An Explicit-Implicit Hybrid Solvent Model for Grand 1 Canonical Simulations of the Electrochemical Environment 2 3 Duy Le^{1,2} 4 ¹ Department of Physics, University of Central Florida, Orlando, FL 32826 5 ² Renewable Energy and Chemical Transformations Cluster, University of Central Florida, 6 7 Orlando, FL 32826 8 Email: duy.le@ucf.edu 9 10 Abstract: The development of *ab initio* methods for atomistic simulations of the electrochemical 11 environment is essential for obtaining a mechanistic understanding of the fundamental reactions. 12 We propose here an explicit-implicit solvent model, SOLHYBRID, that enables grand-canonical 13 ensemble simulations of the electrochemical environment with the popular Vienna Ab initio 14 Simulation Package (VASP), extending its capabilities beyond the commonly used the implicit 15 solvent model VASPSol. We further present a subroutine, TPOT, that allows control of the 16 electrode potential, thereby enabling simulations at constant electrode potential to mimic the experimental electrochemical cell. We demonstrate the key points of our approach for the case of 17 18 CO_2 adsorption on Au(110) in the presence of K⁺ cation.

19 **1 INTRODUCTION**

Because of the rising demand for clean energy and CO_2 recycling, electrochemical reduction of CO₂ under ambient conditions (room temperature and atmospheric pressure) has emerged as a promising route for *hydrogen-free* sustainable fuel production, as it is compatible with the 23 intermittency of renewable energy sources and enables utilization of renewable electricity without 24 the need for expansion of transmission capacity. With the development of an efficient 25 electrochemical process for CO₂ recycling, such *clean energy* technology can ultimately close the 26 carbon cycle for the utilization of carbon-based fuels and reduce carbon emissions to significantly 27 mitigate the anticipated climate and environmental damage. While great progress has been made in recent times, as so well-articulated in a collection of Road Maps,¹ a number of challenges need 28 29 to be overcome before controlled electrocatalytic conversion of gases such as CO₂ becomes an 30 industrial reality.

Electrochemical reactions occur inside an environment that is complex because of the presence of multiple interfaces involving the liquid solvent, ions, gaseous reactants, and solid electrodes. To add to the complexity, the potential applied to the electrode makes the electrocatalysis process potential-dependent. Because of these heterogeneous interfaces, an atomistic understanding of the electrocatalysis process is still unclear partially because of the lack of an adequate theoretical method that incorporates the above complexities in a realistic manner.

37 As noted in several excellent reviews, significant advances have been made in theoretical and computational modeling of electrocatalysis.²⁻⁴ In a nutshell, with reliance on *ab initio* techniques 38 39 based on density functional theory (DFT), simulations of the electrochemical system may proceed 40 along several strategies. The most computationally efficient approach is the implicit solvent model⁵⁻⁹ in which a continuum description is applied to the electrolyte and the ions. Next are the 41 explicit solvent models, which treat water and ions atomistically.¹⁰ Then there is the H-shuttling 42 method,¹¹⁻¹³ in which a water molecule shuttles protons to and from the electrode surface. It may 43 44 also be possible to include the effect of the electrode potential in the calculations through the socalled capacitor model,¹⁴ but such an approach fails to account accurately for the response of the 45

46 charge on the electrode to the dynamical evolution of the adsorbate, electrolyte, and the catalyst 47 surface to keep the electrode potential constant. Fortunately, such dynamical effects are naturally 48 included in the self-consistent grand canonical DFT (GC-DFT)^{15, 16} method, but its implementation 49 has been hampered by high computational cost. Efforts^{17, 18} have also been devoted to the develop 50 the grand-canonical AIMD (or constant potential) simulations, but so far these methods have not 51 been widely applied.

Explicit inclusion of all components of the electrolyte, that is all water molecules, the cations, and anions, could faithfully describe the electrochemical environment. However, such a simulation with *ab initio* methods would be computationally prohibitive. In our opinion, an implicit-explicit hybrid model, in which the region near the electrode is modeled explicitly while an implicit solvent model is used elsewhere, would provide insights into electrochemical processes while keeping the computational cost manageable. In addition, such a simulation could be performed at a constant electrode potential to reflect the conditions of the "real" world electrochemical reactions.¹⁹

59 VASPSol^{5, 6} is a popular implicit solvent model used for simulations of the chemistry in 60 electrochemical environments. It treats the solvation effect via a continuum model. Intuitionally, 61 one would think that the hybrid model mentioned above could be attained by simply adding 62 additional explicit water molecules and ions to the simulations. Could we possibly use the implicit 63 solvent model VASPSol as the basis for creating a hybrid solvent model?

