

1 **Title**

2 Sample preparation free tissue imaging using Laser Desorption – Rapid Evaporative
3 Ionisation Mass spectrometry (LD-REIMS)

4 **Authors**

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22
23 **Abstract**

24 Laser desorption ionisation (LDI) is generally considered to be an inferior ionisation modality to
25 matrix assisted LDI (MALDI), providing information solely on lipids with low sensitivity. The

26 current study demonstrates that the combination of ambient LDI with in-source surface-induced
27 declustering provides sensitivity and chemical coverage comparable to MALDI. The setup was
28 characterised for infrared laser desorption using two different laser systems and was successfully
29 used for ambient mass spectrometric imaging. 20 μm spatial resolution was achieved with
30 oversampling, approaching single-cell resolution, while metabolites and lipids ranging from
31 amino acids through carbohydrates and nuclear bases to complex glycolipids were successfully
32 detected. The technique was also tested as a platform for MS-guided surgery, raising the
33 possibility of using a single technique for generating histological and in-vivo data. The results
34 suggest that the new method can form the basis for a new histological classification system for
35 surgery and pathology environments closing this 150 year old diagnostic gap.

36 **Teaser**

37 A new mass spectrometric imaging method bridging histopathology and in-vivo interventional
38 tissue identification methodology.

39 **Introduction**

40 Laser desorption ionisation (LDI) mass spectrometry was among the first analytical applications
41 using lasers in the late 1960s (1, 2). As one of the first techniques with the aim to analyse
42 biomolecules, laser desorption ionization (LDI) was introduced in the 1960s; also for imaging the
43 laser microprobe mass analysis (LAMMA) technique which was introduced in the 1970s (3).
44 While the rationale of this application for tissue analysis is straightforward, LDI has never gained
45 true popularity as an organic MS method due to a handful of inherent problems stemming from
46 the fundamentals of biological matter-laser interactions. These problems include (1) thermally
47 induced condensation reactions which turn biological samples into a three-dimensional, highly
48 cross-linked polymer structure (eventually leading to carbonisation) efficiently trapping analyte
49 molecules, (2) formation of a large amount of neutral molecules in the gas phase relative to the
50 formation of ions (3) aerosolization of samples leading to particles containing up to $10^8 - 10^9$

51 molecules (4–6). All these three effects act against ionisation efficiency, deeming the approach
52 insensitive for most bioanalytical applications. The first significant breakthrough regarding
53 solving these problems was the development of matrix assisted LDI (MALDI), where soluble
54 analytes are extracted into a crystalline layer of small organic molecules termed as ‘matrix’ (7).
55 The basic requirements for a matrix include high absorbance at the wavelength of the laser and
56 the formation of only gaseous degradation products on laser irradiation (8). The MALDI sample
57 preparation process transfers soluble analytes into an environment close to ideal for laser
58 desorption, giving several orders of magnitude improvement on the ionisation efficiency of small
59 molecules (metabolites, lipids) and allowing the ionisation of macromolecular species like intact
60 proteins. Nevertheless, MALDI predominantly tackles problem 1 (see above) as it still produces a
61 high amount of neutrals and clusters while introducing a set of new problems revolving mainly
62 around the deposition of the matrix layer (9). Following the development of MALDI, a number of
63 different approaches have been tested for tackling problem 2 (formation of gas-phase neutrals),
64 either in combination with matrix assistance or without it. These methods, collectively termed
65 ‘post desorption ionisation’ or shortly just ‘post-ionisation’ methods, ranged from photoionization
66 to discharge ionisation approaches used in conjunction with LDI. One of the most sophisticated
67 ones, MALDI-2, revolutionised MALDI applications in course of the last five years by further
68 improving the sensitivity of the technique by 2-3 orders of magnitude (10). We also demonstrated
69 recently that a low-temperature plasma post-ionisation results in similar sensitivity improvement
70 even in the absence of matrix (11, 12). However, none of these approaches provided a reasonable
71 solution for problem #3, the formation of large molecular clusters resistant to post-ionisation
72 techniques or the heating of certain parts of the ion path (e.g. atmospheric interface in case of AP-
73 MALDI where the cluster formation is particularly profound due to the limited expansion of
74 desorption plume under high-pressure conditions). The problem of cluster formation is not limited
75 to LDI; practically all methods developed for the ionisation of condensed phase samples suffer

76 from this phenomenon, including electrospray, the single most widely used ionisation method (13,
77 14). The idea of utilizing surface-induced dissociation to tackle this problem originated from the
78 electrospray droplet impact technique, an approach developed for enhancing the ionisation
79 efficiency of electrospray ionisation (15). Clusters accelerated by the adiabatic expansion
80 following the first gas conductance limit in an atmospheric interface can be impacted against a
81 solid target to induce their dissociation in a straightforward manner. Advantages of the approach
82 include the short timescale compared to slow heating methods and the independence of cluster
83 size compared to gas collision methods. As it is discussed below, the method works in practice
84 and fulfils the expectations outlined above.

85 Atmospheric pressure laser desorption ionisation methods – despite the poor sensitivity – have
86 gained momentum following our initial publication on the utilisation of a number of different
87 laser systems for direct tissue analysis (16). The subsequently developed Spider Mass system
88 utilising resonant IR laser at 3 μm as well as the CO₂ laser system were successfully used for bulk
89 tissue analysis with a clear perspective for direct surgical utilisation, especially in the case of
90 surgical lasers (17). Further noteworthy development was the use of picosecond infrared lasers for
91 the same purpose resulting in clear surgical advantages, including minimal thermal spread and
92 ultrafine dissection (18). The current project was driven by the motivation of developing a system
93 which is equally capable of supporting mass spectrometry-guided surgical interventions and
94 molecular histopathology. Although tissue identification has been critically important in both
95 surgery and pathology, the in-situ decision making in surgery has almost always been based on
96 sensory information (colour, physical consistency) in contrast to histopathology, where tissue
97 classification is based on cellular morphology. Consequently, the ‘macroscopic’ tumour margin in
98 the case of cancer surgery may be significantly different from the ‘microscopic’ one, potentially
99 resulting in incomplete tumour resections. Developing a technology equally capable of
100 histological imaging with a view of co-registering the resulting data with classical histopathology

101 and collecting in-situ, in-vivo data would eventually close this gap and enable surgeons to operate
102 with histological precision.

