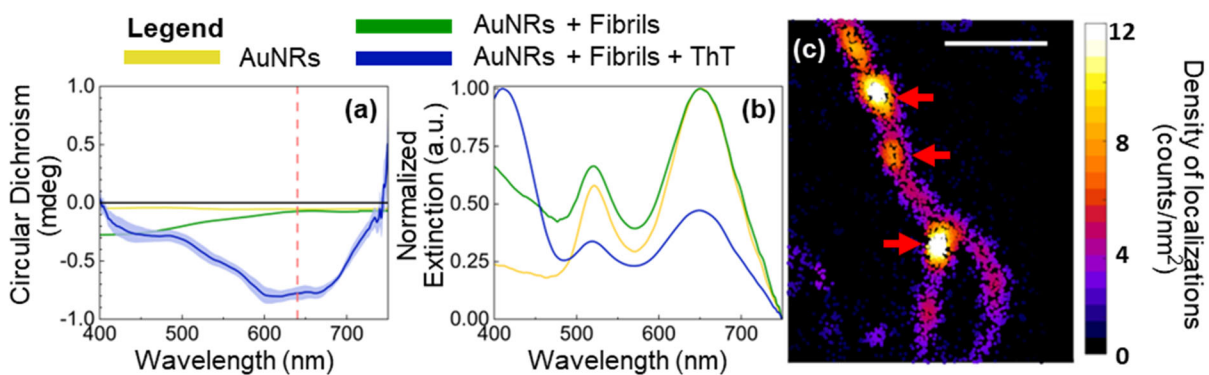


Figure 1. Spectroscopic properties of AuNRs, fibrils, and ThT. (a) Circular dichroism spectra and (b) normalized extinction spectra for solutions at pH 6. Red dotted line: 633-nm excitation wavelength for subsequent single-particle PL measurements. Light blue shading in panel ‘a’: standard deviation from five replicates. (c) Super-resolution image of ThT molecules adsorbed onto fibrils near AuNRs (red arrows) at 488-nm excitation. Scale bar: 1 μm .

that the net CD signal cannot be attributed to the sum of CD from small aggregates. Thus, the net CD signal must be dominated by a few non-representative AuNR aggregates. Our results point to the need for a high-throughput single-particle sampling detection scheme to attain a molecular-scale understanding given the heterogeneous distribution of handed aggregation of AuNRs in the presence of fibrils and ThT¹⁷.

We prepared insulin fibrils at pH 6 and mixed 100 μL of the fibrils with 300 μL of 80 nm \times 40 nm AuNRs in 40 μM of the blue intercalating dye ThT. When ThT is added to the solution of fibrils and AuNRs, there is an increase in dissymmetry around the AuNR plasmon resonance frequency of 640 nm based on ensemble circular dichroism (CD) and absorbance measurements (SI methods; Figure 1a-b). This change in absorbance at 640 nm upon the addition of a non-resonant dye is unexpected and particularly interesting because it allows us to study the coupling effects with negligible background from the dye molecules, which are much more blue: ThT molecules emit strongly upon excitation at 488 nm but are not detected upon excitation at 640 nm, whereas the AuNR PL is significant upon excitation at 640 nm.

Since only the AuNRs are optically active around 640 nm, this dissymmetry change is characterized as induced CD from the AuNRs. This increased dissymmetry has been seen in similar plasmonic nanoparticle/protein interactions^{18,19}. In solution, several factors can lead to a net change in CD, thus, we took a single-particle approach to understand the changes in the absorption of single AuNRs with changes in the surrounding media. We implemented a polarization-sensitive single-molecule microscope to measure PL-detected CD from single emitters (SI Figure S2)¹⁰. We compensated for phase changes in the optical path (primarily due to multiple reflections and the dichroic filter) and alternated between left- and right-handed circularly polarized excitations with a variable waveplate (Thorlabs LCC1223T-A). The incident beam was left-handed circularly polarized when the phase between the orthogonal axes was $+90^\circ$ and right-handed circularly polarized when the phase between the orthogonal axes was -90° . The polarization was monitored at the stage using a power meter and a linear polarizer mounted in a rotation mount and optimized to circularity $\geq 95\%$ for both handednesses in all experiments. We used a wide-field epifluorescence microscope for imaging.