The answer to this question is YES and NO because of the definition of the shape function,^{5, 6} which is used to distinguish the solute and electrolyte in a simulation cell in VASPSol. As we explain in detail below, the shape function defines the electrolyte region through its low electron density. If the explicit solvent and ionic regions extending to a couple of angstroms from the electrode surface are pasted in VASPSol, there are regions inside the explicit electrolyte with low or no electron density because the solvent molecules and ion do not bond to each other via covalent
bonds. In VASPSol these regions will be filled with the implicit electrolyte, which would be a
problem.

In this work, we propose our SOLHYBRID model which utilizes a padding charge for the explicit solvent region and an economical algorithm (TPOT) for performing simulation at a constant potential. We will also use the proposed methods to perform the *ab initio* molecular dynamics (AIMD) simulation of the adsorption of CO_2 on Au(110) surface in the presence of potassium cation (K⁺) to demonstrate the viability of our proposed SOLHYBRID model and TPOT algorithm.

78 2 COMPUTATIONAL DETAILS

We perform density functional theory (DFT) based calculations using the VASP package^{20, 21} with 79 the implicit solvent model VASPSol,^{5, 6} employing the plane-wave supercell and the projector 80 augmented-wave methods.^{22, 23} We used the generalized gradient approximation (GGA) in the 81 form of Perdew-Burke-Ernzerhof (PBE)^{24, 25} functional and the DFT-D3 van der Walls 82 correction²⁶ for evaluating exchange-correlation of electrons. We set the cutoff energy for plane-83 84 wave expansion at 500 eV. We chose a Debye length of 3.04 Å, which corresponds to a bulk 85 electrolyte concentration of 1 M, and a relative permittivity of the solvent of 78.4 (for water at the 86 ambient condition) for all calculations with the implicit solvent model VASPSol. Our simulation supercell consists of a 5-layer, 2×3 Au(110) slab as the modeled electrode, 26 H₂O molecules, 87 one K^+ , a CO₂ molecule, and a vacuum of 15 Å to separate the periodical images along the normal 88 89 direction of the electrode. We refer to this supercell as CO_2 -K⁺/Au(110) in what follows. We

90 sampled the Brillouin zone with a $3 \times 3 \times 1$ grid and used Gaussian Smearing with $\sigma = 0.1$ eV for integration over the Brillouin zone. The threshold for convergence of electronic step is 10^{-6} eV. 91 AIMD simulations were performed at 298.15 K in the NVT ensemble with the Nosé-Hover 92 thermostat.²⁷ During the simulations, the bottom three layers of the Au electrode are held fixed. 93 94 Note that the initial supercell was prepared in a way that binds CO_2 on the Au(110) surface with 95 Au-C bond length of 2.15 Å, CO₂ bending angle (\widehat{OCO}) of 128.5°, above which there is a K atom with a distance of 3.58 Å to an O atom of CO₂, and 27 H₂O molecules. We performed AIMD 96 simulations for 2 ps with low accuracy, i.e., 300 eV energy cutoff, 10⁻⁵ eV threshold for electronic 97 98 convergence, Gamma point sampling for Brillouin zone, fixing all Au, C, and K atoms. After this 99 step, we observed that one H₂O molecule was above all other H₂O molecules. We remove this 100 molecule and continue the AIMD simulation for 2.5 ps. The resulting structure, after releasing the 101 C and K atoms and atoms in the top two layers of the Au(110) slab, was the initial configuration 102 of this work.

103 3 MODEL AND IMPLEMENTATION

In this section, we will introduce our hybrid solvent model, SOLHYBRID, that mitigates the issues with VASPSol in simulating the explicit solvent as discussed in the Introduction. In addition, we will present our TPOT algorithm which allows simulations at a target potential at an economical computational cost. The results of the simulations of CO_2 -K⁺/Au(110) system demonstrate the novelty and effectiveness of our proposed methodology.

109 3.1 Explicit-implicit hybrid solvent model SOLHYBRID

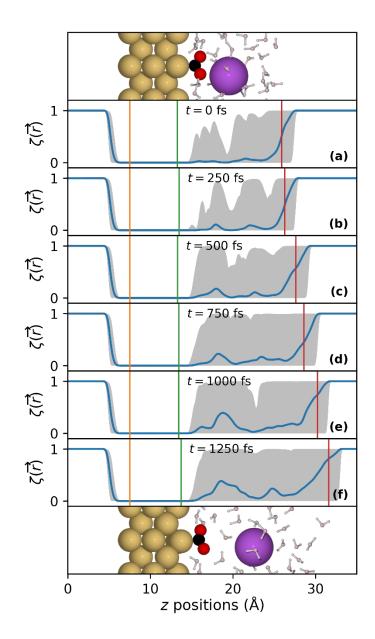


Figure 1. Evolution of shape function (VASPSol) during AIMD simulations of CO_2 -K+/Au(110). The planar112average of the shape function is shown as blue lines. The shaded area indicates the range of value of the113shape function. Orange, green, and red vertical lines indicate the position of the lowest Au atom, the highest114Au atom, and the highest atoms of the explicit solvent molecules, respectively. The ball-stick models of the115atomic configurations at 0 fs and 1250 fs are shown at the top and bottom of the figures, respectively.