103 **Results**

104 *LD-REIMS setup development*

105 A multipurpose LD-REIMS platform was constructed to perform LD-REIMS tissue imaging
106 experiments (see **Figure 1A**). The setup consists of a laser source, a delivery system of the laser
107 beam (open beam guided with mirrors in case of OPO, optical fibre waveguide in case of CO₂)
108 and focal optics, where the final focusing is performed using a 20 mm focal distance lens. The
109 imaging setup and focal optics for the CO₂ laser are fundamentally identical. After carefully
110 optimising the system, the best focal spot sizes achieved with the OPO laser were approximately
111 60 μm and 70 μm for the CO₂ laser; improving the optical performances (by using shorter focal
112 distance lenses with suitable aberration correction, e.g. aspherical lenses) might allow further
113 reduction in focal spot sizes. At the measured spot sizes, the laser power density for both lasers
114 was around 1 kW/cm² (5 J/cm² fluence in the case of the pulsed OPO laser). Due to the pulsed
115 operation mode of the OPO laser and the maximum 20 Hz pulse frequency, the sampling
116 frequency of the mass spectrometer can only be set to a value which is dividable by 20. The
117 imaging speed was set to 10 pixels/second as a good compromise between sensitivity and speed.
118 The laser ablation/ionisation setup was coupled with a metal capillary atmospheric interface
119 custom-built for the purpose of this study. The interface (**figure 1B**) incorporated a matrix
120 introduction element, and it was prepared to accommodate the impact target used for declustering.

121 *Implementation of surface-induced declustering in the atmospheric interface:*

122 Our primary objective was to achieve sensitivity and robustness that is suitable for routine tissue
123 profiling applications. Alternatively to previous laser desorption ionisation setups, we employed a
124 heated impactor element in the vacuum interface, which enhances ion formation by facilitating
125 the declustering of the sample aerosol particles. The fundamental questions regarding this

126 instrumental setup were whether the particles accelerated by the free jet expansion would have
127 sufficient velocity to dissociate and whether the ion optics of the mass spectrometer could capture
128 ionic species formed on the impact event. Further questions were raised on the energetics of the
129 process. To simulate the impact phenomenon of the cluster on the impactor surface, a series of
130 numerical simulations were conducted to design and characterise the optimal impactor surface
131 **(Figure 2A-O)**.

132 The primary particle trajectory simulations revealed that the majority of particles entering
133 the intermediate vacuum regime of the mass spectrometer impact into the jet disruptor surface at
134 an average impact velocity of approximately 600 m/s. The specific kinetic energy of the particles
135 (100 kDa, 10.2 nm diameter) is actually very low – it is ~ 2 eV/kDa, insufficient to break covalent
136 bonds, but certainly sufficient to break up molecular clusters. The trajectory simulations for
137 secondary particles (shown on **Figure 2V**) clearly showed that even individual molecular ions of
138 complex lipids (represented by 1kDa) particles are readily entering the travelling wave ring
139 electrode ion conduit.

140 In line with the expectations, presence of a spherical collision surface in the free-jet flow region
141 enhanced the signal significantly as it is demonstrated by **Figure 2V**. The optimum distance
142 between the capillary inlet and the surface was experimentally determined (based on obtained
143 signal intensity) to be 5 mm. The results of the simulations also show that optimal ion yields can
144 be achieved with a 3 mm diameter spherical collision surface positioned 5 mm from the MS
145 transfer capillary. This geometric setup allows the free jet expansion and acceleration of the
146 sample droplets, which collide with the optimally located impactor surface for maximum
147 ionisation efficiency. The optimal distance shows a good correlation with the position of the
148 Mach-disk region of the free jet. The velocity and the kinetic energy of the individual particles is
149 the maximum before the Mach-disk, thus the positioning of the impactor surface at this region
150 explains the highest ion yields obtained by the high velocity impact at the surface. Although the

151 introduction of a solid collision target resulted in a dramatic improvement of signal intensity, the
152 time-dependent signal intensity in this case also showed a remarkable decline on the timescale of
153 minutes. We were able to recover full signal intensity by thoroughly cleaning the surface,
154 however regular cleaning after few 10 minutes of use was deemed to be incompatible with the
155 envisioned biomedical applications. In order to overcome this problem, a surface with controlled
156 temperature in the range of ambient to ~ 1400K was implemented by using a Kanthal D coil.
157 Although the coil (or band) does not provide a well-characterised surface geometry, any coating
158 (e.g. a ceramic tube) results in significant drop in the temperature of the surface at the collision
159 point due to the cooling effect of the free jet expansion. Experiments show that heating the
160 surface of the impactor to around 1200K dramatically increase signal intensity and decrease
161 memory/carry over issues. Besides providing stable signal for several hours, the introduction of
162 heated surface increased the signal intensity by further three orders of magnitude (**Supporting**
163 **Figure 1**) compared to the cold surface. The elevated temperature of the collision surface
164 enhances the declustering and provide additional thermal energy to complete the subsequent
165 desolvation process to yield individual molecular ions for analysis in the mass spectrometer. The
166 elimination of memory effects and increase in sensitivity allowed us to utilize the laser desorption
167 method for imaging experiments which require instrument stability over extended periods of
168 measurement time.

169 *Laser parameter optimisation:*

170 The laser ablation process producing primary aerosol is expected to be sensitive to the laser
171 parameters including wavelength and fluence among others. Several groups published the
172 successful application of different infrared lasers for tissue ablation and these lasers are
173 commercially available for medical tissue manipulation. However, the mechanism of the ablation
174 process remains poorly understood with the exception of MALDI applications (19–21). We have
175 tested two laser systems including a surgical CO₂ laser emitting at 10600 nm and a mid-infrared

176 Optical Parametric oscillator (OPO) (tuneable between 2700 – 3100 nm) laser, both coupled with
177 a handheld sampling device for human breast tumour tissue analysis. With the shorter pulse width
178 (5 ns vs. 100 μ s) OPO laser the optimised spectral quality (Signal-to-noise ratio and signal
179 intensity) was found to be better in the complex lipid region (600 – 1000 m/z) (**Figure 3 A-D**)
180 which has been widely used to establish histological classification (22). Univariate analysis of the
181 dataset also confirms the previous assessment, a number of significant features observed were
182 larger (286 in case of CO₂ laser and 1082 in case of OPO laser) using the OPO laser system
183 (**Figure 3 E-F**). Visualising individual significant metabolic and lipidomic species (**Figure 3 G-**
184 **K**) also confirms this, as the same features observed with both lasers show a clearer distribution
185 using the OPO laser compared to the CO₂ laser. The two main differences between these laser
186 systems are the output wavelength and pulse width. The differences observed were associated
187 with the different levels of thermal confinement, which is a phenomenon observed when laser
188 pulses interact with solid surfaces (23). Short (5 -7 nanoseconds in the current case) laser pulses
189 cause negligible thermal degradation of the biological material because most of the laser energy is
190 utilised for the cavitation-induced ablation/explosion, while in the case of longer pulse width (0.1
191 milliseconds in the current case) laser, the irradiation extends well beyond the explosion, even
192 inducing thermal degradation of molecular species. Slow thermal degradation of biological
193 macromolecules generally involves condensation reactions via loss of water and ammonia from
194 hydroxyl and amino moieties (cf. Mallard reaction) yielding a cross-linked covalent matrix
195 eventually turning into amorphous carbon (carbonisation) (24). This process effectively shifts the
196 aerosol formation from aqueous droplets towards soot. The downstream declustering process is
197 only effective for liquid droplets as the energy regime is not sufficient to break up the covalent
198 matrix of soot/carbon particles, resulting in poor sensitivity. Based on these findings, further
199 experiments with even shorter (picosecond – femtosecond) pulsed lasers are planned. Picosecond
200 laser pulse widths would allow the ablation to be fully in the stress-confined regime, increasing

