# Sella, an open-source automation-friendly molecular saddle point optimizer

Eric D. Hermes, Khachik Sargsyan, Habib N. Najm, and Judit Zádor<sup>∗</sup>

Combustion Research Facility, Sandia National Laboratories, Livermore, CA 94551-0969 USA

E-mail: jzador@sandia.gov

#### Abstract

We present a new algorithm for the optimization of molecular structures to saddle points on the potential energy surface using a redundant internal coordinate system. This algorithm automates the procedure of defining the internal coordinate system, including the handling of linear bending angles, e.g. through the addition of dummy atoms. Additionally, the algorithm supports constrained optimization using the nullspace sequential quadratic programming formalism. Our algorithm determines the direction of the reaction coordinate through iterative diagonalization of the Hessian matrix, and does not require evaluation of the full Hessian matrix. Geometry optimization steps are chosen using the restricted step partitioned rational function optimization method, and displacements are realized using a high-performance geodesic stepping algorithm. This results in a robust and efficient optimization algorithm suitable for use in automated frameworks. We have implemented our algorithm in Sella, an open source software package designed to optimize atomic systems to saddle point structures. We also introduce a new benchmark test comprising 500 molecular structures that approximate saddle point geometries and show that our saddle point optimization algorithm

outperforms the algorithms implemented in several leading electronic structure theory packages.

## 1 Introduction

The recent and upcoming availability of exascale computational resources has created renewed interest in high-performance automated software frameworks for the development of molecular reaction networks. A crucial step in such frameworks is the determination of saddle point geometries for the calculation of reaction rate coefficients under transition state theory. Automation of saddle point optimization requires robust optimization algorithms that converge quickly and reliably to saddle points. This has typically been a source of difficulty in the development of automated frameworks due to the computational expense involved and the rate of convergence failure many algorithms exhibit. Researchers often have to adjust saddle point initial guess structures manually to avoid pathological configurations and improve convergence of the optimization algorithm. There is therefore great need for development of new saddle point optimization algorithms whose reliability and performance is suitable for deployment in automated frameworks.

We have previously developed a method for saddle point optimization of densely connected systems (such as metals and their oxides in condensed phase) in Cartesian coordinates.<sup>1</sup> Our method was shown to converge in less than half the number of steps as compared to the next best performing code on one of the optbench.org saddle point optimization benchmarks. <sup>2</sup> This method was implemented in Sella, <sup>3</sup> an open source Python software package we have developed that has bindings to more than 20 popular electronic structure theory software packages due to its integration with the ASE software package.<sup>4</sup> However, Cartesian coordinates and the related algorithms are not efficient for saddle point searches on the potential energy landscapes of sparsely connected atomic systems, such as molecules, because of the strong coupling between atomic degrees of freedom. Molecular optimization (both for minima and saddle points)

is generally performed in a basis of internal coordinates, which typically comprise bond stretch distances, bending angles, and dihedral angles, as internal coordinates are significantly less energetically correlated than Cartesian coordinates. 5–9 These chemically relevant features provide a more effective coordinate system for optimization at the cost of additional algorithmic complexity. We have recently developed a new method for realizing displacements in a basis of redundant internal coordinates which dramatically reduces the number of steps required to reach convergence on a minimization benchmark.<sup>10</sup> In this work, we extend this algorithm to optimization of saddle point geometries.

Our optimization algorithm is designed to be used in automated frameworks without the need for manual intervention. However, the complexity of optimization in a basis of internal coordinates frequently frustrates attempts at automation. For example, when a bending angle of a molecule becomes linear or nearly linear during the course of optimization, displacement of the bending angle coordinate becomes ill-defined. To solve this problem, it is necessary to redefine the internal coordinate system to replace the linear bending angle with a different coordinate. Occasionally, this may require the addition of fictional "dummy atoms" that act as scaffolding off of which the replacement coordinate can be defined. The process of adding dummy atoms and dealing with the additional degrees of freedom they bring to the optimization problem poses a difficulty for automation efforts.

To address these issues and to enable efficient optimization of molecular saddle points, our algorithm defines internal coordinates and replaces pathological linear bending angles with more well-defined coordinates. If necessary, our algorithm will add dummy atoms to the system automatically at the beginning of the optimization for the purpose of defining new coordinates to replace linear bending angles. Moreover, if previously well-defined bending angles become linear mid-optimization, our algorithm will automatically replace those coordinates without interrupting the optimization procedure. This resolves the common problem of initially well-defined coordinate systems becoming invalid partway through optimization, resulting in a crash or other error.

When adding dummy atoms to the optimization problem, our algorithm also adds constraints to counteract the added Cartesian degrees of freedom generated by the dummy atoms. This enables the optimization of a wide range of molecular geometries in an efficient redundant internal coordinate system. This capability renders our algorithm a highly effective and robust geometry optimizer for automated workflows that explore reactive potential energy surfaces, such as KinBot.<sup>11</sup>

Further, constrained optimizations, such as relaxed scans, are often used to explore aspects of the potential energy surface and are thus also part of automated workflows. Additionally, the presence of dummy atoms necessitates the addition of constraints, as previously described. However, constrained calculations often fail to converge and are an impediment to automation. To address this challenge, we have implemented constrained saddle point optimization using sequential quadratic programming (SQP), resulting in reliable and efficient constrained optimization trajectories.<sup>12,13</sup>

In the following, we present our methodology to largely overcome the problems mentioned above. In section 2, we describe our saddle point optimization algorithm. This algorithm features the ability to automatically construct a basis of redundant internal coordinates from the Cartesian molecular structure, including the addition of dummy atoms when necessary. Furthermore, we have extended several methods to a basis of redundant internal coordinates, including iterative diagonalization of the Hessian, SQP-based constrained optimization, and restricted-step partitioned rational function optimization (RS-PRFO) for saddle point structures. In section 3 we introduce a new benchmark, demonstrate the performance of our algorithm, and give guidance about its use. We conclude in section 4.

### 2 The Optimization Algorithm

This section describes our redundant internal coordinate saddle point optimization algorithm, illustrated in figure 1. Section 2.1 describes the novel aspects of our algorithm for constructing internal coordinate system, in particular the addition of dummy atoms



Figure 1: A diagram of our internal coordinate saddle point optimization algorithm as implemented in Sella.

and how their associated degrees of freedom are handled. Section 2.2 discusses the construction and transformation of the various projected internal coordinate bases that are used in this work, including the nonredundant internal coordinates, the constrained internal coordinates, and the unconstrained nonredundant internal coordinates. Section 2.3 describes our iterative Hessian diagonalization algorithm, as well as the multi-secant Hessian update formula that is used to construct and update the approximate Hessian matrix. Section 2.4 describes how we have combined the null-space SQP constrained optimization formalism with the RS-PRFO saddle point optimization procedure.

#### 2.1 Handling linear angles using dummy atoms

An internal coordinate system describes the geometry of a molecule using the relative positions and orientations of atoms within the molecule, in contrast to the Cartesian coordinate system which describes the geometry using the absolute positions of each atom. The most common types of internal coordinates are the two-atom bond distances, three-atom bending angles, and four-atom dihedral angles.<sup>5,8</sup> A sufficiently large set of these three types of internal coordinates can fully represent all internal degrees of freedom of a molecule. To facilitate automation, our algorithm is capable of automatically determining an appropriate set of internal coordinates. Notably, the algorithm handles systems with linear bending angles by replacing those coordinates with alternate coordinates, adding dummy atoms where necessary. Without this functionality, optimization algorithms in internal coordinates will often crash due to derivative discontinuities in bending angle degrees of freedom.

There are many ways to construct an internal coordinate system for a molecule. Perhaps the simplest approach is the Z-matrix representation, which defines a minimal internal coordinate system. However, the Z-matrix representation necessarily omits chemically relevant features for non-trivial molecules, particularly molecules with rings. This has a deleterious effect on performance, as the excluded chemically relevant features may change drastically after even a small displacement in the Z-matrix coordinate space.<sup>6</sup>

For this reason, our algorithm uses a redundant internal coordinate system that includes all chemically relevant bond stretch, angle bend, and dihedral angle coordinates. Redundant internal coordinate systems increase the complexity of the optimization problem, as the coordinate system is by construction of a higher dimensionality than the space of all possible molecular configurations. We have described these problems and the methods we developed and implemented in Sella to ameliorate them in a previous work.<sup>10</sup> Our algorithm constructs a complete set of internal coordinates by mapping all possible pairs, triplets, and quadruplets of sequentially-bonded atoms from a molecular graph to bond distance, bending angle, and dihedral angle coordinates, respectively. This procedure is described in detail in appendix A.

One problem that is frequently encountered when constructing internal coordinate systems is bending angles may become nearly or completely linear. This is a problem because linear bending angles do not have a well-defined Cartesian derivative, and dihedral angle coordinates which contain a linear angle are undefined. This problem is typically solved by manually replacing the pathological angle with one or more new coordinates, usually an improper dihedral or in- and out-of-plane bending angles. 8,14 However, if the central atom of the pathological angle is only bonded to two other atoms, then it is not possible to define a chemically meaningful improper dihedral to replace the angle. In that case, it is necessary to define a new, so-called "dummy" atom, which is used to help define the improper dihedral. For this scenario, we have developed an algorithm for automatically adding dummy atoms to a molecule, potentially even mid-optimization. This is a key aspect of our method, as without automatic reconstruction of internal coordinates, some molecular saddle point optimizations will fail due to the coordinate system becoming ill-defined. To the best of our knowledge, no other optimizer automatically reconstructs the internal coordinate system mid-optimization when the initially constructed internal coordinate system becomes ill-defined.