Yellow, black, red, and purple balls represent Au, C, O, and K atoms. The H₂O molecules are shown by the
pink-while ball-stick molecules.

118 In VASPSol, the spatially dependent shape function is defined as:

119
$$\zeta(\vec{r}) = \frac{1}{2} \operatorname{erfc} \left\{ \frac{\log[n(\vec{r})/n_c]}{\sigma\sqrt{2}} \right\};$$

120 where $n(\vec{r})$ is the pseudo electron density calculated as the sum of core electron density⁶ $n_{core}(\vec{r})$ 121 and valence electron density $n_{val}(\vec{r})$, at coordinate \vec{r} ; n_c is the cutoff electron density; σ is the 122 width of the diffuse cavity or interface region. The shape function is used to identify the region 123 that is treated purely by DFT ($\zeta(\vec{r}) = 0$) or by implicit model ($\zeta(\vec{r}) > 0$).

124 With this definition, the value of shape function in the region with electron density smaller than a 125 few n_c will be not zero and this region will be filled with explicit solvent molecules and explicit 126 ions. Figure 1a shows the shape function for a hybrid solvent model, i.e., explicit H_2O and cation 127 near the Au(110) electrode and implicit solvent elsewhere. Because of the low-density region 128 between H₂O molecules, the value of shape function in the explicit region, between the vertical 129 green and red lines in Figure 1, is not always zero. VASPsol will fill the region with implicit water molecules and ions. Since technically, the explicit water molecules and ions do not have a 130 131 preference to be in either explicit or implicit regions, the explicit H₂O and ions will eventually be 132 separated by implicit solvent. Figure 1a-f shows the evolution of the shape function and the 133 structure of CO₂-K⁺/Au(110) system for about 1.25 ps AIMD simulations. The range (shaded area 134 in Figure 1a-f) of the shape function keeps growing and the explicit solvent layers continue to 135 expand to fill the vacuum region. This picture is not fundamentally correct because the implicit 136 solvent should be expelled away from the explicit region.

137 To mitigate the above problem, we propose our SOLHYBRID model by introducing a padding density $n_{pad}(\vec{r})$ to the pseudo electron density to ensure that the implicit solvent does not fill the 138 139 explicit region. We have tried three strategies. In the first strategy, we recycled the part of the VASPSol code that calculates Gaussian core electron density $n_{core}(\vec{r})$ to calculate $n_{pad}(\vec{r})$ as the 140 sum of Gaussian charge centered at atoms of the system, similar to the calculation of $n_{core}(\vec{r})$ but 141 142 with the number of valence electrons instead of the number of core electrons and a large gaussian 143 width. We found that strategy is not an ideal choice as the majority of the padding charge localizes 144 at the center of each atom and that an appropriate gaussian width must be chosen to ensure that the 145 padding charge extends to cover the space between molecules. The latter is not always satisfied 146 because of unforeseen movements of molecules during simulations. We thus decided to remove 147 this strategy from our implementation. The other two strategies are summarized as the two 148 Schemes below.

149 Scheme 1: $n_{pad}(\vec{r})$ is the planar average of valence electron $n_{val}(\vec{r})$ which is smoothened with a 150 Gaussian filter with a width σ_{SH} and the pseudo electron density is defined as:

151
$$n(\vec{r}) = n_{core}(\vec{r}) + (1 - \alpha_{SH})n_{val}(\vec{r}) + \alpha_{SH}n_{pad}(\vec{r});$$

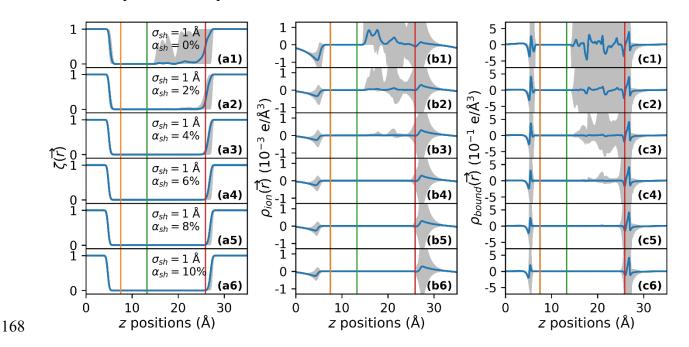
where α_{SH} is a parameter that controls the percentage of valence density is replaced by padding density; $n_{pad}(\vec{r})$ is the planar average of $n_{val}(\vec{r})$. The optimizations of α_{SH} and σ_{SH} are necessary for building a suitable SOLHYBRID model for a specific system. In our opinion, this Scheme is the ideal one as it is calculated from the electron density of the system and has been tested to be stable during simulations.