201 the efficiency of ablating intact biomolecules (25). The effects of wavelength were also tested
202 with the tuneable OPO laser in the available wavelength range using pork liver samples. From the
203 data, four mass bins, 279.25, 303.25, 766.55 and 885.55 (associated with the deprotonated
204 negative ion of linoleic acid, arachidonic acid, phosphatidylethanolamine PE (38:4) and
205 phosphatidylinositol PI (38:4) respectively) were selected for the analysis of laser wavelength,
206 shown on **fig. 4 A-D**. The results show that there is an optimal wavelength between 2950 – 3000
207 nm for intensities of important molecular species, which wavelength range falls close to the
208 wavelength of 2940 nm, which is the O-H bond stretch absorption maximum wavelength. To
209 characterise the spectral pattern-level changes associated with different wavelengths, multivariate
210 statistical analysis (Principal Component Analysis - PCA) was performed on the obtained data.
211 The PCA plot (**supporting fig. 2**) shows no significant separation or differences among data
212 points except from the data obtained at 2750 nm. Looking closer at the data, the signal-to-noise
213 ratio for the 2750 nm measurement point was significantly lower than at any other wavelength,
214 resulting in the separation from the rest of the group. These results suggest that the laser
215 interaction with the tissue is limited to the desorption/mobilization due to the rapid heating of
216 endogenous water and there are no additional interactions (e.g. photochemical ionisation) between
217 the laser radiation and the sampled material. These results fall in line with results obtained by
218 other groups that also utilise endogenous water content to ablate tissues using laser desorption
219 (25, 26).

220 The effect of laser fluence was tested by using different focal optic lenses and the laser operated
221 at different power levels. Ø1 ZnSe aspheric lenses ($f = 12.7$ mm; $f = 25$ mm), Ø1/2 CaF₂ plano-
222 convex 20 mm and $f = 6$ mm) lenses were used to test the ablation characteristics at different laser
223 fluences using pork liver tissues, the tested parameters can be found in **table 1**. The best SNR was
224 achieved with the highest laser fluence using the shorter focal ZnSe lens (at 6.02 J/cm²), and a
225 characteristic ammonia loss (previously observed with REIMS technique) (24) was observed

226 within the phospholipid region in the spectra, specifically in case of the phosphatidylethanolamine
227 (PE) species. The ratio of observed $[M-NH_4]^+$ to $[M-H]^+$ gradually shifted towards a more
228 prominent ammonia loss for PE molecules, as can be seen in **figure 4 E-F**. A PCA model was
229 also generated using this dataset, where the most prominent differentiating factor in PC1 was
230 observed to be the ammonia loss detailed above. Significant separation was observed in spectra
231 obtained between fluence above and below 4.47 J/cm^2 which separation is primarily driven by the
232 effects explained above. The loading and the 3D PC models are shown in **figure 4 G-H**. Higher
233 energy ablation caused a more stable and reproducible ammonia loss effect, which improved the
234 reproducibility and robustness of the technique overall. Based on these results the minimal
235 ablation threshold for LD-REIMS reproducible data was determined to be at 4.50 J/cm^2 , and the
236 laser fluence of 5 J/cm^2 was chosen for the imaging experiments below.

237 *Imaging mass spectrometry*

238 Mouse brain mass spectral images were acquired using both lasers at $70 \mu\text{m}$ raster size, the results
239 of the imaging are shown on **supporting figure 3** (for the OPO laser) and **supporting figure 4**
240 (for the CO_2 laser). The data reveals that both methods are capable of acquiring information from
241 spatially heterogeneous samples, however, the imaging data obtained with the OPO laser had
242 better sensitivity and SNR ratios, which was expected due to the previously mentioned physical
243 differences between the different lasers. The spectral profile mainly consists of small molecules
244 (metabolites, fatty acids and phospholipids), which falls in line with results obtained during
245 previous iKnife experiments. A list of tentatively identified molecules is given in **table 2**. The
246 annotation was performed using accurate mass analysis and using the Chemical Abstract Services
247 database. The set of observed molecular species show good overlap with species described by
248 previous literature on applied laser desorption or REIMS-based MS techniques (27).

249 To challenge the resolution limit of the experiment, the concept of oversampling imaging was
250 tested using the OPO laser with the LD-REIMS imaging setup. At current spot sizes ($60 \mu\text{m}$)

251 different imaging raster sizes were tested on coronal mouse brain tissue sections to determine the
252 maximum achievable resolution. The improvement of the spatial resolution of MSI techniques is
253 one of the main objectives of the field. To achieve this, several different approaches have been
254 described in the literature. Oversampling using the analytical beam has been proposed and
255 successfully demonstrated using MALDI technique (28, 29). Other approaches use post-ionisation
256 steps like MALDI-2 after sampling to achieve sub-micron imaging resolution (10). Using the
257 oversampling approach with the OPO laser setup, different raster size images were acquired at 70,
258 50, 30 and 20 μm pixel size (see **supporting figure 2**). Using the molecular ion of adenine ($[\text{M}-$
259 $\text{H}]^-$ 134.034 m/z) the fine structure of the hippocampal region was identified on the images. As
260 the pixel size decreases, the details of the dentate gyrus region become better defined, and the
261 contrast between the polymorph and the granule cell layer increases. At 30 μm or smaller raster
262 sizes, the fine structures of the stratum radiatum and stratum lacunosum-moleculare regions in the
263 mouse brain become observable as well. The currently achieved 20 μm resolution limit makes the
264 LD-REIMS technique a competitive technique with other MSI techniques. With further planned
265 optimisation, the current imaging setup could acquire images with single-cell resolution. The
266 higher-resolution images provide superior imaging data; however, this comes at the price of
267 significantly increased analysis time. The imaging experiment of a coronal mouse brain section at
268 70 μm took 40 minutes using 10 pixel/second acquisition rate, and a similar-sized tissue section
269 analysed at 20 μm using the same speed requires more than six hours.

270 *Clinical application of mass spectral imaging*

271 A clinical case study was performed with both laser setups using cancerous human breast tissue
272 samples. Human breast tumour samples were chosen to demonstrate the molecular pathology
273 capabilities of the system. 12 μm thick breast tumour sections were mounted on standard glass
274 slides and imaged using the prototype imaging setup with both lasers without any additional
275 sample preparation steps. The data obtained from the imaging experiments were compared with

276 the gold standard Hematoxylin-Eosin staining method, the results are shown on **Figure 5**
277 evaluated using hyperspectral correlation methods (30). The characteristic differences in
278 molecular profiles found in the imaging data show good correlation with the histologically
279 separated tissue regions that differentiate cancerous and healthy tissue regions. Comparing data
280 from the different lasers, the OPO laser operated at 2940 nm was found to be superior, which
281 observation falls in line with expectations.

282 The main detected class of molecules are glycerophospholipids and triglycerides (mainly cell
283 membrane compounds), additionally several fatty acids, small molecules and metabolites were
284 observed as well, which may provide valuable insight on clinically important questions, such as
285 differentiating between breast fibroadenoma and cancer tissues.

