

## Hydrophobic nanoconfinement enhances CO<sub>2</sub> conversion to H<sub>2</sub>CO<sub>3</sub>

Nabankur Dasgupta,<sup>1</sup> Tuan A. Ho,<sup>1\*</sup> Susan Rempe,<sup>2</sup> and Yifeng Wang<sup>3</sup>

<sup>1</sup>*Geochemistry Department, Sandia National Laboratories, Albuquerque, NM 87185, USA*

<sup>2</sup>*Center for Integrated Nanotechnologies, Sandia National Laboratories, Albuquerque, NM 87185, USA*

<sup>3</sup>*Nuclear Waste Disposal Research and Analysis Department, Sandia National Laboratories, Albuquerque, NM 87185, USA*

\* Corresponding author: taho@sandia.gov

### Abstract

An understanding of the formation of H<sub>2</sub>CO<sub>3</sub> in water from carbon dioxide is important in many environmental, biological, and industrial processes, and in the global carbon cycle. Although numerous computational and experimental investigations have focused on understanding these interactions, the conversion of CO<sub>2</sub> to H<sub>2</sub>CO<sub>3</sub> in nanopores, and how this conversion differs from that in bulk water, has not been understood. In this study, we use ReaxFF metadynamics molecular simulations to demonstrate striking differences in the free energy of CO<sub>2</sub> conversion to H<sub>2</sub>CO<sub>3</sub> in bulk water compared with that in water confined in pyrophyllite nanopores. We find that the nanoconfinement not only reduces the energy barrier, but also reverses the reaction from endothermic in bulk water to exothermic in nanoconfined water. In addition, charged species are observed more often under nanoconfinement than in bulk water. The higher number of reactive encounters, stronger solvation, and more favorable proton transfer with increasing confinement enhance the thermodynamics and kinetics of the reaction. As carbonation in nanopores is important in the carbon cycle and a complicated problem that depends on confinement, surface chemistry, pore chemistry, and concentration of CO<sub>2</sub>, our results provide a mechanistic understanding to an important step in this process.

### Introduction

Carbon dioxide (CO<sub>2</sub>) is a greenhouse gas that makes a major contribution to global warming and climate change.<sup>[1]</sup> To reduce CO<sub>2</sub> emissions and mitigate its adverse effects on climate change, several CO<sub>2</sub> management technologies have been developed in recent years, including carbon capture, conversion, and sequestration (CCCS).<sup>[2-4]</sup> Many of these technologies involve the interaction of clay with CO<sub>2</sub>. For example, in geological sequestration, numerous experimental

investigations<sup>[5,6]</sup> concluded that clay adsorbs a significant amount of CO<sub>2</sub>. Natural clay minerals are also potential materials<sup>[7,8]</sup> for capturing CO<sub>2</sub> due to their large surface area, high porosity, abundant basic sites, excellent thermal and chemical stability, and low cost.<sup>[9]</sup> Expansive clay minerals, such as montmorillonite, can adsorb a significant amount of water into the interlayers. Therefore, investigation of CO<sub>2</sub> interactions with clay are incomplete without considering the interactions of CO<sub>2</sub> with interlayer water, including the dissolution of CO<sub>2</sub> to form H<sub>2</sub>CO<sub>3</sub>, which then undergoes dissociation to form bicarbonate (HCO<sub>3</sub><sup>-</sup>) and carbonate (CO<sub>3</sub><sup>2-</sup>) ions, as given by Eq. 1,<sup>[10]</sup>



An understanding of the formation of H<sub>2</sub>CO<sub>3</sub> in water from carbon dioxide, and its decomposition, is important in many environmental, biological, and industrial processes, and in the global carbon cycle.<sup>[11–15]</sup> Several studies have examined the formation and dissociation of H<sub>2</sub>CO<sub>3</sub> in bulk water by quantum mechanical methods, including density functional theory (DFT), Hartree–Fock (HF), perturbation theories, and enhanced sampling methods.<sup>[16] [17]</sup> The reaction,