157 Scheme 2: $n_{pad}(\vec{r})$ is a uniform density in the region of solute and explicit solvent and decays to 158 zero at the interface with the implicit region and the pseudo density is defined as:

159
$$n(\vec{r}) = n_{core}(\vec{r}) + n_{val}(\vec{r}) + \alpha_{SH}n_{pad}(\vec{r});$$

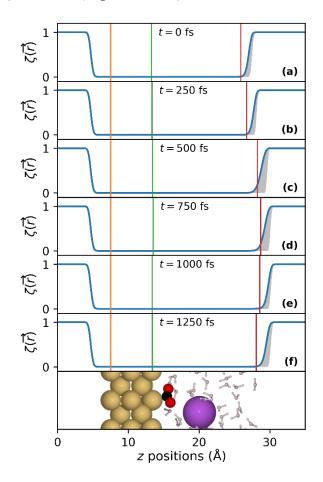
160 where padding charge density $n_{pad}(\vec{r})$ is defined as $n_{pad}(\vec{r}) = \operatorname{erfc}\left\{\frac{z_{min}-z}{\sigma_{SH}\sqrt{2}}\right\} - \operatorname{erfc}\left\{\frac{z_{max}-z}{\sigma_{SH}\sqrt{2}}\right\};$

161 where z_{min} and z_{max} are the *z* coordinates of the lowest and highest points of the region where 162 the padding charge is needed (to define the explicit solvent region). α_{SH} and σ_{SH} are parameters 163 that control the amount of padding charge density to be added to pseudo density and the steepness 164 of padding density near z_{min} and z_{max} . Our tests show that Scheme 2 generally works well but 165 the choice does not reflect the electron density of a physical system. While we have implemented 166 Scheme 2 as an option, we focused mainly on Scheme (1) and used it to produce the data presented 167 here as our proof of concept.



169 **Figure 2**. The dependence of the shape function in SOLHYBRID model on α_{SH} (a1-a6) with $\sigma_{SH} = 1$ Å. 170 Corresponding implicit ion density (ρ_{ion}) and bound charge density (ρ_{bound}) for each value of α_{SH} are 171 shown in (b1-b6) and (c1-c6), respectively. See **Figure 1**'s caption for more descriptions of lines and 172 shaded areas.

Figure 2(a1-a6) shows the shape function in SOLHYBRID model with $\sigma_{SH} = 1$ Å and $\alpha_{SH} =$ 173 174 0, 2, 4, 6, 8, 10% of the same structure, CO₂-K⁺/Au(111), that was used to calculate the shape function shown in Figure 1a. In the case of $\alpha_{SH} = 0$, the new shape function is identical to the 175 176 original one, as expected. With increasing α_{SH} , the value of the shape function in the explicit region reduces (the shaded area diminished). We can achieve zero value in the explicit region with α_{SH} = 177 178 4%. However, our goal is not only to achieve that but also to repel implicit ion⁵ and implicit bound charge^{5, 28, 29} from the explicit region. We found that the implicit ion density in the explicit region 179 approaches zero with $\alpha_{SH} = 6\%$ (Figure 2b1-b6) and that we can only get rid of the bound charge 180 181 in explicit region with $\alpha_{SH} = 10\%$ (Figure 2c1-c6).



183 Figure 3. Evolution of shape function with SOLHYBRID model during AIMD simulations of CO₂-

184 $K^+/Au(110)$. The ball-stick model of the atomic configuration at 1250 fs is shown at the bottom of the figure.

185 See *Figure 1*'s caption for more descriptions of lines, shaded areas, and color codes for atoms.

We used the SOLHYBRID model with $\sigma_{SH} = 1$ Å and $\alpha_{SH} = 10\%$ to perform AIMD simulations for the CO₂-K⁺/Au(110) system, with the same starting configuration of the AIMD with original VASPSol (shown in **Figure 1**). **Figure 3** summarizes the evolution of the shape function and the structure. Unlike the continuing expansion of the explicit solvent during the simulation shown in **Figure 1**, the system expands at the beginning, which is normal, and stops expanding at around 750 fs into the simulation. More importantly, values of shape function in the explicit region consistently remain zero, effectively expelling implicit solvent out of this region.

193 It is important to emphasize that while we discussed only the implementation of SOLHYBRID 194 model into VASPSol, the approach, i.e., using a padding density, could also be applied to other 195 implicit solvent models that use electron density to define the regions of solute or electrolyte, such 196 as those implemented in ENVIRON,^{8, 30-32} those originated from joint DFT framework,³³⁻³⁵ SaLSA 197 solvation model,³⁶ CANDLE solvation model,⁷ to make them work for simulations of explicit-198 implicit solvents.