As dummy atoms are fictitious, the energy of the molecule does not depend on their position, and, therefore, the gradient of the potential energy with respect to the position of dummy atoms is always zero. Because of this, it is possible for the position of the dummy atoms to drift during the course of optimization, reducing the effectiveness of the improper dihedral for which it was added. Our algorithm counteracts this problem by automatically applying three constraints to the position of the dummy atom to ensure it remains in a position that is useful for the optimization. The addition of these constraints also serves to avoid an increase in the dimensionality of the optimization problem upon addition of dummy atoms, though at the cost of turning an otherwise unconstrained optimization into a constrained optimization problem.



Figure 2: A schematic of how dummy atoms are added to near-linear angles. An angle  $\alpha_{abc}$ is considered near-linear when it is above 165°. So long as  $\alpha_{abc}$  is not perfectly linear, the dummy atom (x) is added in a direction perpendicular to the plane defined by the angle's three atoms. The bond stretch coordinate  $\delta_{bx}$  is constrained to 1 Å and the two bending angles  $\alpha_{abx}$  and  $\alpha_{cbx}$  are constrained to 90°. The nearly-linear angle is replaced by the improper dihedral  $\tau_{abxc}$ .

Our algorithm automatically places dummy atoms following a precise set of rules that are designed to ensure consistent performance even when a molecule is translated, rotated, or if the atom order is permuted. If a proposed angle  $\alpha_{abc}$  is found to be within 15° of linear (0° or 180°) during the construction of the internal coordinate system, then the angle is not added and an improper dihedral is added instead. If atom b is bonded to at least three atoms, then the improper dihedral  $\tau_{abdc}$  is added to the coordinate system, where  $d$  is chosen to be the bonded atom closest to  $b$  excluding  $a$  and  $c$ . If atom b is bonded only to a and c, a dummy atom is added to the coordinate system, as illustrated in figure 2. The dummy atom x is placed  $1 \text{\AA}$  away from atom b in the direction of  $\hat{\delta}_{ba} \times \hat{\delta}_{bc}$ , where  $\hat{\delta}$  is the unit bond vector. To counteract the dummy atom's additional Cartesian degrees of freedom, the  $\delta_{bx}$  bond stretch coordinate is constrained

to 1 Å and the  $\alpha_{abx}$  and  $\alpha_{cbx}$  bending angles are constrained to 90°. The nearly-linear angle  $\alpha_{abc}$  is then replaced with the improper dihedral angle  $\tau_{abxc}$ .

However, if the  $\alpha_{abc}$  angle is already too close to linear when the algorithm detects the need for the dummy atom, then this procedure for determining the orientation of the dummy atom will become numerically unstable. For this reason, if  $\|\hat{\boldsymbol{\delta}}_{ba} \times \hat{\boldsymbol{\delta}}_{bc}\|_2 < 10^{-4}$ , an alternate procedure is used instead. This procedure places the dummy atom in the direction  $(\mathbf{I} - \hat{\boldsymbol{\delta}}_{ac} \hat{\boldsymbol{\delta}}_{ac}^T) \hat{\mathbf{w}}$ , where  $\hat{\mathbf{w}}$  is the unit vector of the Cartesian axis that is nearest to being orthogonal with the vector  $\hat{\delta}_{ac}$ . While this approach loses rotational invariance of the dummy atom positioning, the positioning does maintain translational and permutational invariance.

During the course of an optimization trajectory, if a previously-defined bending angle is found to be within 15° of linear, then the internal coordinate system is discarded and subsequently regenerated by re-running the molecular graph algorithm described in this section.

#### 2.2 Definition and conversion of delocalized coordinate sys-

#### tems

In this work, we extend a method that we have previously developed for saddle point optimization in a Cartesian coordinate basis to a basis of redundant internal coordinates.<sup>1</sup> When using redundant internal coordinates or when optimizing a system with constraints, our algorithm uses a variety of projected internal coordinate bases as well. This section outlines how these projected bases are constructed and the relationships between them.

The Cartesian coordinate vector of a molecule is represented by the 3n-dimensional vector  $x$ , where  $n$  is the number of atoms in the molecule. The redundant internal coordinate vector is represented by the m-dimensional vector **q**, where  $m \geq 3n - 6$ . Each element of q corresponds to one of the internal coordinates determined by the algorithm described in appendix A, and  $m$  is the total number of internal coordinates

that have been determined. In practice, only the Cartesian coordinate vector  $\bf{x}$  behaves as an independent variable for the optimization algorithm, and the internal coordinate vector  $q$  is calculated from  $x$ . Therefore, all changes to the internal coordinate vector  $q$ are realized as changes to  $x$ ; this has important implications for the iterative diagonalization algorithm described in section 2.3 and the constrained saddle point optimization algorithm described in section 2.4 below. In this work, all undecorated vector and matrix quantities are defined to be in a basis of redundant internal coordinates with the exception of the Cartesian coordinate vector x. Quantities in other coordinate systems will be decorated with an accent character, as described below.

In addition to the Cartesian and redundant internal coordinate bases, our algorithm also uses a nonredundant basis which is a linear combination of the redundant internal coordinates.<sup>7</sup> We define the nonredundant internal coordinate basis using the singular vectors of the  $m \times 3n$  internal–Cartesian Jacobian matrix **J** (sometimes called the Wilson B-matrix<sup>15</sup>) as projection matrices,

$$
\mathbf{J} = \frac{d\mathbf{q}}{d\mathbf{x}} = \begin{bmatrix} \mathbf{N} & \mathbf{R} \end{bmatrix} \begin{bmatrix} \mathbf{T} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} \end{bmatrix} \begin{bmatrix} \mathbf{Z}^T \\ \mathbf{Z}'^T \end{bmatrix},
$$
(1)

where **T** is the  $(3n-6) \times (3n-6)$  diagonal matrix of non-zero singular values, N is the  $m \times (3n-6)$  orthonormal matrix whose columns span the nonredundant space of q, R is the  $m \times (m - 3n + 6)$  orthonormal matrix whose columns span the redundant space of q, and Z and Z' are the  $3n \times (3n-6)$  and  $3n \times 6$  matrices whose columns are the right singular vectors corresponding to  $N$  and  $R$  respectively. We treat singular values of **J** below  $10^{-6}$  as zero and assign the corresponding singular vectors to the redundant space.

In this work, vectors and matrices in the nonredundant internal coordinate space will be decorated with a check accent (e.g.  $\check{\mathbf{v}}$ ). Using our definition of the nonredundant internal coordinate system, vectors  $\bf{v}$  and matrices  $\bf{M}$  can be transformed between the redundant and nonredundant internal coordinate systems through the relations

$$
\check{\mathbf{v}} = \mathbf{N}^T \mathbf{v} \quad \text{and} \quad \mathbf{v} = \mathbf{N} \check{\mathbf{v}} \tag{2}
$$

$$
\tilde{\mathbf{M}} = \mathbf{N}^T \mathbf{M} \mathbf{N} \quad \text{and} \quad \mathbf{M} = \mathbf{N} \tilde{\mathbf{M}} \mathbf{N}^T. \tag{3}
$$

The nonredundant internal coordinate system is used primarily in the iterative Hessian diagonalization algorithm described in section 2.3 below.

For constrained optimizations using SQP, our algorithm further subdivides the nonredundant coordinate system into an unconstrained nonredundant internal coordinate system and a constrained internal coordinate system. Constraints are represented by the equation

$$
\mathbf{c}(\mathbf{q}) = \mathbf{0},\tag{4}
$$

where each element of the *l*-dimensional vector  $\bf{c}$  corresponds to an individual constraint expression. The constraints, which may have redundancy, reduce the dimensionality of the optimization problem to  $3n - 6 - p$ , where  $p \leq l$  indicates the number of linearly independent constraints. We define the unconstrained and constrained subspaces using the right singular vectors of the  $l \times (3n-6)$  constraint–nonredundant Jacobian matrix as projection matrices,

$$
\frac{d\mathbf{c}}{d\check{\mathbf{q}}} = \begin{bmatrix} \check{\mathbf{U}} & \check{\mathbf{U}}' \end{bmatrix} \begin{bmatrix} \check{\mathbf{W}} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} \end{bmatrix} \begin{bmatrix} \check{\mathbf{P}}^T \\ \check{\mathbf{Q}}^T \end{bmatrix},
$$
(5)

where  $\check{\mathbf{W}}$  is the  $p \times p$  diagonal matrix of non-zero singular values,  $\check{\mathbf{P}}$  is the  $(3n-6) \times p$ orthonormal matrix whose columns span the constrained subspace of  $\check{\mathbf{q}}, \check{\mathbf{Q}}$  is the  $(3n 6) \times (3n-6-p)$  orthonormal matrix whose columns span the unconstrained subspace of  $\check{\mathbf{q}}$ , and  $\check{\mathbf{U}}'$  and  $\check{\mathbf{U}}'$  are the  $l \times p$  and  $l \times (l-p)$  matrices whose columns are the left-singular vectors corresponding to  $\dot{\mathbf{P}}$  and  $\dot{\mathbf{Q}}$  respectively. The constraint–nonredundant Jacobian can be calculated from the constraint–Cartesian Jacobian and the singular values and