specifically the forward reaction for  $n=1-4$  at 0 K, encountered significant energy barriers. For example, when  $n = 1$ , the energy barrier is about 50 kcal/mol.<sup>[18]</sup> Increasing  $n$  decreases the energy barrier (at 0 K) to ~33 ( $n=2$ ), ~29 ( $n=3$ ), and ~20 kcal/mol ( $n=4$ ). Independent of  $n$ , the energy of H<sub>2</sub>CO<sub>3</sub>-water complexes is about 6-7 kcal/mol higher than that of the CO<sub>2</sub>-water complexes,<sup>[18]</sup> leading to a more facile reverse Eq. 2 reaction (H<sub>2</sub>CO<sub>3</sub> +  $n$ H<sub>2</sub>O → CO<sub>2</sub> and ( $n+1$ )H<sub>2</sub>O). Similar to the H<sub>2</sub>CO<sub>3</sub> formation reaction from CO<sub>2</sub>, adding more water to H<sub>2</sub>CO<sub>3</sub> reduces the energy barrier of decomposition significantly,<sup>[19–21]</sup> from ~44 ( $n=0$ ) to ~27.1 ( $n=1$ ) and ~20.9 ( $n=3$ ) kcal/mol.<sup>[21]</sup> Other studies<sup>[22–24]</sup> also found that energy barriers for both forward and reverse hydration reactions decrease to be within the range of 16~20 kcal/mol for  $n>3$ , which is in good agreement with experimental values.<sup>[25–28]</sup>

Even though the thermodynamic and kinetic data for the formation, from CO<sub>2</sub>, and dissociation of H<sub>2</sub>CO<sub>3</sub> in bulk water are well-known, similar data in nanoconfined water is scarce. It is well-accepted that nanoconfinement can affect properties of water, such as the equation of state,<sup>[29,30]</sup> phase behavior,<sup>[31,32]</sup> adsorption,<sup>[33,34]</sup> dielectric constant,<sup>[35–37]</sup> and diffusion constant.<sup>[38,39]</sup> As a result, the reactivity of solutes under nanoconfinement may differ from that in bulk solutions.<sup>[40–43]</sup> For example, the dimensional reduction and increased fluid density could enhance reactions between small solutes in nanoconfinement,<sup>[44,45]</sup> whereas reactions involving

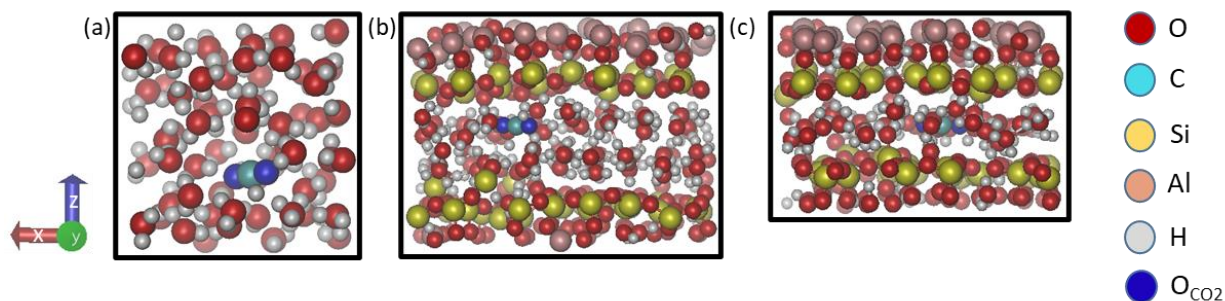
large reactants may be sterically hindered.<sup>[40]</sup> Furthermore, the changes in the parallel and perpendicular component of dielectric tensors of confined water enhance self-dissociation of water even though the quantitative relationship between self-dissociation of water with changes in dielectric tensor remains to be established.<sup>[36]</sup> Therefore, it is important to elucidate the reactivity of CO<sub>2</sub> under nanoconfinement, which can differ substantially from that in bulk water.

Interlayers of the naturally occurring clay minerals provide a confined environment to augment chemical reaction rates, increase selectivity of species, and stabilize reactive species.<sup>[46–48]</sup> The clay structural charges and interlayer cations promote a range of physical and chemical processes, including intercalation, swelling, adsorption, and reaction. Experimental<sup>[8,49,50]</sup> and theoretical (ab initio and classical molecular simulations)<sup>[51–56]</sup> studies on the interactions among CO<sub>2</sub>, water, and clay have focused mainly on adsorption, intercalation, and swelling of clay interlayers. Reactions in Eq. 1 and 2 are poorly investigated in clay nanopores. Limited available data provided a conflicting conclusion about the nanoconfinement effects on CO<sub>2</sub> conversion. For example, Loring et al.<sup>[57]</sup> and Krukowski et al.<sup>[58]</sup> found no evidence of carbonate-forming reactions in montmorillonite, while Hur et al.<sup>[59]</sup> observed the opposite. Only a few studies<sup>[60–62]</sup> explored the reactivity of CO<sub>2</sub> under different clay interlayers. Those studies reported that the carbonation reaction readily occurs near the electron-rich terminal oxygen sites adjacent to cation vacancies of the surface. Nevertheless, our understanding of the formation of carbonic acid within clay interlayers is incomplete. Further investigation of carbonic acid formation under nanoconfinement, as pursued here, will help advance our understanding and enable exploitation of CO<sub>2</sub>-clay interactions for CCCS.