First, we quantified how AuNRs affect the fluorescence of ThT bound to the fibrils. ThT is a fluorogenic dye that has an increased fluorescence quantum yield when bound to fibrils due to immobilization on these β -sheet structures. We detected the fluorescence pattern of each isolated ThT molecule as it transiently adsorbed to the fibril, used single-molecule localization to determine the apparent position of that molecule, and constructed a super-resolution map of the positions along the fibrils. (Figure 1c)^{14,20,21}. We detected ThT molecules along the fibrils, thus confirming that AuNRs do not hinder the transient adsorption of ThT onto fibrils. There appears to be a higher density of ThT molecule localizations near the AuNRs (red arrows in Figure 1c), but this observation is consistent with expected mislocalization toward the plasmonic antenna in plasmon-coupled single-molecule measurements^{22–24}. Moreover, under handed excitation, ThT fluorescence intensity does not significantly change, because the AuNRs do not have a strong optical response resonant to the dye.

We then used PL-detected CD measurements of single particles to probe if the induced CD measured at 640 nm by bulk spectroscopy is consistent with the sum of induced CD from a collection of monomeric AuNRs or if the net CD signal is dominated by a subset of AuNRs. Because PL is directly proportional to absorbance, comparing the PL intensity of the particles under left- and right-handed excitation indicates their differential absorbance. First, monomeric AuNRs were drop-cast onto a coverslip coated with poly-L-lysine. The coverslip was washed with water to remove excess AuNRs in solution and to ensure a good coverage of immobilized AuNRs. Monomeric AuNRs were identified using dark-field scattering spectroscopy²⁵ (SI Figure S3). The PL intensities of the AuNRs in water were then measured under left- and right-handed excitation (Figure 2b). An excess of fibrils was added to the AuNRs on the microscope coverslip and the fibrils were left to settle and stick to the slide for an hour before measuring the PL intensity of the same AuNRs under left- and right-handed excitation (Figure 2c). We then added an excess of ThT to the AuNRs and fibrils on the microscope coverslip. The ThT molecules adsorbed onto the fibrils immediately, so we measured the PL intensity of the same AuNRs under left- and right-handed excitation right after adding ThT (Figure 2d). We compared the PL brightness (PL intensity normalized by the acquisition time) under left- and right-handed

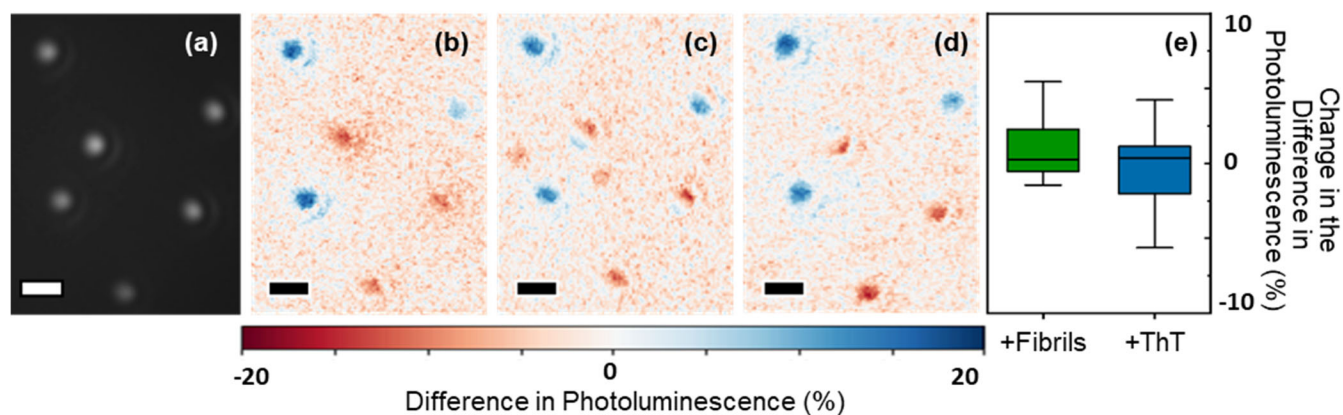


Figure 2. No consistent change is measured for the differential photoluminescence (PL) of AuNRs in the vicinity of fibrils and ThT. (a) Integrated PL image (excitation wavelength: 635 nm). (b – c) Plots of percentage difference of PL intensity between LCP and RCP excitation for the same field of view: (b) only AuNRs, (c) sample in ‘b’ after the addition and adsorption of fibrils, (d) sample in ‘c’ after the addition of 40 μM ThT. Scale bars: 1 μm . (e) Change in percentage difference in PL of AuNRs with the addition of fibrils (i.e., panel ‘c’ minus panel ‘b’) and ThT (i.e., panel ‘d’ minus panel ‘c’) for $N = 18$ AuNRs. Box limits: 25 – 75% range; center lines: means; whiskers: data extrema.

excitation over each diffraction-limited spot. The percentage difference in PL, ΔPL (Figures 2 and 3) was calculated as:

$$\Delta PL = \frac{2(I_L - I_R)}{(I_L + I_R)} \times 100 \% \quad (1)$$

where I_L and I_R are the PL brightnesses upon left- and right-handed excitation, respectively.