right singular vectors of the redundant–Cartesian Jacobian from equation 1,

$$
\frac{d\mathbf{c}}{d\check{\mathbf{q}}} = \frac{d\mathbf{c}}{d\mathbf{x}} \left(\frac{d\check{\mathbf{q}}}{d\mathbf{x}}\right)^+ = \frac{d\mathbf{c}}{d\mathbf{x}} \left(\frac{d\mathbf{q}}{d\mathbf{x}}\right)^+ \mathbf{N} = \frac{d\mathbf{c}}{d\mathbf{x}} \mathbf{Z} \mathbf{T}^{-1} \mathbf{N}^T \mathbf{N} = \frac{d\mathbf{c}}{d\mathbf{x}} \mathbf{Z} \mathbf{T}^{-1},\tag{6}
$$

where a superscript plus symbol indicates the Moore–Penrose pseudoinverse. <sup>16,17</sup> As before, we treat singular values below  $10^{-6}$  as zero and assign the corresponding singular vectors to the unconstrained space. Using the projection matrices defined in equations 1 and 5 above, we define the  $m \times p$  matrix **P** and the  $m \times (3n - 6 - p)$  matrix **Q** for projection directly from the redundant internal coordinate system to the constrained and unconstrained nonredundant internal coordinate spaces respectively as

$$
\mathbf{P} = \mathbf{N}\check{\mathbf{P}} \tag{7}
$$

$$
Q = N\check{Q}.
$$
 (8)

The projection matrix  $N$  projects from the redundant internal coordinates to the nonredundant internal coordinates, and  $\check{P}$  and  $\check{Q}$  project from the nonredundant internal coordinates to the constrained and unconstrained coordinates, respectively.

In this work, vectors and matrices in unconstrained internal coordinates will be decorated with a tilde accent (e.g.  $\tilde{\mathbf{v}}$ ). Using our definition of the unconstrained internal coordinate system, vectors  $\bf{v}$  and matrices  $\bf{M}$  can be converted between redundant internal coordinates and unconstrained internal coordinates through the relations

$$
\tilde{\mathbf{v}} = \mathbf{Q}^T \mathbf{v} \quad \text{and} \quad \mathbf{v} = \mathbf{Q}\tilde{\mathbf{v}} \tag{9}
$$

$$
\tilde{\mathbf{M}} = \mathbf{Q}^T \mathbf{M} \mathbf{Q} \quad \text{and} \quad \mathbf{M} = \mathbf{Q} \tilde{\mathbf{M}} \mathbf{Q}^T.
$$
 (10)

The unconstrained internal coordinate system is used primarily in the geometry optimization algorithm described in section 2.4 below.

A summary of the relationships between the various internal coordinate systems is illustrated in figure 3, where the length of the bars correlates with the dimensionality of the coordinate space. As the name implies, the nonredundant internal coordinate



Figure 3: A graphical representation of the relationship between the different internal coordinate systems in this work. The full redundant internal coordinate space q (dark blue) is projected into the smaller nonredundant space  $\check{\mathbf{q}}$  (green) with the projection matrix  $\mathbf{N}$ . This space is further split into the constrained space (yellow) and the unconstrained space  $\tilde{q}$  (red), which are obtained from q with the projection matrices P and Q, respectively.

space spans only the nonredundant part of the redundant internal coordinates. The constrained and unconstrained nonredundant spaces together also span the nonredundant space. The redundant internal coordinates can be projected into the nonredundant space, the constrained space, and the unconstrained space with the  $N$ ,  $P$ , and  $Q$  matrices respectively. The nonredundant internal coordinates can be projected into the constrained and unconstrained nonredundant space with the  $\dot{P}$  and  $\dot{Q}$  matrices respectively. Note that when performing an unconstrained optimization, the nonredundant and unconstrained nonredundant spaces are identical and the constrained space is zerodimensional.

## 2.3 Iterative Hessian diagonalization in nonredundant internal coordinates

To optimize the geometry of a molecule to a first order saddle point, some information about the curvature of the potential energy surface is required. This information is contained within the Hessian matrix, some (or all) of which must be known to step towards a saddle point. Optimization to a first order saddle point requires stepping uphill on the potential energy surface in the direction of the reaction coordinate and

downhill in all other directions. We define the reaction coordinate as the direction in which the curvature is smallest (most negative or least positive), which is by construction the leftmost eigenvector of the Hessian matrix. Therefore it is at least necessary to know this eigenvector to step towards a first order saddle point.

For molecular saddle point optimization, it is common to evaluate the full Hessian matrix at least once at the beginning of optimization. Many electronic structure theory methods such as Hartree–Fock and DFT have efficient analytical Hessian implementations. Additionally, transformation of Cartesian Hessians into internal coordinates is straightforward. When using a level of theory for which efficient analytical second derivatives are not available, the Hessian matrix is usually evaluated in full using finite difference. For computationally expensive levels of theory or for very large molecules, the evaluation of the initial Hessian matrix may become more costly than all subsequent geometry optimization steps. Even when using levels of theory with efficient analytical second derivatives, evaluation of the Hessian matrix will eventually dominate the overall cost of optimization as system sizes grow larger due to the higher computational scaling of Hessian evaluations relative to gradient evaluations. For this reason, saddle point optimization methods which do not require full evaluation of the Hessian matrix may be useful for expensive levels of theory and for large systems.

We have previously developed a saddle point optimization algorithm that efficiently identifies the leftmost eigenvector of the Hessian matrix using iterative diagonalization.<sup>1</sup> Furthermore, this algorithm uses the information obtained by the iterative diagonalization procedure to construct an approximate Hessian that is exact in the subspace searched by the eigensolver. Our method can be compared to that of Sharada, Bell, and Head-Gordon,<sup>18</sup> with the key difference being that *all* curvature information obtained by the eigensolver is incorporated into the approximate Hessian, not just the lowest curvature mode. This results in a much more accurate approximate Hessian matrix, and therefore a substantially more efficient optimization algorithm. However, as our method was originally designed only for optimization in a Cartesian coordinate basis, it was not suitable for optimization of molecular saddle point geometries. In this work,

we extend our iterative diagonalization and approximate Hessian construction methods to optimization in a basis of redundant internal coordinates, potentially with added constraints. Our new method diagonalizes the Hessian and constructs the approximate Hessian directly in an internal coordinate basis. This is in contrast to the method of Sharada et al.,<sup>18</sup> which diagonalizes the Hessian and constructs the approximate Hessian in a Cartesian basis before transforming it into an internal coordinate basis.

Our optimization algorithm enforces constraints using the SQP formalism, which is based on the method of Lagrange multipliers. The Lagrangian  $\mathcal L$  is defined as

$$
\mathcal{L}(\mathbf{q}, \mathbf{w}) = \mathcal{E}(\mathbf{q}) - \mathbf{w}^T \mathbf{c}(\mathbf{q}),
$$
\n(11)

where  $\mathcal E$  is the potential energy, c is the vector of constraints, and w are the corresponding Lagrange multipliers. 19,20

The goal of the iterative diagonalization procedure for saddle point optimization, outlined in algorithm 1, is to find the leftmost eigenvector of the Hessian of the potential energy H. For constrained optimization, the goal is instead to find the leftmost eigenvector of the Hessian of the Lagrangian  $H_{\mathcal{L}}$ , as this is the Hessian (or approximation thereof) that is used in the optimization procedure (see section 2.4). When performing an optimization in an internal coordinate basis, our algorithm iteratively diagonalizes the Hessian in the *nonredundant* internal coordinate basis  $\dot{q}$ . One could alternately perform iterative diagonalization of the Hessian in the unconstrained nonredundant space  $\tilde{\mathbf{q}}$ . However, it is useful to know the energetic coupling between the constrained and unconstrained degrees of freedom, particularly for extrapolation of the gradient as in equation 26 below. For this reason, we include constrained degrees of freedom in the iterative Hessian diagonalization procedure.