286 LD-REIMS is a powerful technology for sample preparation free, untargeted tissue imaging,
287 suitable for label free molecular histopathology, and has the potential to be an efficient tool in
288 pathology and drug discovery laboratories around the world. Beyond these advantages the
289 technology also bridges imaging MS with in-vivo MS. The REIMS based intelligent surgical
290 device (iKnife) technology has been demonstrated for *in vivo* tissue characterisation, however the
291 technique is strongly dependent on the underlying classification models (31–33). To perform any
292 tissue classification, a histologically validated database is required to provide adequate references
293 for accurate diagnosis. The required data is usually produced by analyzing excised tissue
294 specimens using the surgical energy device, where bulk tissue is ablated and the resulting aerosol
295 is analyzed by means of REIMS. The remaining tissue is sent for histological assessment to
296 determine pathological status of the ablated cells. This process however carries a certain degree of
297 uncertainty since the ablation is performed using a handheld probe, which causes significant
298 variance in the ablated tissue volume that can lead to sampling a mixture of cells belonging to
299 different histological classes. The pathologist also provides an assessment not based on the
300 ablated tissue, but on the margin remaining after the ablation process which brings another level

301 of uncertainty to the database, as the exact histological composition of the ablated tissues cannot
302 be properly determined. Since LD-REIMS imaging and the iKnife uses the same mass
303 spectrometric experimental setup, translating the data from tissue sections to bulk *in vivo* tissue
304 analysis should be feasible. As the laser can give better spatial resolution than the diathermy and
305 therefore a lower chance of mixed tissue signal, the laser data is expected to carry higher
306 histological specificity. If a tissue section can be imaged with LD-REIMS not only are the
307 resolution and specificity increased but also the number of data points. Most importantly, the
308 histological analysis of consecutive sections gives proper classification data down to single cell
309 accuracy. Thus, LD-REIMS imaging is expected to solve all problems associated with the
310 production of training data for iKnife applications.

311 A proof-of-concept study was conducted where a database and a multivariate model was
312 constructed using the imaging CO₂ setup and was validated by the iKnife instrument (**figure 6 A-**
313 **B**). For laser sampling, surgical CO₂ laser was chosen as this laser allowed us to collect surgical
314 iKnife data as well as MSI data. The combined method model showed good sensitivity (92.3%)
315 and specificity (90.6%). Good separation was observed between healthy and cancerous tissue
316 based on the first principal component. Cancerous laser and iKnife data points overlapped on the
317 PCA plot; however, minor separation could be observed between laser and iKnife normal tissue
318 spectra on the second component probably due to the improved signal-to-noise ratio with the
319 laser. As the first component represents the highest variance, and cancerous and healthy data were
320 separated along the PC1 axis regardless of whether it was acquired with the laser or the iKnife,
321 the models were exchanged to test classification accuracy (**Figure 6 C-D**). Diathermy data
322 evaluated with the laser model had 100% correct classification for normal tissue and 92% for
323 cancerous tissue. The laser data evaluated with the diathermy model gave 97% correct
324 classification for normal tissue and 100% for cancerous tissue. These results show that these
325 models can be used across different ablation modalities, raising the possibility of creating

326 method-independent models universally applicable across all surgical energy devices and
327 corresponding imaging modalities.

328

329 **Discussion**

330 It was demonstrated that the ion yield of atmospheric pressure desorption ionisation methods can
331 be significantly improved by the low energy surface-induced dissociation of molecular clusters in
332 the atmospheric interface region of the mass spectrometer. The simulation results show that if
333 singly charged 1kDa particles with a diameter of 2.2nm are generated on the surface of a spherical
334 collision surface positioned in the Mach disk region, the ion optics of a commercial mass
335 spectrometer can efficiently capture these secondary ions.

336 Additionally, heating the impactor surface was found to yield a robust analytical setup with
337 minimal carryover or surface fouling and further increased the overall ion yield of the desorption
338 ionization technique. This latter effect was associated with further energy deposition into the
339 system, facilitating the complete dissociation of supramolecular adducts. While one of the
340 potential drawbacks of heating the impactor surface is the thermal degradation of biomolecules,
341 as the surface reaches temperatures up to 1400K, no such effect was observed. This effect was
342 associated with the adiabatic expansion of the gas jet after the first conductance limit, which
343 exerts an efficient cooling effect on the transported particles, lowering the internal energy to the
344 effective temperature range of 10-40K. Furthermore, the timeframe of the particle-surface
345 interaction is relatively short, as individual molecules spend less than 10 μ s in the vicinity of the
346 collision surface. Thermal conductivity and thermal convection are negligible due to the vacuum
347 environment around the impactor surface, and heat transfer through infrared radiation is
348 insufficient to deliver enough energy in this timeframe to thermally degrade ions of interest. The
349 described impact-based ionization enhancement setup intrinsically accelerates clusters due to the
350 forementioned adiabatic expansion effect, resulting in a simple, robust and compact embodiment

351 not requiring any electrostatic field gradient. The kinetic energy of the high velocity clusters is
352 utilized in the generation of individual molecular ions as the clusters break up resulting in a
353 relatively low kinetic energy ion population, which is efficiently captured by a ring electrode or
354 multipole ion guide.

355 The phenomena of surface – cluster interactions and the effect of it on ion formation have been
356 less studied compared to other ion formation mechanisms. However, both electrospray ionisation
357 and laser desorption-based methods generate clusters with sufficient kinetic energy to reach and
358 collide with parts of the ion optics in spite of the electrostatic or pseudopotential fields utilized for
359 ion guidance. These collisions likely result in additional uncontrolled ionization phenomena,
360 materializing as memory effects for the user (34). In certain commercial atmospheric interface
361 setups these cluster-surface collisions may significantly contribute to the observed signal.

362 We have implemented atmospheric interface SID in conjunction with laser desorption ionization
363 to gain sensitivity for mass spectrometric tissue analysis. The utilization of infrared laser
364 desorption has been an obvious choice for a method generating primarily aerosols when used for
365 tissue ablation. While diathermal ablation involves macroscopic thermal damage including
366 charring, UV laser desorption results in the excessive fragmentation of gaseous molecules and
367 ions. In contrast, infrared lasers (especially short pulse width resonant mid-infrared lasers)
368 transfer the molecular content of tissues into the gas phase efficiently and without significant
369 thermal damage or fragmentation (35, 36). These properties make the technique one of the most
370 promising candidates for a universal tissue profiling method equally capable of tissue imaging as
371 well as bulk or even in-vivo tissue analysis. The most serious shortcoming of LDI methods (in
372 contrast to e.g. MALDI) has traditionally been the low ion yield, which is less critical for bulk
373 analysis where milligrams of tissue can be ablated from tens of thousands of μm^2 area (17, 24).
374 However, histological imaging requires significantly better sensitivity, hitherto inaccessible for
375 routine applications. The importance of universal applicability is particularly important in the

376 aspect of generating histologically annotated imaging data to train shallow or deep learning-
377 driven tissue classification methods as it is described above. Furthermore, directly linking MS-
378 guided surgery techniques with MS-based histology methods provides a unique selling point for
379 the technique in both environments. Current histopathology tools cannot be used for real-time,
380 interventional tissue classification or detection and similarly, current intraoperative tissue
381 identification tools ranging from fluorescent labelling to impedance spectroscopy do not provide
382 sufficient information for histological assessment of tissue specimens.