In this work, we apply an enhanced sampling method, metadynamics, in molecular dynamics simulations using the reactive force field, ReaxFF,<sup>[63,64]</sup> to investigate the thermodynamics and kinetics of H<sub>2</sub>CO<sub>3</sub> formation under nanoconfinement. The ReaxFF molecular simulations allow molecular bonds to be broken and reformed without using expensive quantum mechanical methods. ReaxFF force fields are parameterized according to DFT data and are capable of accurately defining bond formation and bond breaking in large systems (> 10<sup>6</sup> atoms).<sup>[65–69]</sup> When combined with an enhanced sampling method, ReaxFF molecular simulations can sample high energy states of a relatively large chemically reactive system that would not be possible using ab initio or non-reactive molecular simulations.<sup>[70]</sup> Therefore, for the present study, ReaxFF metadynamics molecular simulation is deemed to be a reasonable approach for understanding the kinetics and thermodynamics of H<sub>2</sub>CO<sub>3</sub> formation in clay interlayers, and comparing with the reaction in the bulk, which has already been studied extensively in the past. Our results reveal striking differences for the reaction in nanopores compared with that in bulk water.

## Results and discussions

In Fig. 1, we present the cells used to simulate the reaction of a CO<sub>2</sub> molecule with bulk (Fig. 1a) and confined (Figs. 1b and 1c) water. In Figs. 1b and 1c, there are two (2W) and one (1W) water layers, confined in pyrophyllite (pph) nanopores. Pph is a dioctahedral phyllosilicate (Al<sub>2</sub>Si<sub>4</sub>O<sub>10</sub>(OH)<sub>2</sub>) comprising an octahedral (O) Al-centered sheet sandwiched between two tetrahedral (T) Si-centered sheets (i.e., T-O-T structure). Only one pph layer is explicitly represented in the simulation boxes (Figs. 1b and 1c). When the periodic boundary condition is applied in the z direction, water is confined in a pph nanopore. Note that pph does not swell in water; that is, the systems in Figs. 1b and 1c are selected as hypothetical models to investigate the reaction of CO<sub>2</sub> with confined water. This is a reasonable choice as our focus is on the effect of hydrophobic nanoconfinement on the free energy landscape of the conversion of CO<sub>2</sub> to H<sub>2</sub>CO<sub>3</sub>. In addition, the T-O-T structure of pph can be found in many swelling clays, including montmorillonite.



**Figure 1:** Simulation systems used to investigate the reaction of a CO<sub>2</sub> molecule (cyan-blue colors) with (a) bulk water, and confined water in pyrophyllite nanopores consisting of (b) 2W, and (c) 1W layers. The simulation box size for Fig. 1a is  $12.5 \times 12.5 \times 12.5 \text{ \AA}^3$ , with one CO<sub>2</sub> and 64 water molecules. The simulation cell sizes for Figs. 1b and c are  $20.64 \times 17.93 \times 15 \text{ \AA}^3$  (1 CO<sub>2</sub>, 92 H<sub>2</sub>O molecules, and a pph surface) and  $20.64 \times 17.93 \times 12 \text{ \AA}^3$  (1 CO<sub>2</sub> and 54 H<sub>2</sub>O molecules, and a pph surface).

The systems in Fig. 1 undergo equilibrium simulations with the ReaxFF force field,<sup>[68]</sup> as detailed in the computational methods section of the supporting information (SI). After equilibration, the free energy calculations are performed using well-tempered metadynamics,<sup>[71]</sup> available in the COLVARS<sup>[72]</sup> package built in the LAMMPS software package.<sup>[73]</sup> The free energy surfaces (FES, two-dimensional contour plots) represent the free energy for the formation of H<sub>2</sub>CO<sub>3</sub> from CO<sub>2</sub> and H<sub>2</sub>O as a function of two collective variables, CV1 and CV2, in each environment (bulk and confined water) (see Fig. 2).