If the bulk signal is consistent with the effects of single AuNRs, we would expect to see a decrease in the ΔPL of the monomeric AuNRs (Figure 2a) after adding fibrils and ThT, because there is a strong dip in the ensemble CD signal at 640 nm (Figure 1a). However, ΔPL does not change significantly when fibrils are added (Figure 2c), nor when ThT binds to the fibrils (Figure 2d). This characterization was performed for $N = 18$ particles (Figure 2b-d; SI Figure S4). Thus, we observed no significant change in ΔPL between the three conditions, and no consistent change was recorded over $N = 18$ particles (Figure 2e). Thus, we conclude that the net CD is not the sum of the induced CD from each monomeric AuNR.

Previous work has also shown that fibrils and small complexes favor the formation of AuNR aggregations that tend to be handed^{19,26}. When mixed with fibrils, the 80 nm \times 40 nm AuNRs in this experiment are observed in transmission electron microscopy (TEM) images to aggregate in a variety of configurations (Figure 3e), thus our second hypothesis was that the bulk CD signal is dominated by a subset of AuNRs that form small chiral aggregates with a favored handedness upon mixing with fibrils and ThT. This hypothesis is consistent with similar chiral templating that has been observed for other small proteins²⁷⁻²⁹. In earlier experiments, we avoided aggregation by sequentially adding components to focus on measuring the effect of nearby fibrils and ThT on the induced CD from monomeric AuNRs. Here, we instead formed agglomerates and tested their handedness by mixing AuNRs in an Eppendorf tube either with a solution of only fibrils (Figure 3a-b) or with a solution of fibrils and ThT (Figure 3c-d). By mixing in a tube before drop-casting onto the coverslip, we allowed aggregates to form in solution; these aggregates were then immobilized onto a poly-L-lysine-coated coverslip. The coverslip was washed gently with Milli-Q water to remove any fibrils or AuNRs that were not immobilized. PL images were taken under left and right circularly polarized excitation for both these solutions (Figure 3a,c). The ΔPL for each sample was calculated using Equation (1) (Figure 3b,d). We