The iterative diagonalization procedure works by generating an orthonormal matrix  $\check{\mathbf{S}}_i$  column-by-column using the Olsen method,  $^{21}$  which is described below. The columns of the matrix  $\check{\mathbf{S}}_i$  are used to define a projected coordinate system, and in this section vectors and matrices in this space are decorated with a grave accent (e.g.  $\dot{\mathbf{v}}$ ). The

Algorithm 1 Iterative diagonalization algorithm

1: procedure RAYLEIGHRITZ 2: input: q, g,  $B_k$ 3: output:  $B_{k+1}$ 4: if first diagonalization then 5:  $\check{\mathbf{s}}_1 \leftarrow \frac{\check{\mathbf{g}}}{\|\check{\mathbf{g}}\|_2}$ 6: else 7:  $\check{\mathbf{s}}_1 \leftarrow \text{leftmost eigenvector of } \check{\mathbf{B}}_{\mathcal{L}}$ 8:  $\check{\mathbf{S}}_1 \leftarrow [\check{\mathbf{s}}_1]$ 9:  $\check{\mathbf{Y}}_0 \leftarrow \begin{bmatrix} \cdot \end{bmatrix}$ 10:  $i \leftarrow 1$ 11: loop 12: evaluate  $\check{\mathbf{y}}_i = \check{\mathbf{H}}_L \check{\mathbf{s}}_i$  through forward finite difference (see equation 12) 13:  $\check{\mathbf{Y}}_i \leftarrow \begin{bmatrix} \check{\mathbf{Y}}_{i-1} & \check{\mathbf{y}}_i \end{bmatrix}$ 14: diagonalize  $\check{\mathbf{Y}}_i^T \check{\mathbf{S}}_i$  to obtain leftmost eigenpair  $\theta_i$  and  $\dot{\mathbf{z}}_i$ 15:  $\check{\mathbf{z}}_i \leftarrow \check{\mathbf{S}}_i \grave{\mathbf{z}}_i$ 16:  $\check{\mathbf{r}}_i \leftarrow \check{\mathbf{Y}}_i \dot{\mathbf{z}}_i - \theta_i \check{\mathbf{z}}_i$ 17: **if**  $\|\check{\mathbf{r}}_i\|_2 < \gamma |\theta_i|$  or  $i = 3n - 6$  **then** 18: **exit** and calculate  $B_{k+1}$  using equations 17 through 20 19: Solve equation 14 for  $\check{\mathbf{t}}_i$  $20:$  $\vec{\Sigma}_i \gets \left(\check{\mathbf{I}} - \check{\mathbf{S}}_i\check{\mathbf{S}}_i^T\right)$  $\left(\begin{matrix} T \\ i \end{matrix}\right)$   $\tilde{\mathbf{t}}_i$  (in practice, use modified Gram-Schmidt) 21:  $\check{\mathbf{s}}_{i+1} \leftarrow \frac{\check{\mathbf{t}}'_i}{\|\check{\mathbf{t}}'_i\|_2}$ 22:  $\mathbf{\check{S}}_{i+1} \leftarrow \begin{bmatrix} \mathbf{\check{S}}_i & \mathbf{\check{s}}_{i+1} \end{bmatrix}$ 23:  $i \leftarrow i + 1$ 

Hessian of the Lagrangian  $\check{H}_{\mathcal{L}}$  is evaluated in this space using forward finite difference for the energy term, while the constraint term is evaluated analytically,

$$
\check{\mathbf{H}}_{\mathcal{L}}\check{\mathbf{s}} \approx \frac{\check{\mathbf{g}}(\check{\mathbf{q}} + \eta \check{\mathbf{s}}) - \check{\mathbf{g}}(\check{\mathbf{q}})}{\eta} - \sum_{i} w_{i} \frac{\partial^{2} c_{i}}{\partial \check{\mathbf{q}}^{2}} \check{\mathbf{s}},
$$
(12)

where  $\check{\mathbf{s}}$  is one column of  $\check{\mathbf{S}}_i$ ,  $\check{\mathbf{g}}$  is the gradient of the potential energy,  $\eta$  is a small number that controls the magnitude of the displacement, and the Lagrange multipliers w are chosen to be the least-squares solution to the first-order optimality condition,

$$
\frac{\partial \mathcal{L}}{\partial \mathbf{q}} = \mathbf{g} - \mathbf{w}^T \frac{\partial \mathbf{c}}{\partial \mathbf{q}} = \mathbf{0}.
$$
 (13)

Note that least-squares solutions to equation 13 will generally have non-zero residual at all points except the final converged geometry. The Hessian-matrix products from equation 12 are stored in the matrix  $\check{\mathbf{Y}}_i = \check{\mathbf{H}}_{\mathcal{L}} \check{\mathbf{S}}_i$ . Additionally, to update the approximate potential energy Hessian matrix B after the diagonalization algorithm has converged, the *potential energy* Hessian-vector products  $\check{H}\check{S}_i$  are stored separately in the matrix  $\check{\mathbf{Y}}_{\mathcal{E}}$ . The iterative diagonalization procedure is considered converged once the space spanned by the columns of  $\check{\mathbf{S}}_i$  is found to contain the leftmost eigenvector of  $\check{\mathbf{H}}_{\mathcal{L}}.$ 

As mentioned in section 2.2, displacements in the internal coordinate space must be realized as changes to the Cartesian coordinates x, as these are the independent coordinate variables. In our algorithm, these displacements are realized using a geodesic internal coordinate stepping approach we have previously developed for minimization in a redundant internal coordinate basis.  $^{\rm 10}$ 

At each iteration of the diagonalization algorithm, the part of the Hessian which has already been evaluated  $\hat{\mathbf{H}}_i = \check{\mathbf{Y}}_i^T \check{\mathbf{S}}_i$  is exactly diagonalized to find the leftmost eigenvalue  $\theta_i$  and eigenvector  $\dot{\mathbf{z}}_i$ . This eigenvector, when projected into the larger nonredundant internal coordinate space through the relation  $\check{\mathbf{z}}_i = \check{\mathbf{S}}_i \hat{\mathbf{z}}_i$ , approximates the leftmost eigenvector of  $H_{\mathcal{L}}$ . The diagonalization procedure is repeated until the convergence criterion  $\|\check{\mathbf{r}}_i\|_2 < \gamma |\theta_i|$  is met, where  $\check{\mathbf{r}}_i = \check{\mathbf{Y}}_i \dot{\mathbf{z}}_i - \theta_i \check{\mathbf{z}}_i$  is the residual vector

and  $\gamma$  is a unitless convergence parameter which defaults to 0.1 in Sella. This procedure is a particular implementation of the Rayleigh-Ritz method. <sup>22</sup>

Implementations of the Rayleigh-Ritz method differ in how the matrix  $\check{\mathbf{S}}_i$  is constructed. Our implementation uses the Olsen method, <sup>21</sup> a method related to the generalized Davidson<sup>23,24</sup> and Jacobi-Davidson<sup>25–27</sup> diagonalization algorithms. Olsen's method extends the matrix  $\check{\mathbf{S}}_i$  with the vector obtained by orthonormalizing the vector  $\check{\mathbf{t}}_i$  against  $\check{\mathbf{S}}_i$ , where  $\check{\mathbf{t}}_i$  is the solution to the equation

$$
\left(\check{\mathbf{I}} - \check{\mathbf{z}}_i \check{\mathbf{z}}_i^T\right) \left(\check{\mathbf{B}}_{\mathcal{L}} - \theta_i \check{\mathbf{I}}\right) \left(\check{\mathbf{I}} - \check{\mathbf{z}}_i \check{\mathbf{z}}_i^T\right) \check{\mathbf{t}}_i = -\check{\mathbf{r}}_i,\tag{14}
$$

where  $\check{\mathbf{I}}$  is the identity matrix and  $\check{\mathbf{B}}_{\mathcal{L}}$  is the current approximate Hessian of the Lagrangian,

$$
\mathbf{B}_{\mathcal{L}} = \mathbf{B} - \sum_{i} w_{i} \frac{\partial^{2} c_{i}}{\partial \mathbf{q}^{2}},
$$
\n(15)

where **B** is the approximate Hessian of the potential energy used in the optimization algorithm. In our implementation, the Hessians of the constraints are evaluated in a Cartesian coordinate basis using automatic differentiation and then converted to an internal coordinate basis through the chain rule relation

$$
\frac{\partial^2 c_i}{\partial \mathbf{q}^2} = \mathbf{J}^{+T} \left( \frac{\partial^2 c_i}{\partial \mathbf{x}^2} - \sum_j \left[ \frac{\partial c_i}{\partial \mathbf{x}} \mathbf{J}^+ \right]_j \frac{\partial^2 q_j}{\partial \mathbf{x}^2} \right) \mathbf{J}^+.
$$
 (16)

Once the iterative diagonalization algorithm has converged, the matrix of displacement vectors S and the corresponding potential energy Hessian-vector products  $Y_{\mathcal{E}}$ are projected into the full redundant internal coordinate basis and used to update the approximate potential energy Hessian B. This is done with the multi-secant TS-BFGS Hessian update formula,  $1,28-31$ 

$$
\mathbf{B}_{k+1} = \mathbf{B}_k + \mathbf{U} (\mathbf{Y}_{\mathcal{E}} - \mathbf{B}_k \mathbf{S})^T + (\mathbf{Y}_{\mathcal{E}} - \mathbf{B}_k \mathbf{S}) \mathbf{U}^T - \mathbf{U} (\mathbf{Y}_{\mathcal{E}} - \mathbf{B}_k \mathbf{S})^T \mathbf{S} \mathbf{U}^T
$$
 (17)

$$
\mathbf{U} = \mathbf{M}\mathbf{S} \left[ \mathbf{S}^T \mathbf{M} \mathbf{S} \right]^{-1} \tag{18}
$$

$$
\mathbf{M} = \mathbf{Y}_{\mathcal{E}} \mathbf{Y}_{\mathcal{E}}^T + |\mathbf{B}_k| \mathbf{S} \mathbf{S}^T |\mathbf{B}_k|
$$
\n(19)

$$
|\mathbf{B}_k| = \sum_i |\lambda_k^{(i)}| \mathbf{v}_k^{(i)} \mathbf{v}_k^{(i)T},\tag{20}
$$

where  $\lambda_k^{(i)}$  $_k^{(i)}$  and  $\mathbf{v}_k^{(i)}$  $\mathbf{R}_k^{(i)}$  are respectively the *i*th eigenvalue and eigenvector of  $\mathbf{B}_k$ . At the beginning of optimization, the approximate potential energy Hessian B is initialized as a diagonal matrix using the method of Fischer and Almlof. <sup>32</sup>