383 The presented results serve as a basis for future work regarding sample preparation-free molecular
384 imaging of clinically important tissues and other samples. Since the technique requires no
385 labelling, sample preparation or other user interaction with the process, the method has a high
386 potential for non-targeted, automated molecular profiling of human disease samples.

388 **Materials and Methods**

389 *Experimental Design:*

390 A Waters Xevo G2-XS QToF (Waters, Wilmslow, UK) mass spectrometer was used for
391 all experiments. The instrument was equipped with a modified REIMS source described by Balog
392 et. al.(31). 100 µl/min MS grade 2-propanol (Merck, Gillingham, UK) was injected in front of the
393 MS inlet capillary to achieve Matrix Assisted REIMS as described by Jones et. al.(37). A 1.5 m
394 long, 1.6 mm I.D. PTFE tube (Merck, Gillingham, UK) was used to aspirate the aerosol from the
395 sampling position. For the iKnife diathermy experiments, a ForceTriad (Medtronic, Watford, UK)
396 electrosurgical generator was used in monopolar mode. The electrosurgical unit was coupled with
397 a modified electrosurgical handpiece (Waters Research Center, Budapest, Hungary), the output
398 power was 20W. Pork liver tissue were obtained from commercial suppliers. Ethical approval was
399 gained from the South East London Research Ethics Committee Reference 11/LO/0686, the East
400 of England - Cambridge East Research Ethics Committee Reference 14/EE/0024 and the project

401 was registered under the Imperial College Tissue Bank. Data were only obtained from patients
402 who had consented to the use of tissue for research.

403 *Numerical simulation:*

404 Characterization of the flow structure was performed sequentially using different solvers
405 and tools to cover the wide pressure range of the flow field and perform particle tracing
406 respectively. The first step involved the application of the Reynolds-Averaged Navier-Stokes
407 method to solve the high-pressure region inside a 70 mm, 0.5 mm I.D. inlet capillary at room
408 temperature. A structured high-density mesh was generated to describe the laminar flow across
409 the capillary. Zero slip velocity was imposed on the inner walls of the capillary. Mass averaged
410 velocity, temperature and gas density values were determined at 0.2 mm from the capillary exit as
411 335 m/s, 240 K and 243 mbar and subsequently imposed as inlet boundary conditions for the
412 Direct Simulation Monte Carlo (DSMC) calculations. The DSMC domain was constructed with a
413 5 μm cell size progressively increased to 10 μm in regions distant from the impingement target to
414 increase computational speed without affecting the accuracy of the solution near the target. The
415 background pressure was 3 mbar. DSMC calculations were performed in the low-pressure region
416 of the domain using the SPARTA solver. The flow field solutions were then further processed in
417 MATLAB and introduced into SIMION for tracing particles. A thousand 100 kDa neutral
418 particles with 10.2 nm diameter were used to simulate solvent clusters. Their position and impact
419 velocity were recorded and then used to generate an ion cloud of 1 kDa particles with 2.2 nm
420 diameter.

421 *Laser parameter characterisation and imaging*

422 A commercially available Opolette HE2731 Optical Parametric Oscillator (Opotek,
423 Carlsbad, USA) and a FELS25A Intelliguide CO₂ laser (Omniguide, Cambridge, USA) laser were
424 used during the experiments. Optomechanical components were obtained from Thorlabs for the
425 optical cage system construction (Thorlabs, Newton, USA). LA7733-E4 and LA7477-E4 ZnSe

426 and LA5315-E CaF₂ lenses were used for the OPO laser and only LA7733-E3 ZnSe was used for
427 the CO₂ laser due to lack of appropriate anti-reflection coating. To test the effect of wavelength
428 on the tissue desorption efficiency, 12 μm fresh frozen pork liver slides were sampled using the
429 tuneable OPO laser between 2700 – 3100 nm and the laser power was normed to 5 J/cm². The
430 laser energy was measured using an EnergyMax-RS J-10MB-HE Energy Sensor (Coherent,
431 Saxonburg, USA) after the last focal lens on the defocused beam (around 10 mm from the focal
432 spot). The fluences were calculated by measuring the area of the ablated spot sizes and the used
433 energy for ablation using optical microscopy. The imaging experiments were done using a
434 modified two-dimensional stage setup (Prosolia). For imaging data processing and visualization
435 of the ion heatmap images, HDImaging (version 1.4, Waters) software was used.

436 *Statistical Analysis:*

437 Data processing for modelling was performed using Abstract Model Builder (AMX, version
438 1.1967.0, Waters). This software was used to select was used to define the spectra used for data
439 analysis. Mass drift correction was performed against the leucine enkephalin lock mass compound
440 (negative mode $m/z = 554.2615$), and mass binning was done to 0.1 Da. For univariate analysis,
441 an in-house data processing pipeline was used that was written in python. Principal Component
442 Analysis (PCA) was also performed to evaluate spectral differences using the AMX software.
443 Principal Component Analysis (PCA) and PCA- Linear Discriminant Analysis (PCA-LDA)
444 models were built and the PCA-LDA model was used for tissue classification. Spectral
445 comparison for the translational study was performed with multivariate statistical approaches.
446 Both diathermy and laser data were plotted on the same PCA-LDA model, which model was
447 cross validated using leave – one – out cross validation method.

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576 Numerical simulations: DP

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578 Sample collection, histology: HH, AM