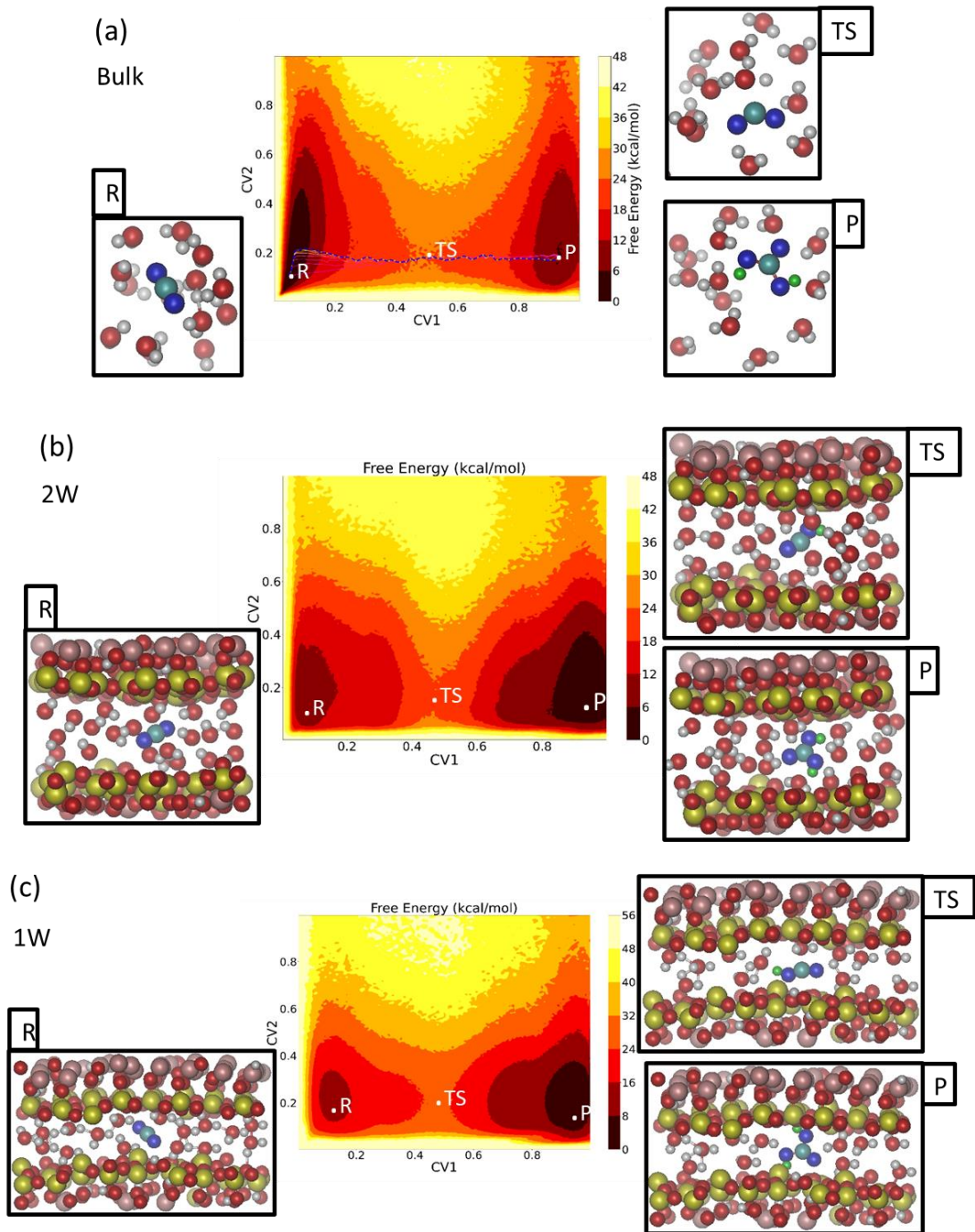
The first collective variable,  $CV1$ , is the coordination number of the carbon in  $CO_2$  (C) and oxygen of water ( $O_w$ ). The second collective variable,  $CV2$ , is the coordination number of the oxygen of  $CO_2$  ( $O_c$ ) and hydrogen of water (H). Variables  $CV1$  and  $CV2$  are defined in Eq. 3 and Eq. 4,

$$CV1 = \sum_{k \in O_w} \frac{1 - \left(\frac{r_{CO_wk}}{r_o}\right)^m}{1 - \left(\frac{r_{CO_wk}}{r_o}\right)^n} \quad (3)$$