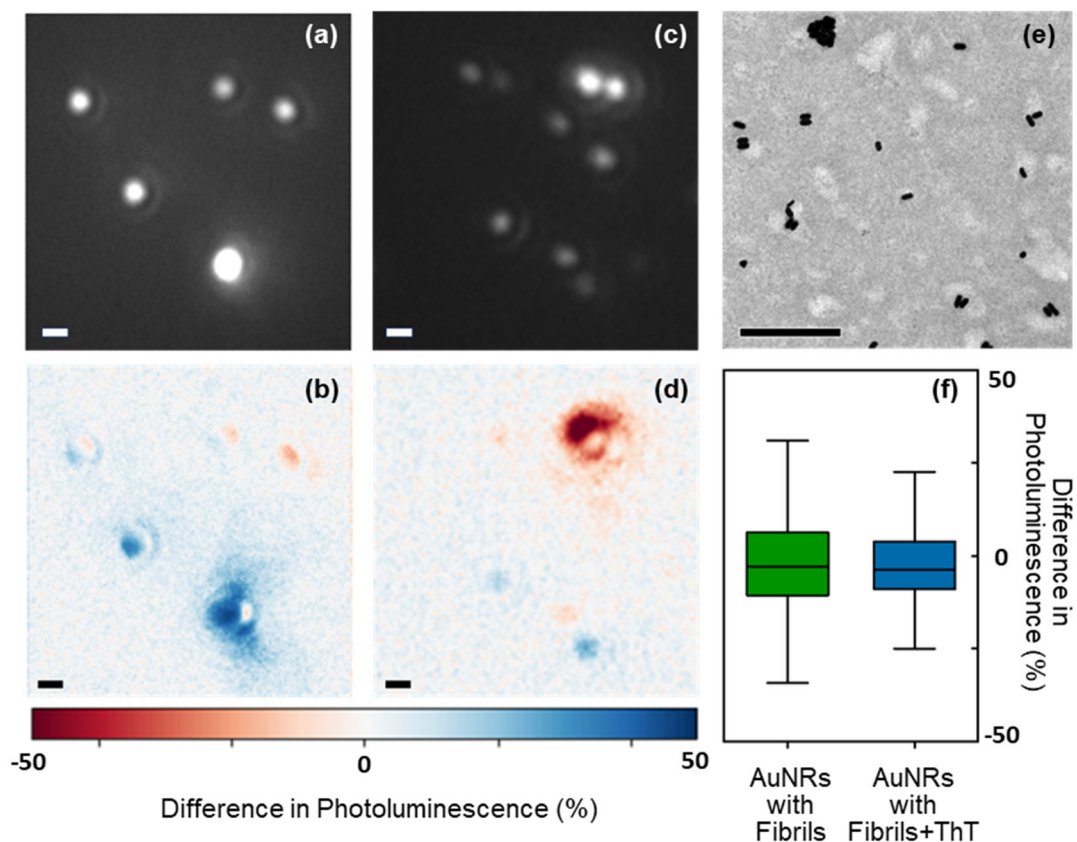


Figure 3. AuNR aggregation upon mixing with fibrils and ThT in an Eppendorf tube at pH 6. (a – b) AuNR aggregates in a fibril solution: (a) Integrated photoluminescence (PL) image (excitation wavelength: 635 nm). (b) Plot of percentage difference of PL intensity between LCP and RCP excitation for the same field of view. (c – d) AuNR aggregates in a fibril and ThT solution: (c) Integrated PL image (excitation wavelength: 635 nm). (d) Plot of percentage difference of PL intensity between LCP and RCP excitation for the same field of view. (e) TEM image of AuNR aggregates formed in a solution of fibrils and ThT at the same conditions. Scale bars: 1 μm . (f) Percentage difference in PL intensity between LCP and RCP excitation of AuNR aggregates formed in solution with: fibrils at pH 6; or fibrils and ThT at pH 6 for $N = 30$ aggregates. Box limits: 25 – 75% range; center lines: means; whiskers: data extrema.

observed that these aggregates have much higher magnitudes of ΔPL , and are thus optically more chiral, than the monomeric AuNRs in Figure 2, but they are not preferentially templated into right- or left-handed aggregates (Figure 3f); any observed bias is within our measurement error (SI Figure S5). From the ensemble CD measurement (Figure 1a), we had expected a skew towards more right-handed aggregates when the AuNRs were mixed with the fibrils and the ThT, but our single particle measurement suggests that the ThT does not play such a big role in the chirality of small aggregates, and in both solutions, the net effect of the small aggregates on the ensemble CD would thus be zero.

Since our experimental setup only captures the effect of fibrils and ThT on AuNR monomers (Figure 2) and small AuNR aggregates with diffraction-limited point spread functions immobilized on poly-L-lysine coated coverslips (Figure 3), our observations exclude any larger 3D chiral structures. Thus, our single-particle results indicate that the CD signal is dominated by the chiroptical response from only a subset of large AuNR aggregates and is not representative of the sum of the signals from monomeric AuNRs coupled to fibrils and constricted ThT molecule, nor to small aggregates of AuNRs.

In conclusion, we have designed and implemented a single-particle polarimeter sensitive to a 5% difference in PL to conclude that the introduction of an excess of intercalating dye is not enough to enable the detection of the chirality of fibrils at the single-particle level. There is no significant and consistent enhancement of differential absorption by the monomeric AuNRs due to ThT bound to nearby fibrils, and there is a heterogeneous distribution of handedness in the aggregates. The skew toward a specific fibril handedness is negligible for small aggregates. Additional aggregation factors, such as salt concentrations and pH, need to be explored before using these 80 nm \times 40 nm gold nanorods to build single-particle chirality sensors. In addition to directly probing chirality in biological systems, the microscope setup can also be used to study the chiroptical properties of more elaborate substrates like chiral antenna arrays that can be deployed for sensing by redirecting the emission from molecules.

Supporting Information

Supporting methods and figures detailing spectroscopic properties, single-particle polarimeter set-up, and PL intensity differences.

Conflict of Interest Disclosure

The authors declare no competing financial interest.

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TOC Figure