579 Supervision: IF, JB, ZT

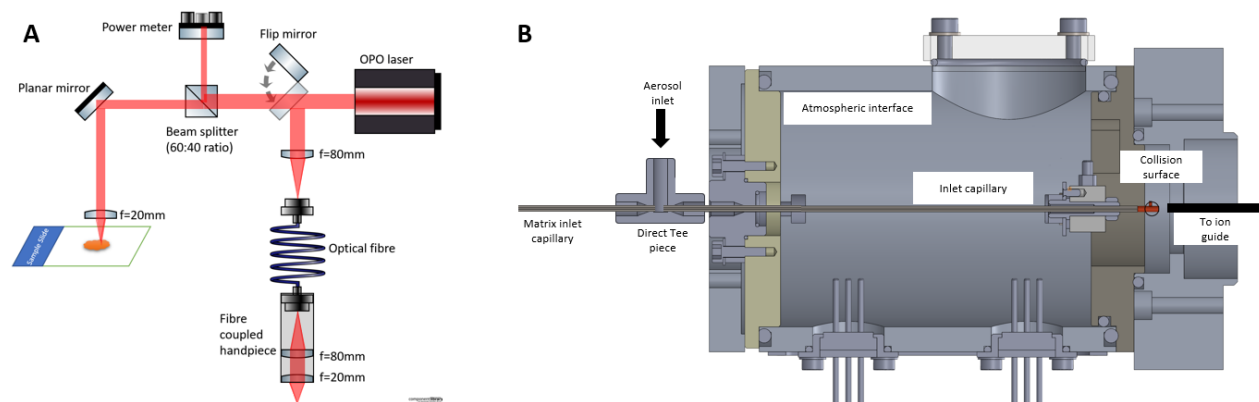
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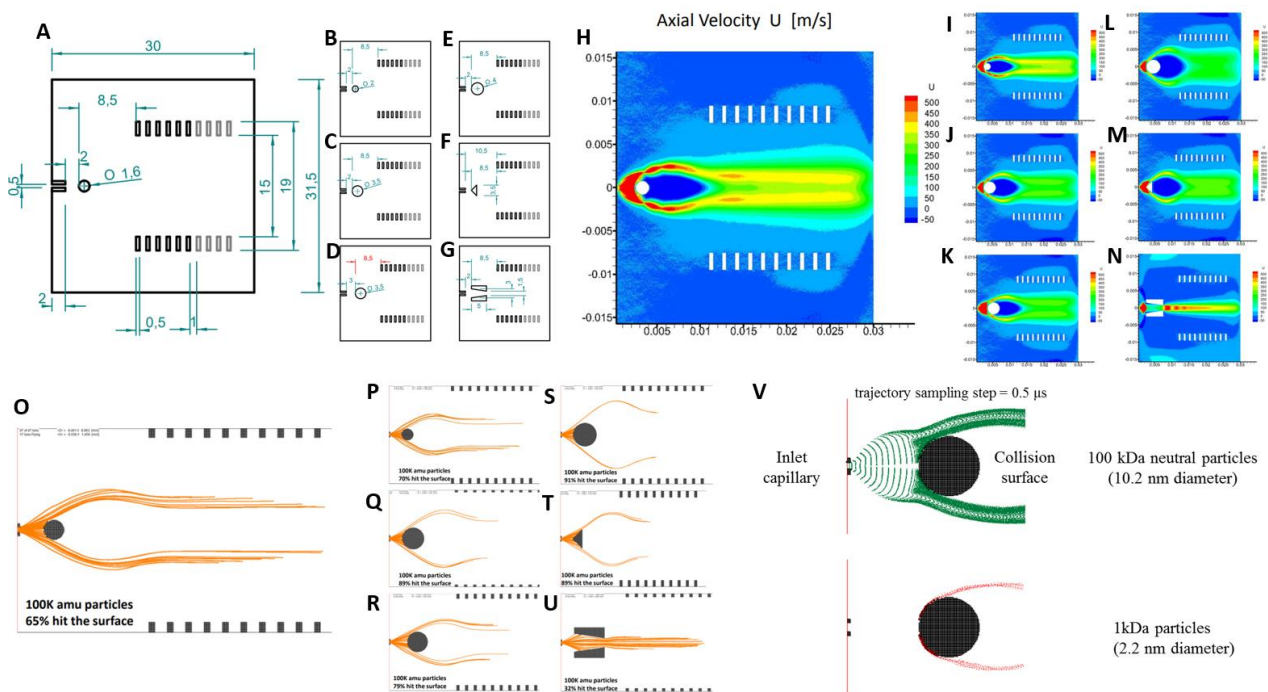
586 Figures and Tables



587
588 **Fig. 1.: Schematics of the LD-REIMS sampling and ionisation setup.**

589 The schematics of the prototype LD-REIMS imaging platform can be seen on figure A. The laser
590 source (in this case the OPO laser) is emitting the beam, which can be directed using a flip mirror
591 towards an optical fibre-coupled handpiece that allows the free analysis of bulk tissue samples
592 without the spatial information. The second option if the mirror is not used is the high-resolution
593 imaging setup, where the beam is first split with a power meter (for continuous power
594 monitoring) then directed and focused onto the glass slide holding the sample, which is positioned
595 on a 2D XY motorized stage. The generated aerosol is evacuated and introduced to the
596 atmospheric interface (B) through the Direct Tee piece. A continuous matrix solvent is added

597 through the matrix inlet capillary, and the droplet-fused aerosols enter the vacuum through the
 598 inlet capillary. After exiting the inlet capillary, the droplets go through an acceleration via the
 599 carrier gas' adiabatic expansion, and hit the high temperature collision surface, where free gas
 600 phase ions are liberated from the aerosol droplets, which can be collected by the ion guide and
 601 analyzed by the mass spectrometer downstream.

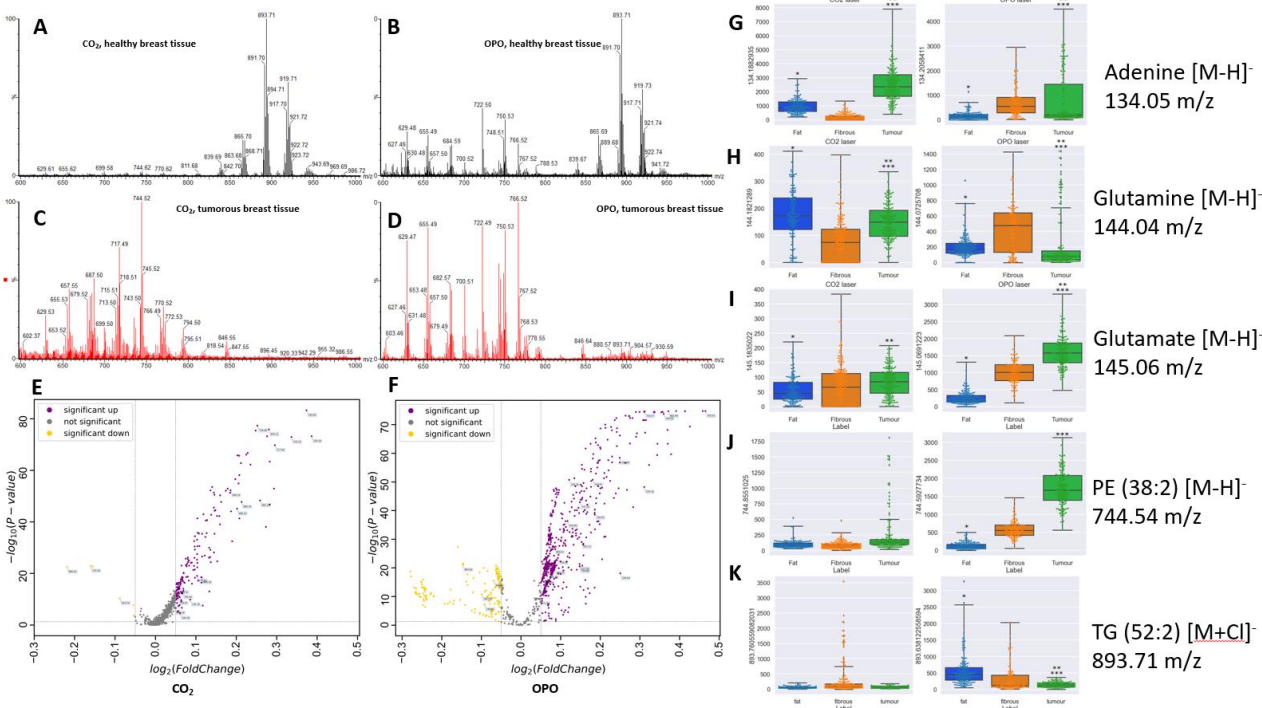


602

603 **Fig. 2.: Optimisation of the impactor surface using DSMC numerical simulations.**