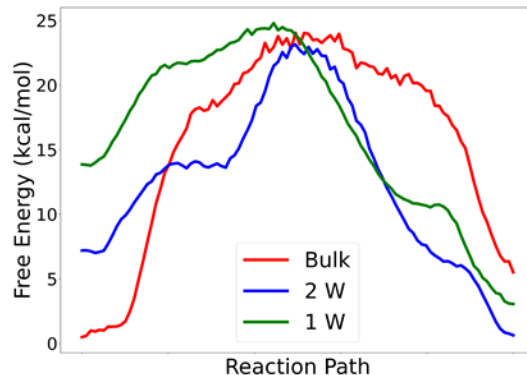
$$CV2 = \sum_{i \in O_c} \sum_{j \in H} \frac{1 - \left(\frac{r_{ij}}{r_c}\right)^m}{1 - \left(\frac{r_{ij}}{r_c}\right)^n} \quad (4)$$

where  $r_{CO_wk}$  in Eq. 3 represents the C- $O_w$  ( $k$  is the number of water molecules) and  $r_o$  is the cut-off distance,  $r_o = 1.6$  Å. The  $r_{ij}$  in Eq. 4 indicates the  $O_c$ -H distance [ $(i=1,2)$  and  $(j=1,2k)$ ] and  $r_c$  is the cut-off distance for  $r_{ij}$ ,  $r_c = 1.3$  Å. The exponential factors  $m$  and  $n$  have the values of 8 and 16. Small  $CV$  values (near 0) mean atoms are weakly coordinated while larger values (near 1) indicate strong coordination. More simulation details can be found in the SI.

The FES for the reaction in bulk water (Fig. 2a) comprises two distinguished wells, labeled as R and P, representing the reactant ( $CO_2$  and water complex) and product ( $H_2CO_3$  and water complex). The reaction coordinates ( $CV1$ ,  $CV2$ ) of points R and P are (0.051, 0.103) and (0.936, 0.186). The  $CV1$  value at point R is small, indicating that the coordination between C and  $O_w$  is weak, rendering  $CO_2$  as a mild solvation state ( $CO_2$  is a straight molecule at point R, snapshot R). A high  $CV1$  value at point P indicates that a C- $O_w$  bond has formed. Once the C- $O_w$  bond forms, proton transfer from any surrounding water molecule to form  $H_2CO_3$  is usually swift and the  $H_2CO_3$  is in a strongly solvated state (snapshot P). Note that formation of carbonic acid is thermodynamically more favorable than bicarbonate and carbonate ions formation.<sup>[76,77]</sup> We observe the *cis-trans* conformation for  $H_2CO_3$  regardless of the environment, which agrees with the fact that *cis-trans* is the lowest formation free energy out of the three conformers, including *trans-trans* and *cis-cis*.<sup>[23]</sup> The point TS at (0.497, 0.178) in Fig. 2a is the transition state. The method used to determine TS is discussed below, and presented in the SI. The TS snapshot in Fig. 2a shows the bending of the O-C-O angle, with a  $O_w$  close to the C atom. The O-C-O angle bending is important for  $H_2CO_3$  formation, regardless of solvation environment, and triggered by the C- $O_w$  coordination due to increasing  $CV1$  value driven by metadynamics simulation. The bending of O-C-O angle in TS of  $CO_2$  and  $H_2O$  reaction under bulk environment (see TS in Fig. 2a) agrees well with the overwhelming studies on  $CO_2$  hydration using ab-initio calculations in literature.<sup>[18,24,78]</sup> There has not been any studies on  $CO_2$  hydration under nanoconfinement as such which can be compared with our present TS configuration inside pph nanopore (see TS in Fig. 2b and Fig. 2c).



**Figure 2:** Free energy surface as a function of CV1 and CV2 obtained for  $\text{H}_2\text{CO}_3$  formation in (a) bulk water, and nanoconfined in (b) 2W and (c) 1W layers. The R, P, and TS labels on the FESs and snapshots represent the reactant, product, and transition states in each water environment. The color code is the same as in Fig. 1, except that the  $\text{H}_2\text{CO}_3$  and TS hydrogens are labelled in green. The blue dashed line on the FES in Fig. 2a is the minimum energy path obtained from the finite temperature string method<sup>[74,75]</sup> (see SI).