3.2 Controlling the electrode potential with TPOT

200 TPOT, Target **POT**ential, is a routine that runs in conjunction with the VASP code and VASPSol 201 to control the number of electrons during a simulation required to reach a target electrode potential. 202 The electrode potential is defined as $\Phi = (\Phi_{Vac} - \mu_e)/e$; where Φ_{Vac} is energy level at the point 203 far away from the electrode, i.e., vacuum level, μ_e is the chemical potential of electrons, and e is 204 the elementary charge. In principle, the numbers of electrons can be optimized during the self-consistent field (SCF) cycles to obtain a target potential at the end of each ionic iteration.^{15, 37, 38} Even though we implemented this approach in TPOT, we found that it requires substantial number of iterations and increases the computational cost.

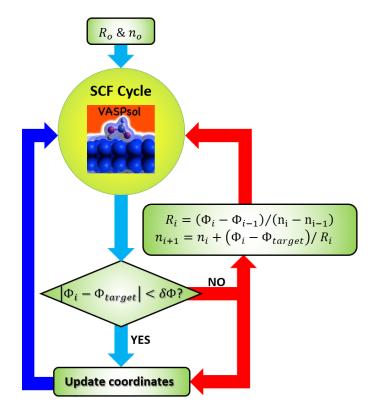


Figure 4. Flowchart of TPOT. R_o and n_o are initial guesses of the rate of change of electrode potential with respect to the change of number of electrons and the initial number of electrons, respectively. Φ_i , Φ_{target} , and $\delta \Phi$ are electrode potential at ionic iteration i, target potential, and the error threshold of electrode potential. R_i and n_i are the rate of change of electrode potential and number of electrons at ionic iteration i.

We also found that the optimization of number of electron during SCF cycle is not necessary because, during ionic relaxations or molecular dynamics simulations, the atomic structures in two consecutive steps are not substantially different, pointing to no substantial change in the electrode

218 potential. We, thus, propose *not to* optimize the number of electrons during SCF cycles but after 219 each ionic iteration. The flowchart in Figure 4 shows that TPOT needs to optimize the number of 220 electrons only if the resulting electrode potential (Φ) differs from the target potential by an amount that is larger than a predetermined threshold $\delta \Phi$ (in practice, we set $\delta \Phi = 0.001$ V). This approach 221 sets TPOT apart from other approaches that have been proposed^{15, 37, 38} for performing constant 222 223 potential calculations. It saves computational time as it requires virtually no additional cost as compared to the constant charge calculation as implemented in VASPSol.^{5, 6} While, in principles, 224 225 advanced methods for optimizing the number of electrons can be developed and implemented, we 226 use a simple approach to update the number of electrons at iteration *i* using the following algorithm $n_{i+1} = n_i + (\Phi_i - \Phi_{target})/R_i$ only if $|\Phi_i - \Phi_{target}| > \delta\Phi$, where $R = \partial\Phi/\partial n$ is the rate of 227 the change of electrode potential with the change in number of electrons. The rate is updated at the 228 last time that the number of electrons is updated. R_i is calculated as $(\Phi_i - \Phi_{i-1})/(n_i - n_{i-1})$ or 229 230 is provided as input (R_o) in first iteration.

231 3.3 Additional notes

232 During the implementation of the SOLHYBRID and TPOT, we noticed two important issues 233 related to VASPSol that should be mentioned. The first issue is about the calculation of core electron density $n_{core}(\vec{r})$. VASPSol uses gaussian type core electrons to repel implicit solvents 234 from the core of atoms.⁶ The default value, set by VASPSol, of the number of core electrons is 0 235 236 for atoms with atomic number that is smaller than 10 and 1 otherwise. If the default values are 237 used, some atoms will have empty cores that may be filled with the implicit solvent by VASPSol. 238 With our SOLHYBRID model, this is not an issue as the padding charges will fill these core 239 spaces. Nevertheless, in our modified version of VASPSol, the values of core electrons for all species, as well as all parameters used in the model, are printed out in the main output file (i.e.,
OUTCAR) for alerting users. We also suggest explicitly specifying core electrons in input (i.e.,
INCAR) to avoid unintended outcomes.

The second issue is the calculation of the vacuum level (Φ_{Vac}). VASPSol calculated a correction for aligning the vacuum level to zero (i.e., FERMI_SHIFT).⁵ This value is regarded as $-\Phi_{Vac}$.^{5, 14} However, we tested and found that for charged systems, there is a small difference between FERMI_SHIFT and $-\Phi_{Vac}$. Thus, in our implementation, we calculated Φ_{Vac} from the electrostatic potential of the system. TPOT offers two options for defining electrode potential $\Phi =$ $(\Phi_{Vac} - \mu_e)/e$ and $\Phi = -(\mu_e + \text{FERMI_SHIFT})/e$.