604 The prototype atmospheric interface design allowed the rapid and efficient experimental
 605 characterisation of different collision surface geometries in addition to the numerical simulation
 606 data. DSMC simulations were performed in the modelled atmospheric interface environment
 607 using the following geometries: 1.6 mm sphere, 2 mm distance from inlet capillary (A), 2 mm
 608 sphere, 2 mm distance from the inlet (B), 3.5 mm sphere, 2 mm from inlet (C), 3.5 mm sphere, 3
 609 mm from inlet capillary (D), 4 mm sphere, 2 mm from the inlet (E), 3.5 mm cone, 2 mm from
 610 inlet capillary (F), 3-1.5 mm decreasing funnel, 5 mm long, 2 mm from inlet capillary (G). The
 611 axial velocity profiles (H-N) were characterised, and 100k amu particle collision simulations (O-
 612 U) were performed. Simulation results for the interaction of particles with a spherical collision

613 surface, derived from DMSC calculations followed by particle tracing were also conducted (V).



614
615 **Fig. 3.: Lipidomic and metabolomic characterisation of breast cancer samples using LD-**
616 **REIMS.**

617 Ex vivo human breast samples were analysed using two different IR lasers (OPO and CO₂) using
618 the LD-REIMS method. The tissues contained healthy and cancerous sections, validated by
619 histopathology analysis. The spectra observed in the lipid region (m/z 600 – 1000) show
620 characteristic differences of elevated levels of triglycerides in the healthy breast tissue (A-B) and
621 phospholipids are more prevalent in the tumorous samples (C-D). Volcano plots generated from
622 the data obtained with the two lasers (E-F) allow us to identify numerous metabolites and
623 phospholipids that show statistically significant fold changes. Univariate plots of identified
624 molecules of Adenine [M-H]⁻ (G), Glutamine [M-H]⁻ (H), Glutamate [M-H]⁻ (I),
625 Phosphatidylethanolamine PE(38:2) [M-H]⁻ (J) and Triglyceride TG(52:2) [M+Cl]⁻ (K) show
626 significant differences between different tissue types. * Statistically significant difference
627 between fat and fibrous tissue ** Statistically significant difference between tumour and fat tissue
628 *** Statistically significant difference between tumour and fibrous tissue.

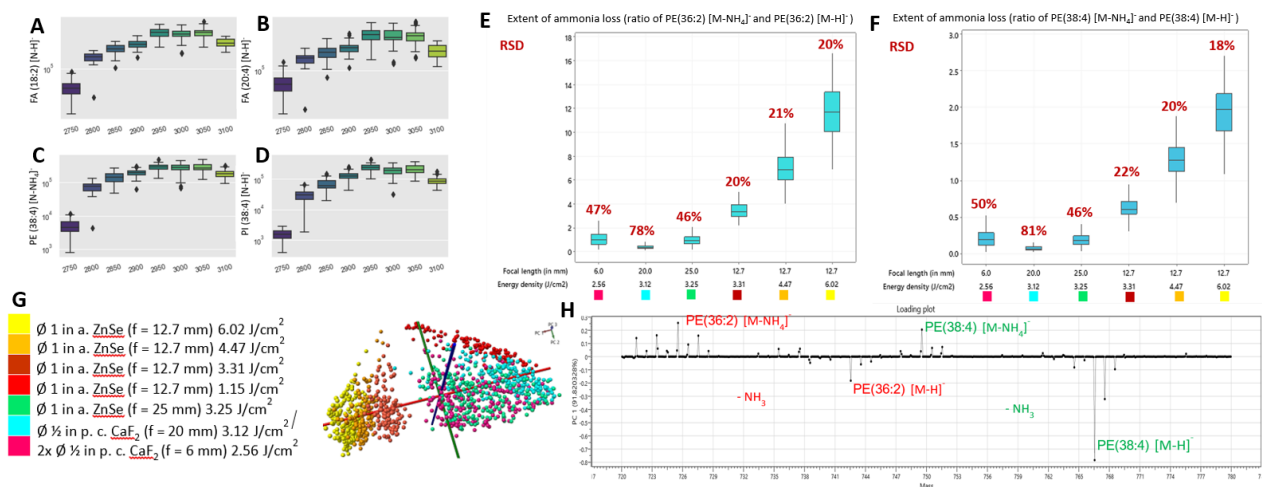


Fig. 4.: Effects of laser fluence on the spectral composition and molecule fragmentation.

The tunable OPO laser allowed the characterisation of the effect of wavelength on the desorption process on pork liver tissue. The available wavelength range (2700 – 3100 nm) was scanned in 50nm resolution; the laser energy was normed to the same fluence. Spectral profiles were recorded in the range of 2750 – 3100, 2700 nm provided no observable spectra due to low laser absorption rates. The intensity profile of four known molecules over the tested range (FA 18:1, FA 20:4, PE 38:4 and PI 38:4) (A-D) show an optimal desorption wavelength at around 2900 – 3000 nm. This falls in line with the water absorption maxima observed at 2940 nm in the mid-IR range originating from the O – H bond stretch absorption. The loss of ammonia from phosphatidylethanolamine lipids can be observed with the LD REIMS process. The ratio of observed ions of two PE species (PE(36:2) and PE(38:4)) originating from pork liver samples is shown to increase as the laser fluence increases (E-F). The PCA model can distinguish between spectra acquired at different fluence levels, as seen in plots G and H. The deamination process of PE species becomes prevalent with laser fluence levels above 3.3 J/cm².