**Figure 3:** One-dimensional minimum free energy path for reactions in bulk water, and 2W and 1W confined water.

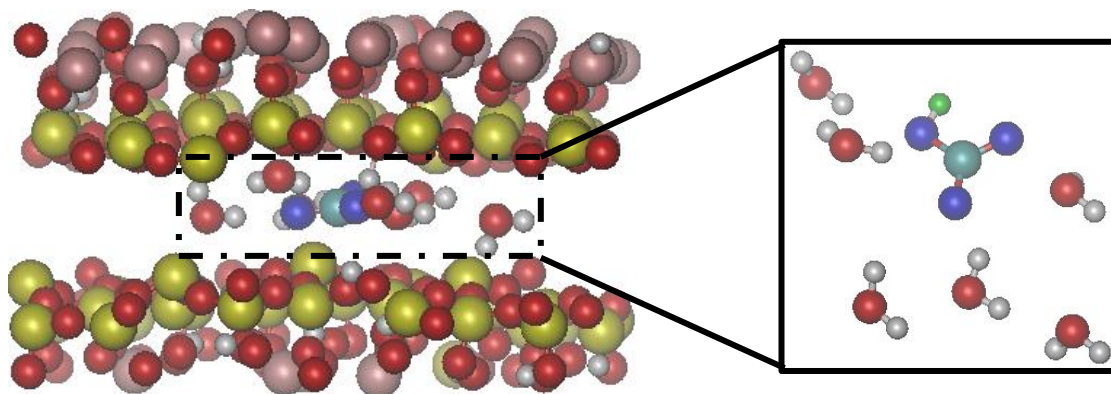
To estimate the free energy barrier and free energy of reaction for  $\text{H}_2\text{CO}_3$  formation in bulk water, we calculate the one-dimensional minimum energy path (MEP) (dashed blue line on FES in Fig. 2a) using the finite temperature string method (see SI for more details).<sup>[74,75]</sup> The one-dimensional free energy profile is also reported by the red line in Fig. 3. The energy barrier (difference between energy at TS and that at R) required for  $\text{H}_2\text{CO}_3$  formation is  $23.2 \pm 0.43$  kcal/mol, which is about 1.4 kcal/mol higher than the experimental value of 21.8 kcal/mol.<sup>[27]</sup> Our predicted barrier also agrees well with energy barriers calculated previously.<sup>[23,27,28]</sup> The product is about  $4.9 \pm 1.36$  kcal/mol higher in energy than the reactant (i.e., the well R is deeper than the well P on FES, Fig. 2a;  $\Delta G_{\text{bulk}} = 4.9 \pm 1.36$  kcal/mol). The geometries and relative free energies between the reactant, transition state, and product computed here are in good agreement with experimental<sup>[79]</sup> and theoretical studies<sup>[18,19,80]</sup> for bulk water and carbon dioxide. For example, the gas phase experimental<sup>[81,82]</sup> and theoretical<sup>[18]</sup> free energy of formation for  $\text{H}_2\text{CO}_3$  from  $\text{CO}_2$  and  $\text{H}_2\text{O}$  found in prior studies is 4.02 kcal/mol and 5.2 kcal/mol, respectively.

The results obtained for the reaction in pph nanopores are stunningly different from those obtained for the reaction in bulk water. Figs. 2b and 2c show the FES for the same reaction taking place in 2W and 1W pph nanopores. The one-dimensional free energy profiles for 2W and 1W nanopores are reported by the blue and green lines in Fig. 3. These results indicate that, when increasing confinement, the energy barriers decrease from  $23.2 \pm 0.43$  kcal/mol in bulk water to  $15.1 \pm 3.48$  kcal/mol in 2W nanopore, and to  $10.4 \pm 5.14$  kcal/mol in 1W nanopore. Furthermore, the reactant has lower energy than the product ( $\Delta G_{\text{bulk}} = 4.9 \pm 1.36$  kcal/mol) in bulk water. In contrast, the reactant has higher energy than the product for the reaction in 2W (i.e., the well P is deeper than well R on FES in Fig. 2b,  $\Delta G_{2\text{W}} = -6.5 \pm 2.71$  kcal/mol) and 1W

pores ( $\Delta G_{1W} = -13.4 \pm 4.33$  kcal/mol). This result suggests that hydrophobic nanoconfinement not only reduces the energy barrier, but also changes the thermodynamics of the reaction in favor of product ( $\text{H}_2\text{CO}_3$ ) formation. Another difference for the reaction in nanoconfined water, compared with the reaction in bulk water, is that charged species are usually observed in nanopores. Snapshots that represent the TS states in Figs. 2b and 2c already show a proton transferred to the oxygen of  $\text{CO}_2$ , indicating that the complex formed at TS is a charged species unlike the neutral complex in bulk water (TS state in Fig. 2a).

There are several reasons for the observed differences described above.