## 2.4 Constrained optimization to a saddle point with SQP and RS-PRFO

Constrained optimization of molecular geometries is commonly used in computational chemistry, for example to scan along specific internal coordinates. There are several algorithms for the enforcement of constraints on molecular geometry optimization. RAT-TLE and SHAKE are two well-known algorithms that rigorously enforce  $\mathbf{c}(\mathbf{x}_k) = \mathbf{0}$  for all geometries  $\mathbf{x}_k$  visited during the course of the optimization trajectory.<sup>33,34</sup> This is important for molecular dynamics simulations, where conservation of energy requires that all structures of the trajectory lay on the constraint manifold. In contrast, optimization trajectories do not need to conserve energy, and, therefore, only the final converged structure must lie on the constraint manifold. For this reason, we have chosen not to use RATTLE or SHAKE for constrained optimization in our framework. Instead, our algorithm performs constrained optimization using null-space SQP, a method of Lagrange multipliers in which a linear step towards the constraint manifold is made at every iteration of the optimization. Using this approach, intermediate steps of optimization may not rigorously satisfy the constraint, though the constraint will be satisfied at convergence.

In principle, our algorithm is capable of enforcing any function of the molecular structure as a constraint on the optimization. Sella natively supports the fixing of arbitrary atomic distances, bending angles, dihedral angles, and one or more Cartesian components of the barycenter of a cluster of atoms. When applied to a single atom, this last constraint fixes the atom in space. Recently, Sella has also implemented the ability to define custom constraint expressions, as well as the ability to enforce inequality constraints.



At each iteration of SQP minimization algorithms, the coordinates are updated with a vector **s** chosen to minimize the quadratic expression<sup>13,35</sup>

$$
\mathbf{s}^T \mathbf{g} + \frac{1}{2} \mathbf{s}^T \mathbf{B}_{\mathcal{L}} \mathbf{s} \tag{21}
$$



Figure 4: An illustration of the null-space SQP method. The initial point q has gradient g (red) and does not lie on the constraint manifold (green). The constraint-correction displacement vector  $s_P$  leads to the closest point on the constraint manifold. The extrapolated gradient  $\mathbf{g}_{ex}$  (orange) is the sum of the original gradient **g** and the approximate gradient  $\mathbf{Bs}_P$ (yellow) created by the displacement assuming constant curvature of the potential energy surface. The optimization displacement  $s_Q$  is calculated with this extrapolated gradient in the space tangent to the constraint manifold and orthogonal to  $s<sub>P</sub>$ . These two displacement vectors are combined into the overall displacement vector s.

subject to the linearized constraint

$$
\mathbf{c}(\mathbf{q}) + \left(\frac{\partial \mathbf{c}}{\partial \mathbf{q}}\right) \mathbf{s} = \mathbf{0},\tag{22}
$$

where  $\mathbf{B}_{\mathcal{L}}$  is the approximate Hessian of the Lagrangian defined in equation 15. This minimization problem can be solved in a variety of ways. Our algorithm uses the nullspace SQP method,<sup>13,35</sup> illustrated in figure 4, which splits the displacement vector  $s$ into two components,

$$
\mathbf{s} = \mathbf{s}_P + \mathbf{s}_Q,\tag{23}
$$

where  $s_P$  is the constraint correction displacement vector and  $s_Q$  is the optimization step displacement vector, and  $s_P \perp s_Q$ . This method was chosen as it reduces the dimensionality of the optimization problem by the size of the space spanned by the constraint equations. The vector  $s_P$  is defined as the least-squares solution to equation

$$
22,
$$

$$
\left(\frac{\partial \mathbf{c}}{\partial \mathbf{q}}\right) \mathbf{s}_P = -\mathbf{c}(\mathbf{q}).\tag{24}
$$

In null-space SQP minimization algorithms, the vector  $s_Q$  is found by minimizing the projected quadratic approximation

$$
\tilde{\mathbf{s}}_Q^T \tilde{\mathbf{g}}_{\text{ex}} + \frac{1}{2} \tilde{\mathbf{s}}_Q^T \tilde{\mathbf{B}}_L \tilde{\mathbf{s}}_Q, \qquad (25)
$$

where

$$
\tilde{\mathbf{g}}_{\text{ex}} = \tilde{\mathbf{g}} + \mathbf{Q}^T \mathbf{B} \mathbf{s}_P \tag{26}
$$

is the gradient which has been first approximately extrapolated to the point  $q + s_P$  and then projected into the unconstrained nonredundant space.

Note that equation 25 is the traditional quadratic approximation to the potential energy surface, but with the gradient  $\tilde{\mathbf{g}}$  replaced by the extrapolated gradient  $\tilde{\mathbf{g}}_{\text{ex}}$  and the approximate Hessian of the potential energy  $\ddot{\mathbf{B}}$  replaced by the approximate Hessian of the Lagrangian  $\tilde{\mathbf{B}}_{\mathcal{L}}$ . Therefore, the procedure for obtaining  $\mathbf{s}_{Q}$  described above is equivalent to standard quasi-Newton minimization with some modifications to the gradient and approximate Hessian. It is possible to perform a constrained minimization using an alternative approach such as the restricted step rational function optimization  $(RS-RFO)$  method by making these same substitutions.  $36-38$  This is useful, as unlike quasi-Newton, RS-RFO guarantees a downhill step even when the approximate Hessian  $\tilde{\mathbf{B}}_{\mathcal{L}}$  is indefinite. Furthermore, making these substitutions in the restricted step partitioned rational function optimization (RS-PRFO) method will result in a constrained saddle point optimization step.  $36-38$  Our algorithm implements constrained saddle point optimization by combining null-space SQP with RS-PRFO in this way.

It is illustrative to explain how null-space SQP can be implemented into RS-RFO before explaining how to implement it in the RS-PRFO. In RS-RFO, the vector  $s_Q$  is chosen to be the minimizer of the rational function

$$
\mu = \frac{\tilde{\mathbf{s}}_Q^T \tilde{\mathbf{g}}_{\text{ex}} + \frac{1}{2} \tilde{\mathbf{s}}_Q^T \tilde{\mathbf{B}}_L \tilde{\mathbf{s}}_Q}{1 + \alpha^{-2} \tilde{\mathbf{s}}_Q^T \tilde{\mathbf{s}}_Q},\tag{27}
$$

where  $\alpha$  is a parameter that controls the magnitude of the displacement. The value of  $\alpha$  is chosen to satisfy our trust region condition, which we describe below. Note that as with the standard SQP minimization method described above, we have replaced the gradient  $\tilde{\mathbf{g}}$  with the extrapolated gradient  $\tilde{\mathbf{g}}_{\text{ex}}$  and the approximate Hessian of the potential energy  $\tilde{B}$  with the approximate Hessian of the Lagrangian  $\tilde{B}_{\mathcal{L}}$ . This will result in a displacement vector  $s_Q$  which is suitable for use in the null-space SQP minimization formalism. The right-hand side of equation 27 can be minimized by solving an eigenvalue problem,  $36$ 

$$
\begin{pmatrix} \alpha^2 \tilde{\mathbf{B}}_{\mathcal{L}} & \alpha \tilde{\mathbf{g}}_{\text{ex}} \\ \alpha \tilde{\mathbf{g}}_{\text{ex}}^T & 0 \end{pmatrix} \begin{pmatrix} \alpha^{-1} \tilde{\mathbf{s}}_{Q} \\ 1 \end{pmatrix} = 2\mu \begin{pmatrix} \alpha^{-1} \tilde{\mathbf{s}}_{Q} \\ 1 \end{pmatrix}, \qquad (28)
$$

where a minimization step is obtained from the leftmost solution to equation 28. In theory, choosing to define  $s_Q$  using the *j*th eigenvector of equation 28 will result in a step towards an order-j saddle point, but in practice this approach often fails to converge.