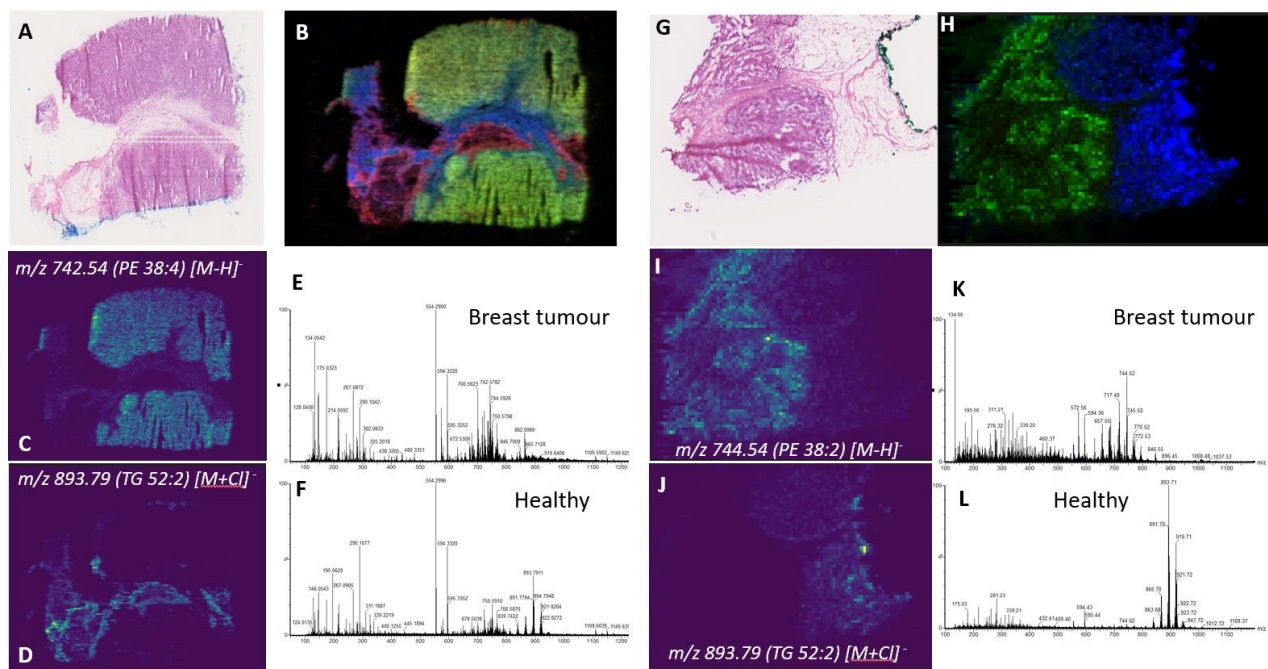
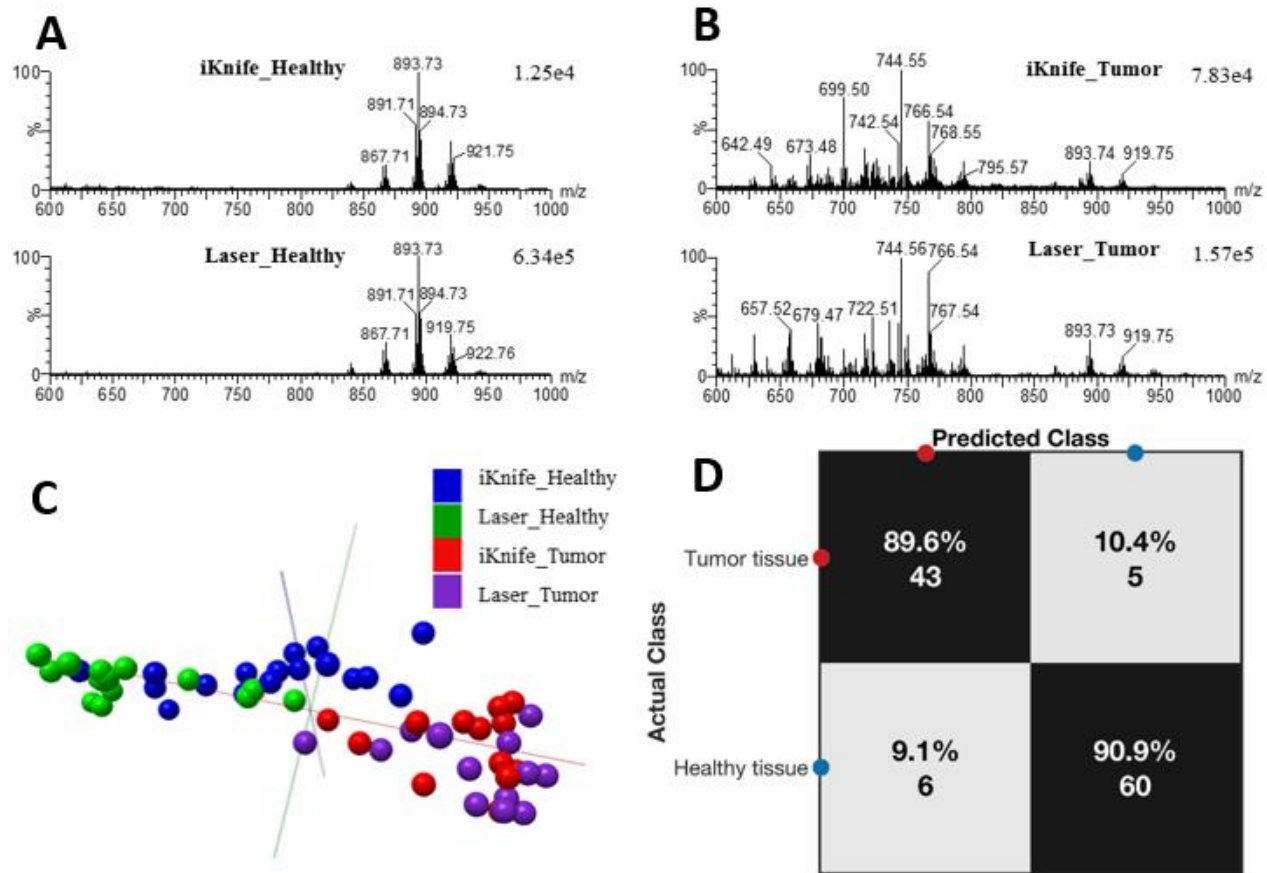


Fig. 5.: Biological imaging study of two breast cancer samples using OPO and CO₂ lasers.

The breast cancer samples were analysed using the two lasers at 70 μ m raster sizes. Consecutive tissue sections were H&E stained and classified by pathologists (A and G). The results of the OPO imaging can be seen in figure B-E, where the composite RGB image (B) is visualised (Red: 893.79 m/z (TG 52:2 [M+Cl]⁻), Green: 742.54 m/z (PE 38:4 [M-H]⁻), Blue: 281.23 m/z (FA18:1 [M-H]⁻)). The results of the CO₂ laser can be seen in figures G-K, where the composite image (H) is visualised (Blue: 893.79 m/z (TG 52:2 [M+Cl]⁻), Green: 742.54 m/z (PE 38:2 [M-H]⁻)). Individual ion images for different ionic species (C-D for OPO and I-J for CO₂) show good differentiation between different histological status tissues, and the 50 – 1200 m/z spectra of different annotated regions (E-F for OPO and K-L for CO₂) show clear differences between the metabolic and lipidomic fingerprint of healthy and cancerous tissues.



656

657 **Fig. 6.: Translational capabilities between intraoperative iKnife technique and digital**
 658 **pathology LD-REIMS experiments.**

659 Spectral comparison between different breast tissue types obtained with the laser and iKnife
 660 methods was performed using breast cancer samples (A-B). A PCA model was built from the
 661 obtained data for tissue classification (B). Leave-one-out cross-validation was performed from the
 662 previously built PCA-LDA model; the model yielded good sensitivity (92.3%) and specificity
 663 (90.6%) (C-D).