249 In addition, when we first realized the issue with VASPSol in simulating the explicit-implicit 250 model, we implemented a penalty potential to prevent the explicit solvent molecules from 251 desorbing from the explicit region. The penalty potential applied to atom I takes the form of a softplus function as $E_i^{Pen} = h_{Pen} \log\{1 + \exp[(z_i - z_o)/w_{Pen}]\}$; where z_i is the z coordinate of 252 atom i, z_o is the center of the softplus function, h_{Pen} and w_{Pen} define the high and width of the 253 254 penalty potential. We later realized that this is not a solution as it does not fix the root cause of 255 VASPSol, i.e., the implicit solvent can fill the explicit region causing it to expand and apply stress 256 to the electrode. We keep this implementation for rare instances in which explicit solvent 257 molecules may desorb from the explicit region into the implicit one because in principles it should 258 have no preference to be in either region. Note that none of the simulations presented in this paper 259 were produced with the penalty potential described above.

260 4 DEMONSTRATION OF GRAND CANONICAL AIMD

In this section, we demonstrate the use of the SOLHYBRID model and TPOT for grand canonical (GC) AIMD simulations. We consider the case of CO_2 adsorption on the Au(110) electrode with the presence of one potassium cation K⁺ as our toy model. We choose this particular system because it has been shown that with one K⁺ the CO₂ does not stay adsorbed on the Au(110) surface during AIMD simulation with pure DFT calculations.³⁹

To prepare the starting point for AIMD simulation, we first performed GC-AIMD for about 9 ps at -1 V vs RHE with the starting configuration described in Computational Details section to obtain a configuration in which CO_2 chemisorbed on the electrode, i.e., bonded to Au surface atom with OCO bending angle that is significantly smaller than 180°, with a K⁺ nearby.

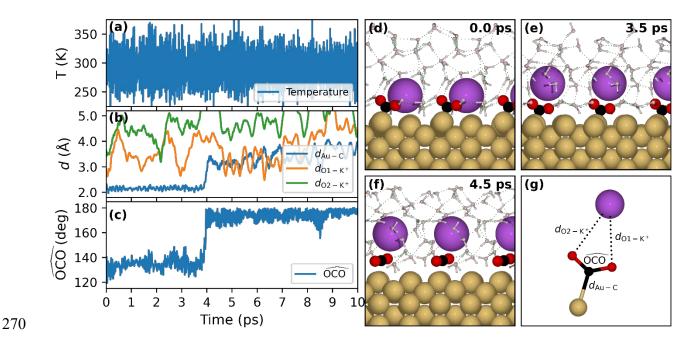


Figure 5. Evolution of (a) Temperature, (b) distance from Au and C atom of CO₂ molecule (d_{Au-C}), distance from K⁺ to two O atoms of CO₂ molecule (d_{O1-K^+} and d_{O2-K^+}), (c) bending angle of the CO₂ molecule (\widehat{OCO}) during the AIMD simulation of CO₂-K⁺/Au(111) with standard DFT. Snapshots of the

simulations are shown in (d-f). Quantities shown in (b-c) are illustrated in (g), where all water molecules and Au atoms that are not the binding site for CO_2 were omitted for clarity. See **Figure 1**'s caption for the descriptions of color codes for atoms. The hydrogen bonding network is shown with green dot lines.

277 Figure 5 shows the results of AIMD simulation of the system with standard DFT (PBE functional, 278 neutral charge supercell). The movie of this simulation is provided in Supplementary Video 279 Movie K+CO2 DFT.mp4. A few snapshots of the movies are shown in Figure 5d-f. We found 280 that the CO₂ chemisorbed configuration is stable on the surface for about 4 ps, after which point, 281 the CO₂ desorbs from the surface, which is evidenced by the increase Au-C bond length from \sim 2.1 Å to > 3.0 Å (Figure 5b) and by the increase of \widehat{OCO} angle from around or below 140° to near 282 283 180° (Figure 5c). The result that the CO₂ desorbs from the Au surface after a short simulation agrees with 284 a similar simulation reported by Qin et al.³⁹