In the RS-PRFO method,  $36-38$  the unconstrained space  $\tilde{q}$  is split into maximization and minimization subspaces which are defined using the eigenvectors of  $\tilde{\mathbf{B}}_{\mathcal{L}}$ . The goal of RS-PRFO is to find a step that maximizes the energy in the maximization subspace while minimizing the energy in the minimization subspace. To optimize to an order- $j$ saddle point, the *j* leftmost eigenvectors of  $\tilde{\mathbf{B}}_{\mathcal{L}}$  are assigned to the columns of  $\tilde{\mathbf{V}}_{\mathcal{L}}^{(\max)}$  $\mathcal{L}$ and the remaining eigenvectors are assigned to the columns of  $\tilde{V}_{\mathcal{L}}^{(\min)}$ . In particular, optimization to a first-order saddle point will assign only the leftmost eigenvector of  $\tilde{\mathbf{B}}_{\mathcal{L}}$ to  $\tilde{\mathbf{V}}_{\mathcal{L}}^{(\text{max})}$ .  $\tilde{\mathbf{V}}_{\mathcal{L}}^{(\text{max})}$  and  $\tilde{\mathbf{V}}_{\mathcal{L}}^{(\text{min})}$  behave as projection matrices from the unconstrained space to the maximization and minimization subspaces, respectively. The displacement in the maximization subspace is found by calculating the rightmost eigenvector of the eigenvalue problem

$$
\begin{pmatrix}\n\alpha^2 \tilde{\mathbf{B}}_{\mathcal{L}}^{(\max)} & \alpha \tilde{\mathbf{g}}_{\text{ex}}^{(\max)} \\
\alpha \tilde{\mathbf{g}}_{\text{ex}}^{(\max)T} & 0\n\end{pmatrix}\n\begin{pmatrix}\n\alpha^{-1} \tilde{\mathbf{s}}_{Q}^{(\max)} \\
1\n\end{pmatrix} = 2\mu^{(\max)} \begin{pmatrix}\n\alpha^{-1} \tilde{\mathbf{s}}_{Q}^{(\max)} \\
1\n\end{pmatrix},
$$
\n(29)

where  $\tilde{\mathbf{B}}_{\mathcal{L}}^{(\text{max})} = \tilde{\mathbf{V}}_{\mathcal{L}}^{(\text{max})} \tilde{\mathbf{B}}_{\mathcal{L}} \tilde{\mathbf{V}}_{\mathcal{L}}^{(\text{max})}$  and  $\tilde{\mathbf{g}}_{\text{ex}}^{(\text{max})} = \tilde{\mathbf{V}}_{\mathcal{L}}^{(\text{max})} \tilde{\mathbf{g}}_{\text{ex}}$  are the approximate Hessian and extrapolated gradient projected into the maximization subspace, respectively. The step in the *minimization* subspace is found by calculating the *leftmost* eigenvector of the eigenvalue problem

$$
\begin{pmatrix}\n\alpha^2 \tilde{\mathbf{B}}_{\mathcal{L}}^{(\min)} & \alpha \tilde{\mathbf{g}}_{\text{ex}}^{(\min)} \\
\alpha \tilde{\mathbf{g}}_{\text{ex}}^{(\min)T} & 0\n\end{pmatrix}\n\begin{pmatrix}\n\alpha^{-1} \tilde{\mathbf{s}}_{Q}^{(\min)} \\
1\n\end{pmatrix} = 2\mu^{(\min)} \begin{pmatrix}\n\alpha^{-1} \tilde{\mathbf{s}}_{Q}^{(\min)} \\
1\n\end{pmatrix},
$$
\n(30)

where  $\tilde{\mathbf{B}}_{\mathcal{L}}^{(\min)} = \tilde{\mathbf{V}}_{\mathcal{L}}^{(\min)} \tilde{\mathbf{B}}_{\mathcal{L}} \tilde{\mathbf{V}}_{\mathcal{L}}^{(\min)}$  and  $\tilde{\mathbf{g}}_{\mathrm{ex}}^{(\min)} = \tilde{\mathbf{V}}_{\mathcal{L}}^{(\min)T} \tilde{\mathbf{g}}_{\mathrm{ex}}$  are the approximate Hessian and extrapolated gradient projected into the minimization subspace, respectively. The final unconstrained displacement vector is calculated as

$$
\tilde{\mathbf{s}}_Q = \tilde{\mathbf{V}}_L^{(\text{max})} \tilde{\mathbf{s}}_Q^{(\text{max})} + \tilde{\mathbf{V}}_L^{(\text{min})} \tilde{\mathbf{s}}_Q^{(\text{min})},\tag{31}
$$

and the overall displacement vector is obtained according to equation 23.

Once a displacement vector s has been chosen, either for a geometry optimization step or for the iterative Hessian diagonalization scheme described in section 2.3, the molecular geometry must be updated. This process is trivial when using Cartesian coordinates, as the displacement vector can simply be added to the current Cartesian coordinate vector x to update the molecular geometry. In redundant internal coordinates, however, the point in coordinate space specified by  $q + s$  may not correspond to a physically possible combination of internal coordinates. This is because the set of all physically possible coordinate vectors forms a  $(3n-6)$ -dimensional manifold embedded in the larger m-dimensional redundant internal coordinate space. Moreover, this manifold has curvature that arises from the coupling between the internal coordinates. As

a result, even if the starting point q lies on the manifold and the displacement vector s lies tangent to the manifold at  $q$ , the point  $q + s$  will generally not lie on the manifold of physically allowed configurations.

In practice, this problem is usually solved by finding the point  $x'$  in Cartesian space which corresponds to the valid coordinate  $q'$  that is closest to  $q + s$  in redundant internal coordinate space. This can be accomplished by iterating the Newton-Raphson root-finding algorithm until convergence,

$$
\mathbf{x}_{k+1} = \mathbf{x}_k + \mathbf{J}_k^+ \left( \mathbf{q} + \mathbf{s} - \mathbf{q}_k \right),\tag{32}
$$

where  $\mathbf{q}_k$  and  $\mathbf{J}_k$  are the internal coordinate vector and internal–Cartesian Jacobian matrix corresponding to  $\mathbf{x}_k$ , respectively. While this approach is computationally facile and generally converges in only a few iterations, it does not account for the curvature of the space of valid internal coordinates, which results in poor quality displacements when using highly redundant coordinate systems, for example for molecules with rings. We have previously developed an alternate approach for realizing displacements in redundant internal coordinate systems based on geodesics of the curved internal coordinate manifold.<sup>10</sup> This geodesic approach accounts for the correlation between internal coordinates not only at the endpoints of the displacement, but also throughout the entire displacement trajectory. We have shown that this geodesic displacement strategy dramatically reduces the number of steps required to reach convergence in a minimization algorithm as compared to the traditional Newton displacement method. Following each geometry optimization step, the approximate Hessian is updated using the geodesic secant update, which is also described in our previous work. <sup>10</sup>

As mentioned previously, the parameter  $\alpha$  in equation 28 (for minimization) and equations 29 and 30 (for saddle point optimization) is chosen to satisfy a trust region condition. In our trust region algorithm, a displacement vector  $s_k$  is accepted only if it

satisfies the trust region condition,

$$
\|\mathbf{s}_k\|_{\infty} \le \epsilon_k,\tag{33}
$$

where  $\epsilon_k$  is the size of the trust region at iteration k of the optimization. Note that this differs from our previous work in that the infinity-norm is used in the trust region condition, rather than the 2-norm.<sup>1</sup> This is important for optimization in a redundant internal coordinate basis, as using the 2-norm would result in smaller displacements upon addition of new redundant coordinates to the coordinate system. This would unnecessarily slow optimization in systems with a high degree of redundancy in their internal coordinate system, as these redundant coordinates will result in larger displacement vectors in the 2-norm sense, but not necessarily in the infinity-norm sense.

The means by which equation 33 is enforced depends on the relative magnitudes of  $\mathbf{s}_P$  and  $\mathbf{s}_Q$ . If  $\|\mathbf{s}_P\|_{\infty} \geq \epsilon_k$ , then the constraint correction step by itself is too large to fit within the trust region, and thus the displacement step is chosen to be

$$
\mathbf{s}_k = \frac{\epsilon_k}{\|\mathbf{s}_P\|_{\infty}} \mathbf{s}_P,\tag{34}
$$

meaning that no optimization step is taken in the unconstrained space. In this scenario,  $s_Q$  is not used or even calculated at this iteration of the optimization algorithm. If  $\|\mathbf{s}_P\|_{\infty} < \epsilon_k$ , then  $\mathbf{s}_Q$  is calculated with an initial value of  $\alpha = 1$ . If  $\|\mathbf{s}_P + \mathbf{s}_Q\|_{\infty} \leq \epsilon_k$ , then the overall displacement  $s_k = s_P + s_Q$  is accepted. Otherwise, a value of  $\alpha$  is determined for which  $\|\mathbf{s}_P + \mathbf{s}_Q(\alpha)\|_{\infty} = \epsilon_k$  using the Newton-Raphson root-finding algorithm.

After every optimization step, our algorithm checks to see whether the trust region size  $\epsilon_k$  needs to be updated. This is accomplished using a scheme we have previously developed for saddle point optimization in a Cartesian coordinate basis.<sup>1</sup>

## 3 Results

#### 3.1 Sella benchmark

To measure the performance of our internal coordinate saddle point optimization algorithm, we have created a benchmark test set consisting of 500 molecular geometries generated by KinBot. <sup>11</sup> This test set consists of small organic molecules with between 7 and 25 atoms in configurations that approximate saddle points for a set of reaction families. Each molecular geometry in the test set was optimized to a saddle point structure in a basis of internal coordinates which were determined automatically by Sella. Calculations were performed with NWChem at the Hartree–Fock level of theory with the 3-21G basis. 39,40 Open-shell molecules were calculated at the unrestricted Hartree–Fock level of theory instead. This level of theory was chosen for computational expedience and to ensure broad compatibility with the widest range of molecular electronic structure theory packages possible. The molecular geometry optimization was considered converged once the maximum atomic force in Cartesian coordinates dropped below  $0.01 \,\mathrm{eV} \,\mathrm{\AA}^{-1}$  (see appendix B). The initial molecular geometries and a Python script for performing the optimizations are available in the supplementary material.