First, the free energy barrier is considerably lowered in nanopores compared to bulk water because a water molecule does not bind strongly to the surface under hydrophobic pph nanoconfinement, making it always available near the  $\text{CO}_2$  molecule for coordination and reaction. We observe that the free energy required to coordinate water molecules with  $\text{CO}_2$  to bend the O-C-O angle of a  $\text{CO}_2$  molecule to the O-C-O angle of  $\text{H}_2\text{CO}_3$  is high, as depicted by the energy barrier (Fig. 2a) to bring the system from R snapshot to TS snapshot. In other words, the reaction requires lower energy to coordinate water molecules with  $\text{CO}_2$  in hydrophobic nanopores compared with the energy required to break water hydrogen bonds and coordinate bulk water molecules with a  $\text{CO}_2$  molecule. This argument raises questions about the effect of hydrophilicity on the reaction. We will investigate this question in future work with a hypothesis that hydrophilic nanoconfinement will increase the energy barrier compared to energy barriers observed in bulk water and in hydrophobic confined water. Furthermore, steric hindrance imposed by the confining walls favor reactive encounters of the reactants.<sup>[40,43]</sup> Roughly speaking, confinement leaves only two translational degrees of freedom for the position of one reactant relative to the other, as opposed to three degrees of freedom available in the bulk.



**Figure 4:** Simulation snapshot demonstrating the formation of a bicarbonate ion in 1W nanopores. The color code is the same as Fig. 2.



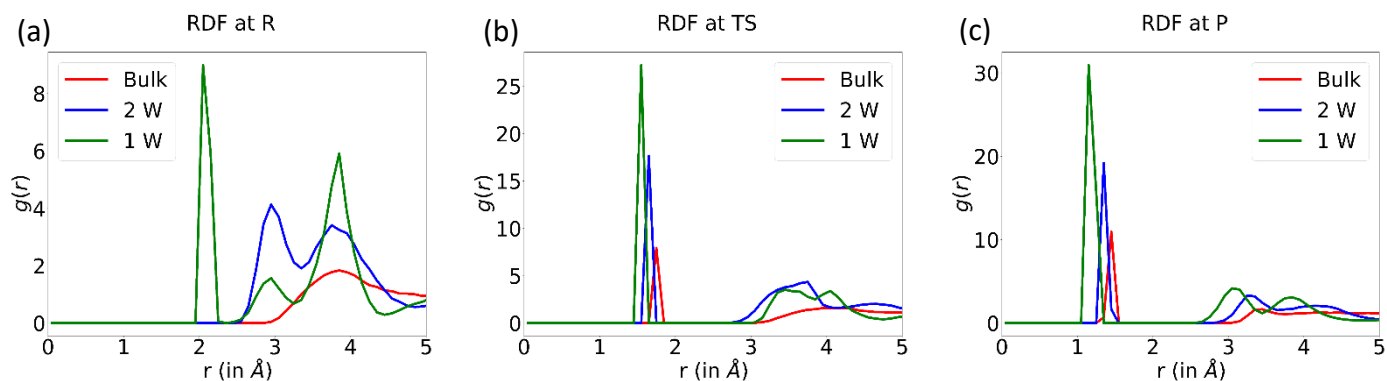
Second, the formation of  $\text{H}_2\text{CO}_3$  from  $\text{CO}_2$  and water involves hydrogen transfer that is likely enhanced under nanoconfinement. There is strong evidence that nanoconfinement enhances water self-dissociation,<sup>[36]</sup> and a qualitative relation between strongly varying components of the dielectric tensor with enhanced self-dissociation.<sup>[36]</sup> It has also been observed that the parallel component of the dielectric tensor of water at interfaces at ambient conditions increases significantly (up to 120) when approaching the surface from the bulk.<sup>[37]</sup> The perpendicular component of the dielectric tensor varies only from 0 to 2 upon approaching the interface.<sup>[40]</sup> We therefore postulate that proton donation to carbon dioxide is more likely in 1W nanopore than in 2W nanopore and in bulk water. In addition, nanoconfinement generically facilitates and stabilizes the formation of charged species.<sup>[9][40,43,83]</sup> We observe charged species at the TS states (Fig. 2) in nanopores and also the formation of  $\text{HCO}_3^-$  and  $\text{H}_3\text{O}^+$  in the nanoconfined regime intermittently before forming  $\text{H}_2\text{CO}_3$  (Fig. 4). The solvation energy  $\Delta G_s$  of charged species can be calculated using Born solvation equation

$$\Delta G_s = \omega(1/\varepsilon - 1) \quad (5)$$

where  $\omega$  is the Born solvation coefficient, and  $\varepsilon$  is the dielectric constant. The increase of the parallel component of the dielectric tensor inevitably decreases solvation energy (more negative), while the decrease of the perpendicular component of the dielectric tensor increases solvation energy (less negative). In the 2-dimensional channels (1W and 2W) solvation likely occur in the parallel direction. Therefore, charge species is more likely formed in the nanopores than in bulk water. However, note that bicarbonate ions formed intermittently in solutions due to their high free energy of formation ( $\Delta G_{\text{rxn}} = 4.7 \text{ kcal/mol}$ ) from carbonic acid.<sup>[76]</sup>