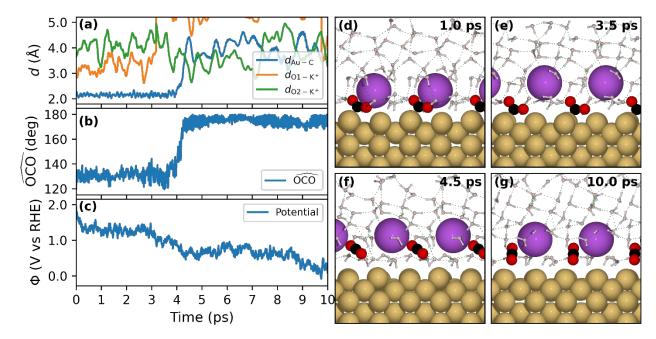
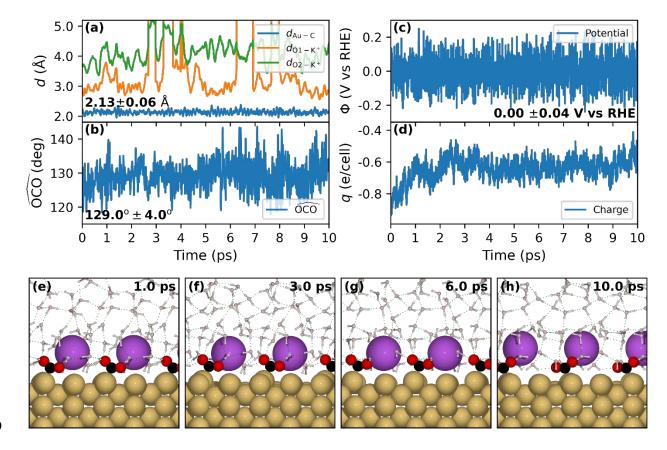


Figure 6. Evolution of **(a)** distance from Au and C atom of CO_2 molecule (d_{Au-C}) , distance from K^+ to two O atoms of CO_2 molecule $(d_{O1-K^+} and d_{O2-K^+})$, **(b)** bending angle of the CO_2 molecule (\widehat{OCO}) , (c) potential of Au(110) electrode (Φ) during the AIMD simulation of CO_2 -K⁺/Au(111) with SOLHYBRID model.

Snapshots of the simulations are shown in (d-g). See Figure 1's caption for the descriptions of color codes
for atoms. The hydrogen bonding network is shown with green dot lines.

291 We next performed a similar AIMD simulation but with SOLHYBRYD model with a neutral 292 charge supercell. The results are summarized in Figure 6. The movie of this simulation is provided 293 in Supplementary Video Movie K+CO2 SOLHYBRID.mp4. A few snapshots of the movies are 294 shown in **Figure 6d-g**. Once again, we see that the CO_2 molecule starts to desorb from the Au(110) 295 electrode at around 4 ps second into the simulation as evidenced by the increased Au-C distance 296 (Figure 6a) and \widehat{OCO} angle (Figure 6b) in a way that is similar to AIMD simulation with standard DFT 297 shown in Figure 5. The agreement between the two simulations validates our SOLHYBRID model which 298 aims to simulate explicit solvent near the electrode surface. Thanks to our new implementation, we can 299 calculate the electrode potential on-the-fly, as shown in Figure 6c. It is worth noting that the potential of 300 the Au electrode at the beginning of the simulation is around 2 V vs RHE and decreases during the 301 simulation to approach 0 V vs RHE at 10 ps. This variation of electrode potential is understandable because 302 CO_2 requires electrons to be donated to its antibonding π^* orbital to form a bent configuration and to bind 303 on the Au surface. The donated electrons are from the Au(110) surface thus increasing its electrode 304 potential. Upon CO₂ desorption, it returns those electrons to the electrode thereby reducing its electrode 305 potential. This behavior, i.e., variation of electrode potential, should not happen in a grand-canonical 306 environment in which the electrode is supplied with electrons from an electron source (reservoir) during 307 electrochemical reactions, i.e., the electrons donated to CO_2 are then replenished with electrons from the 308 reservoir.



309

Figure 7. Evolution of (a) distance from Au and C atom of CO₂ molecule (d_{Au-C}) , distance from K⁺ to two O atoms of CO₂ molecule $(d_{O1-K^+} and d_{O2-K^+})$, (b) bending angle of the CO₂ molecule (OCO), (c) potential of Au(110) electrode (Φ) , and (d) charge in the supercell (q) during the GC-AIMD simulation of CO₂- $K^+/Au(111)$ with SOLHYBRID model at 0 V vs RHE. Snapshots of the simulations are shown in I-h). See **Figure 1**'s caption for the descriptions of color codes for atoms. The hydrogen bonding network is shown with green dot lines.

316 We next performed GC-AIMD simulation at 0 V vs RHE. The results are summarized in Figure 317 The movie this simulation provided Supplementary 7. of is in Video 318 Movie K+CO2 0.0VvsRHE.mp4. A few snapshots of the movies are shown in Figure 7e-h. We 319 found that the CO₂ does not desorb from the surface during the simulation. The Au-C bond length is measured to be 2.13 \pm 0.06 Å (Figure 7a) and the \widehat{OCO} bent angle is measured to be 129° \pm 4° 320

321 (Figure 7b). We also observed a large variation in the distance between K^+ cation and the two 322 oxygen atoms of the CO₂ molecule. More importantly, we showed that we are able to maintain the 323 potential of the Au electrode at 0 V vs RHE with an error of 0.04 V (Figure 7c). This target 324 potential is achieved by controlling the number of electrons in the system with TPOT (Figure 7d).