The overall performance of our optimization algorithm on the 500-molecule benchmark is summarized in figure 5. The left half of this figure is a scatter plot that represents the total number of gradient evaluations required to converge to a first order saddle point plotted against the number of atoms in each molecule of the 500-molecule test set. The right half of figure 5 is a histogram of the number of gradient evaluations across the entire test set. With our chosen optimization parameters, the molecules in this test set require between 19 and 125 gradient evaluations to reach convergence, with an average of 50 gradient evaluations.

The majority of molecules in our benchmark set (81.6 %, green) do not have any linear bending angles at any point along the optimization trajectory. Of the remaining molecules, most  $(9.6\%,$  yellow) have a linear bending angle in the initial geometry, but do not require the addition of a dummy atom. Optimization performance for



Figure 5: A scatter plot and histogram depicting the number of gradient evaluations required to converge to a saddle point against the number of atoms in each molecule in the 500 molecule benchmark set. Green dots correspond to molecules with no linear bending angles, yellow to molecules with at least one linear angle at the beginning of optimization but without the need for dummy atoms, black dots correspond to molecules with at least one linear angle at the beginning of optimization that required a dummy atom, and red dots correspond to molecules that required a new set of internal coordinates be constructed during optimization due to a previously valid angle becoming linear.

these molecules is approximately the same as molecules with no linear bending angles, suggesting our algorithm's approach of replacing linear bending angles with improper dihedrals is highly efficient. Only four of the molecules in this test set  $(0.8\%, \text{ black})$ required the addition of a dummy atom in their initial geometry. There is not enough information available in this benchmark set to indicate how optimization performance is affected by the addition of dummy atoms, but the four molecules with dummy atoms in this test set are not outliers. The remaining molecules  $(8.0\%$ , red) acquired a linear bending angle partway through optimization, and therefore required the internal coordinate system to be reconstructed from scratch. When our algorithm discards and reconstructs the internal coordinate system, the Hessian matrix must also be rediagonalized. This results in an increase in the total number of gradient evaluations required to reach convergence on average for these systems, which can be observed in figure 5, though these points do not appear to be outliers.

It is clear from these data that there is no strong correlation between molecule

size and the number of gradient evaluations required to reach convergence. This is perhaps not surprising, considering the small range of molecule sizes included in this study, though it is not at all a guaranteed property. For optimizations in a basis of redundant internal coordinates, it is crucial to use a trust region condition that scales appropriately with the system size, rather than the dimensionality of the basis. If one were to use a 2-norm based trust region condition, the apparent magnitude of any given displacement will increase if additional redundant coordinates are added to the coordinate system. It is for this reason that we choose to use the infinity-norm for the trust region condition, as displacements are limited by only the largest-changing component, and thus the addition of more redundant coordinates will generally not change the displacement magnitude.

To measure the relative performance of our algorithm, we have optimized the set of 500 molecules using four alternative algorithms. The relative performance is summarized in figure 6, which plots the number of effective gradient evaluations for the alternative algorithms against our algorithm as implemented in Sella with its default parameters. Figure 6a compares the performance of our algorithm using Sella's default parameters versus Sella with full calculation of the Hessian matrix, which is achieved by setting the convergence criterion for the iterative diagonalization routine to be extremely tight. The majority of the structures require more gradient evaluations to converge when the Hessian matrix is calculated in full as compared to the partial diagonalization that is performed using Sella's default parameters. Notably, the red data points, which represent structures for which the internal coordinate system is reconstructed midway through optimization, experience a more significant increase in the number of gradient evaluations compared to other structures. This is to be expected, as the Hessian will be evaluated in full at least twice throughout the course of the optimization. For the data points which lie below the line, full diagonalization actually results in an overall reduction in the number of gradient evaluations required to reach convergence. It is likely that for these structures, the initial Hessian diagonalization failed to fully converge the true leftmost eigenvector of the Hessian matrix. This would



Figure 6: A benchmark comparing the effective number of gradient evaluations of our saddle point optimization algorithm as implemented in Sella using its default parameters against a) Sella with full diagonalization of the Hessian, b) NWChem's Z-matrix saddle point optimization algorithm, c) Q-Chem's redundant internal coordinates saddle point optimization algorithm, and d) Psi4's redundant internal coordinate saddle point optimization (optking). The color scheme is identical to that of figure 5. For Sella, the effective number of gradient evaluations is the number of geometry optimization steps plus the number of Hessian diagonalization iterations, while for NWChem, Q-Chem, and Psi4 it is the number of geometry optimization steps plus a factor of 3n to account for analytical evaluation of the full Hessian.

result in an optimization trajectory which initially moves away from the closest saddle point, thereby slowing convergence and possibly resulting in convergence to a different structure entirely.

Figures 6b, 6c, and 6d compare the performance of our algorithm that that of NWChem,  $39,40$  Q-Chem,  $41$  and Psi $442$  respectively. Each of these programs implements a saddle point optimization algorithm in a basis of internal coordinates that requires full evaluations of the Hessian matrix, either analytically or through finite difference. For these codes, the effective number of gradient evaluations is taken to be the number of geometry optimization steps plus a factor of 3n to account for evaluation of the Hessian matrix in full. We note that the cost of evaluating the Hessian matrix using analytical methods – such as coupled-perturbed Hartree–Fock or density functional perturbation theory – is generally lower than the cost of  $3n$  gradient evaluations. However, when analytical Hessians are not available, the actual cost of evaluating the full Hessian matrix will be at least  $3n$ , rising to  $6n$  if central finite difference is used. Based on the result in figure 6, NWChem, Q-Chem, and Psi4 all show comparable performance to Sella when the Hessian is fully diagonalized. NWChem shows the overall worst performance, likely due to its use of a minimal Z-matrix internal coordinate basis rather than a redundant internal coordinate basis. Nonredundant internal coordinate systems such as the Zmatrix approach necessarily omit some chemically relevant degrees of freedom from the coordinate system, which can have a deleterious impact on optimization performance. As with figure 6a, several of the data points lie below the line, indicating that Sella converges more slowly for these points. This is again likely due to a failure of the diagonalization algorithm to correctly identify the lowest curvature mode, a problem that can be ameliorated by tightening the diagonalization convergence criterion  $\gamma$ .

When bending angles become linear mid-optimization, Q-Chem will by default automatically switch to a Cartesian coordinate optimization in order to avoid crashing. We have disabled this feature in our calculations, as leaving it enabled would result in longer optimization trajectories; instead, these calculations are allowed to crash and are therefore omitted from figure 6c. While this feature improves the robustness of

the optimization algorithm implemented in Q-Chem, it does so at the cost of longer trajectories for these pathological systems. In this scenario, our optimization algorithm will instead automatically reconstruct a new internal coordinate system, rather than falling back on Cartesian optimization. This results in a robust optimization algorithm similar to Q-Chem, without the additional cost of switching to a Cartesian optimization scheme.

We note that while each of the four alternative algorithms resulted in some of the structures failing to converge, we do not consider this to be an appropriate metric for comparison, as the molecules in the test set were specifically chosen to be those that converge in Sella with its default parameters. Consequently, any structures that would fail to converge in Sella are excluded from the test set by construction, and so comparing the number of structures that fail to converge would not be fair. The data points in figure 6 are only those structures which converge in both of the codes that are being compared.



Figure 7: A relaxed dihedral angle scan for a transition state structure. (a) The electronic potential energy of the transition state as a dihedral angle is scanned. A depiction of the molecule is inset with the three bonds that make up the dihedral angle highlighted in orange. (b) The total number of gradient evaluations (green) and the number of gradient evaluations per scan steps (red) versus the number of scan steps. The red dashed line corresponds to theoretical lower bound of 1 gradient evaluations per scan step.

It is common in computational studies of molecules to perform scans over specific

internal coordinates. These scans are a series of constrained optimizations, which are straightforward when the goal is to minimize the energy with respect to the unconstrained degrees of freedom, e.g., when the scan is done on a minimum energy structure. However, special care must be taken for saddle point scans to ensure that the scan trajectory preserves the structure of the reacting region of the molecule and does not fall into a different saddle point geometry. Our algorithm's implementation of null-space SQP saddle point optimization provides an easy and effective way to scan coordinates. Figure 7 depicts a full 360° scan of a dihedral angle of a 5-exo-tet cyclization reaction saddle point. The geometry of the saddle point and relative potential energy of the dihedral scan is shown in figure 7a. The orange-highlighted bonds in figure 7a indicate the dihedral angle that is being scanned, which includes several atoms involved in the reaction corresponding to this saddle point. This scan reveals that the structure found by Sella (the structure at zero degrees) is not the lowest-energy conformer for this reaction. This is not surprising, as our saddle point optimization algorithm tends to converge to the saddle point that is geometrically closest to the initial structure, and this may not be the lowest energy conformer.