Third,  $\text{CO}_2$ , the TS, and  $\text{H}_2\text{CO}_3$  are coordinated by water molecules differently under different environments. We observe that the CV1 value for the state R in nanopores are higher than that in bulk water, indicating a close coordination of C and  $\text{O}_w$  under hydrophobic nanoconfinement. In other words, a  $\text{CO}_2$ -confined water complex (at point R, Fig. 2b and 2c) is thermodynamically favorable when the  $\text{CO}_2$  molecule is closely coordinated by  $\text{H}_2\text{O}$  molecules, while a  $\text{CO}_2$ -bulk water complex (at point R, Fig. 2a) is thermodynamically favorable with a looser  $\text{H}_2\text{O} - \text{CO}_2$  coordination. We can imagine the reaction process as follow: (i) water molecules closely coordinate with  $\text{CO}_2$  molecule, leading to the O-C-O angle bending (see TS in Fig. 2a), and then (ii) the C- $\text{O}_w$  bond forms, followed by proton transfers and hydration of  $\text{H}_2\text{CO}_3$ . The close coordination of  $\text{CO}_2$ -confined water helps lower the free energy of the first step (i). To further demonstrate the close coordination of  $\text{CO}_2$ -confined water, in Figure 5a we report the radial distribution function (RDF) of C and  $\text{O}_w$  at point R for different water environments (see SI for more details on this calculation). The RDF results indicate that, for the 1W case, the first peak (closest C- $\text{O}_w$  coordination) is at 2.12 Å. For the 2W case, the first peak is at 2.97 Å. In contrast,

the first peak is at 3.77 Å for bulk water. Similarly, the RDF results for TS and P states (Figs. 5b and 5c) also indicate a close coordination of these states, decreasing in the order 1W > 2W > bulk system.



**Figure 5:** Radial distribution function  $g(r)$  for C and  $O_w$  atoms in different environments for different states R (a), TS (b), and P (c).

The current research focuses on the  $\text{CO}_2$  to  $\text{H}_2\text{CO}_3$  conversion reaction in hydrophobic nanopores. To complete our understanding of the reactions in Eq. 1 in nanopores, future studies will focus on the same conversion reaction in hydrophilic and heterogeneous nanopores. Certainly, the surface charge distribution and interlayer cations will play a significant role in the thermodynamics and kinetics of the reaction because they can change the solvation and hydrogen bonding properties of the reaction constituents. Finally, the conversion of  $\text{H}_2\text{CO}_3$  to  $\text{HCO}_3^-$  and  $\text{CO}_3^{2-}$  ions in nanopores also needs to be investigated even though those reactions can be much more challenging using the computational method employed here.

## Conclusion

We investigated the thermodynamics and kinetics of the conversion of carbon dioxide to carbonic acid in bulk water and in water confined in hydrophobic pph nanopores using ReaxFF metadynamics molecular simulations. Our results demonstrated striking differences for the  $\text{H}_2\text{CO}_3$  formation in confined water and that in bulk water. The nanoconfined environment not only reduces the energy barrier but also changes the reaction from endothermic in bulk water to exothermic in nanoconfined water. In addition, during the formation of carbonic acid, charged species are observed more often in the nanopores than in bulk water. The increasing number of reactive encounters, improved molecular solvation and coordination, and favored proton transfer with increasing confinement enhances the thermodynamics and kinetics of the reaction.

As carbonation in nanopores is important in the carbon cycle and a complicated problem that depends on confinement, surface and pore chemistry, and concentration of CO<sub>2</sub>, our results provide mechanistic understanding to an important step in this process.

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**Keywords:** CO<sub>2</sub> conversion, ReaxFF, metadynamics, H<sub>2</sub>CO<sub>3</sub> formation, clay interlayer, pyrophyllite, carbon dioxide, bicarbonate ions, carbonate ions, molecular dynamics simulations

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