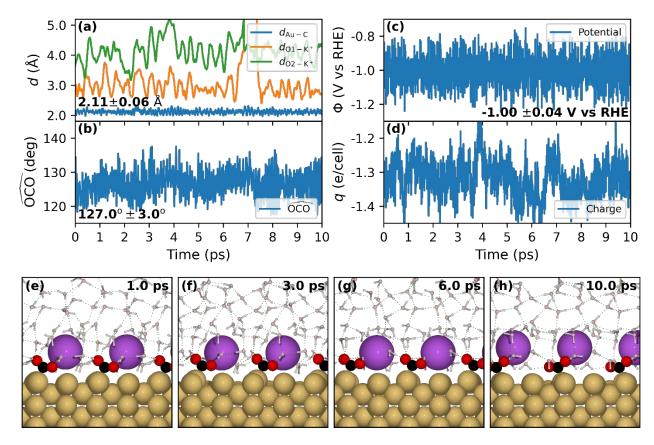


Figure 8. Evolution of (a) distance from Au and C atom of CO₂ molecule (d_{Au-C}) , distance from K⁺ to two O atoms of CO₂ molecule $(d_{O1-K^+} and d_{O2-K^+})$, (b) bending angle of the CO₂ molecule (\overline{OCO}) , (c) potential of Au(110) electrode (Φ) , and (d) charge in the supercell (q) during the GC-AIMD simulation of CO₂- $K^+/Au(111)$ with SOLHYBRID model at -1.0 V vs RHE. Snapshots of the simulations are shown iI(e-h). See Figure 1's caption for the descriptions of color codes for atoms. The hydrogen bonding network is shown with green dot lines.

332 Finally, we performed GC-AIMD simulation at -1.0 V vs RHE. The results are summarized in 333 Figure 8. The movie of this simulation is provided in Supplementary Video Movie K+CO2 -334 1.0VvsRHE.mp4. A few snapshots of the movies are shown in Figure 8e-h. We found that the 335 CO₂ does stay adsorbed on the surface during the simulation. The Au-C bond length is measured to be 2.11 ± 0.06 Å (Figure 8a) and the \widehat{OCO} bending angle is measured to be $127^{\circ} \pm 3^{\circ}$ (Figure 8b). 336 337 We also observed the variations in the distance between K⁺ cation and the two oxygen atoms of 338 the CO₂ molecule are not as large as those in the simulation at 0 V vs RHE (Figure 7a). These 339 results, i.e., shorter Au-C bond length, smaller \overline{OCO} angle, and smaller variations of distance from K⁺ 340 to oxygen atoms of CO_2 than those in 0 V vs RHE case, indicate that the CO_2 binds stronger on the Au(110) 341 electrode and that the configuration and stability of the CO₂ adsorption depends on the potential of the 342 electrode. Once again, we show that we are able to maintain the potential of the Au electrode at the 343 target potential (-1.0 V vs RHE) with a small error of 0.04 V (Figure 8c). This target potential is 344 achieved by the variation of the number of electrons in the system that is done by our TPOT 345 algorithm (Figure 8d).

346 5 CONCLUSIONS AND OUTLOOK

In this work, we have proposed to modify the popular implicit solvent model VASPSol by using a padding charge to enable its capability to perform simulations with the explicit-implicit hybrid solvent models (SOLHYBRID) and we have introduced an economical scheme to control the electrode potential (TPOT) that optimizes the number of electrons in the system during simulation to keep the electrode at a predetermined target potential. We have also demonstrated our methods by performing (GC-)AIMD for CO₂ adsorption on Au(110) in the presence of the potassium cation. Our approach not only allows the simulation of a hybrid explicit-implicit solvent but also facilitatesthe control of electrode potential with minimal error.

While we have demonstrated that the simulations (of hybrid solvent model and constant electrode potential) are feasible with the popular plane-wave DFT package, i.e., VASP, such simulations are computationally demanding mainly because of the cost associated with solving linearized Poisson-Boltzmann equation.⁵ In order to use this method for realistically large length scale and long time scale simulations, it is necessary to either improve the speed of solving linearized Poisson-Boltzmann equation⁵ or to use data generated from these methods to develop machine learning interaction potentials (MLIP) that are electrode-potential dependent.

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368 Data Availability

369 All data are included in the manuscript. Additional data supporting this work are available a zenodo 370 repo (available at the acceptance of this manuscript) including coordinates used to generate 371 movies, data for plotting.

372 Code Availability

VASP code version 5.4.4 is used for this work and is commercially available. SOLHYBRID model
is implemented in our forked version of VASPSol. The code and its manual are available at
https://github.com/zoowe/VASPsol. TPOT code and its manual are available at

- 376 <u>https://github.com/zoowe/tpot</u>. Images shown in this manuscript are generated with matplotlib,
- 377 povray. Movies are created with povray and moviepy.

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