Figure 7b shows how the total number of gradient evaluations required to perform the entire scan (in green) and the number of gradient evaluations per scan step (in red) scale as a function of the number of scan steps used. For this scan, between 12 and 360 steps were used, which corresponds to a dihedral angle step size of between 30° and 1°. Though the total number of gradient evaluations tends to increase with the number of steps used in the scan, the relationship is not perfectly linear. One possible explanation for this is that larger steps in the scan will tend to perturb the unconstrained degrees of freedom more, therefore requiring more gradient evaluations to converge. The upper limit of this plot indicates that when 360 1° steps are used, only 1.61 gradient evaluations are required per grid point. This is very close to the theoretical lower limit of 1 gradient evaluation per step (the dashed red line), in which each step of the scan converges in a single gradient evaluation. Such performance is only possible due to our use of null-space SQP constrained optimization, which projects the gradient to

Listing 1: A minimal Sella script for optimizing a molecule to a saddle point in internal coordinates.

```
#!/usr/bin/env python3
from ase.io import read
from ase.calculators.nwchem import NWChem
from sella import Sella
atoms = read('input_structure.xyz')
atoms.calc = NWChem()
```
opt = Sella(atoms, internal=True) opt.run(fmax=0.01)

> constraint-corrected geometry before performing any gradient evaluations (see equation 26). The success of this projection approach depends entirely on the accuracy of the approximate Hessian, and so the performance observed in this test system indicates that the initial iterative Hessian diagonalization and TS-BFGS Hessian updates are highly effective.

#### 3.2 Using Sella

Sella is an open source Python software package that is available on both github.com<sup>3</sup> and from the Python Package Index (PyPI). Additionally, the exact version of Sella used to run these calculations is available from Zenodo.<sup>43</sup>

Listing 1 contains a minimal Python script that uses Sella to perform saddle point optimization on a molecule in an internal coordinate basis. An initial geometry contained in input\_structure.xyz is read using ASE's IO routines. <sup>4</sup> An ASE calculator for NWChem is attached to the resulting ASE Atoms object using all default settings, which will perform a calculation at the  $HF/3-21G$  level of theory. This level of theory can be changed by adding arguments to the NWChem calculator object constructor, and the code used to perform the calculation can be altered by replacing the NWChem class with a different calculator for an alternative code. A Sella optimizer object is then created for the ASE Atoms object, with optimization in internal coordinates requested

explicitly. Finally, the optimization to a saddle point is run with a convergence criteria of a maximum atomic force of  $0.01 \text{ eV} \text{ Å}^{-1}$ .

The optimization algorithm implemented in Sella has several user-tunable hyperparameters that may be passed to the Sella class. One important parameter is the iterative diagonalization convergence criterion  $\gamma$ , which may be set with the keyword argument gamma. By default,  $\gamma = 0.1$ , though tighter convergence may be requested by reducing  $\gamma$ . Reducing  $\gamma$  will result in additional iterative Hessian diagonalization steps, but will result in a more accurate approximate Hessian and therefore better subsequent geometry optimization steps. The initial trust region size  $\epsilon_0$  defaults to a value of 0.1 (with units of Å for bond stretch coordinates and radians for angle coordinates). This parameter may be changed by using the keyword argument delta0. Increasing the trust region size will result in larger step sizes, though at the cost of increasing the likelihood of a poor step. Sella is also capable of performing minimization or optimization to higher-order saddle points. The desired saddle point order may be specified with the order keyword argument. This defaults to 1, corresponding to optimization to a first order saddle point. Minimization can be requrested by setting order to 0. Additional keyword arguments are outlined in the Sella manual, available on github. <sup>3</sup>

Listing 2 contains a more complex example Python script that uses Sella to optimize a molecule to a saddle point. First, this example uses the i-PI socket protocol to communicate with the NWChem executable. <sup>44</sup> This speeds up the overall calculation time by reusing the same NWChem processes for all energy and gradient calculations, rather than starting a new instance of NWChem for each single point. Other electronic structure theory codes that support the i-PI socket protocol include Quantum Espresso,  $45,46$  Siesta,  $47,48$  Abinit,  $49,50$  and FHI-aims.  $51$  Additionally, this example uses the irun iterator for running the optimization, rather than the run function. The use of irun makes it possible to run custom user-specified code at every iteration of the optimization trajectory. This enables users to perform more sophisticated logging, monitor the structure as it is optimizing, or implement their own custom convergence criteria.

Listing 2: A more complex Sella script for optimizing a molecule to a saddle point in internal coordinates.

#!/usr/bin/env python3

```
import numpy as np
from ase.io import read
from ase.calculators.nwchem import NWChem
from ase.calculators.socketio import SocketIOCalculator
from sella import Sella
socket_name = 'nwchem'
atoms = read('input_structure.xyz')
nwchem = NWChem (
    driver={'socket': {'unix': socket_name}},
    task='optimize',
\mathcal{L}with SocketIOCalculator(nwchem, unixsocket=socket_name) as calc:
    atoms.calc = calc
    opt = Sella(atoms, internal=True)
    for i, converged in enumerate(opt.irun(fmax=0)):
        print(f'The energy at iteration {i} is {atoms.get_potential_energy()}')
        if atoms.get_distance(0, 1) > 2.0:
            raise RuntimeError('The 0-1 bond broke!')
        if np.abs(opt.pes.get_g()) . max() < 1e-2:
            break
```
## 4 Conclusion

We have presented a novel algorithm for the automated optimization of molecules to saddle points in a basis of redundant internal coordinates. Our algorithm automatically constructs internal coordinates from the Cartesian positions of the atoms, and automatically replaces pathological linear angles with improper dihedrals. When necessary, our algorithm automatically adds dummy atoms to the system, and adds constraints to ensure that the dummy atom does not drift unnecessarily over the course of optimization. In the event of a bending angle becoming linear mid-optimization, the coordinate system is re-constructed for the current structure with the above procedure, resulting in a more robust algorithm. To determine the direction of the reaction coordinates, our algorithm iteratively diagonalizes the Hessian in a basis of internal coordinates. The potential energy surface curvature information obtained during the iterative diagonalization procedure is used to construct a partially-exact Hessian matrix. This Hessian matrix is used in a state-of-the-art constrained partitioned rational function optimization algorithm to step towards a saddle point geometry. Our approach enables efficient and robust high-throughput automated saddle point optimization of molecular species. The performance of our algorithm was tested on a newly created benchmark set of 500 molecular geometries that are geometrically close to reactive saddle points. We find that on this benchmark, our algorithm converges to saddle point structures in fewer effective steps than that of popular quantum chemistry codes.

## Acknowledgement

This work was done within the Exascale Catalytic Chemistry (ECC) Project, which is supported by the U.S. Department of Energy, Office of Science, Basic Energy Sciences, Chemical Sciences, Geosciences and Biosciences Division, as part of the Computational Chemistry Sciences Program.

We thank Dr. Mohan Sarovar for reviewing this manuscript and for providing several helpful comments regarding its presentation.

Sandia National Laboratories is a multimission laboratory managed and operated by National Technology and Engineering Solutions of Sandia, LLC., a wholly owned subsidiary of Honeywell International, Inc., for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA0003525. The views expressed in the article do not necessarily represent the views of the U.S. Department of Energy or the United States Government.

## Supporting Information Available

The 500 initial molecular geometries used in figure 5 and an example Python script for optimizing those structures to a saddle point using Sella are available in the SI.

### A Automatic Generation of Internal Coordinates

Our approach for finding internal coordinates, illustrated in figure 8, begins by constructing a molecular graph from the initial Cartesian geometry. This is accomplished by checking all pairs of atoms, and adding a bond if

$$
\|\mathbf{x}_j - \mathbf{x}_i\| < \sigma \left( r_i^{\text{cov}} + r_j^{\text{cov}} \right),\tag{35}
$$

where  $x_i$  and  $x_j$  are the Cartesian position of atoms i and j,  $\sigma$  is a scalar parameter which in Sella defaults to 1.25, and  $r_i^{\text{cov}}$  and  $r_j^{\text{cov}}$  are the covalent radii of atoms i and  $j$ .<sup>52</sup> After all pairs of atoms have been checked, a flood fill algorithm is used to count the number of disconnected fragments in the molecular graph. This algorithm works by walking along the graph and noting which atoms belong to the same fragment, then counting the number of distinct fragments. If only a single fragment is found, the algorithm exits; otherwise,  $\sigma$  is increased by 5% and the above procedure is repeated with the additional rule that only pairs of atoms in different molecular fragments are checked for connectivity. This additional rule helps ensure that geometries with large separation between clusters of atoms do not become densely connected. This procedure



Figure 8: An illustration of our algorithm for identifying internal coordinates. (a) Spheres are drawn around each atom with radius equal to the covalent radius of the element multiplied by a scaling factor. Bonds are added between atoms with overlapping spheres (yellow regions). (b) The molecular graph obtained from this procedure. Nodes of the graph correspond to atoms, while edges (yellow) correspond to bond stretch coordinates. Angle bending coordinates (green) are defined by considering all possible pairs of edges sharing exactly one node. (c) Dihedral angle coordinates (red) are defined by considering all possible pairs of bending angle coordinates sharing exactly one edge, excluding those that would form an improper dihedral.

continues until only a single molecular fragment is found.

Once the molecular graph has been constructed, the internal coordinates are formed from combinations of the edges of the graph. <sup>8</sup> Each unique graph edge corresponds to a bond stretching coordinate. Pairs of bond stretching coordinates that share exactly one atom are combined to form angle bending coordinates, with the common atom at the center of the angle. Pairs of angle bending coordinates that share exactly one edge are combined to form dihedral angles (pairs of angles that would form improper dihedrals, such as e.g. an  $H(1)-N-H(2)-H(3)$  dihedral angle in ammonia, are not added automatically).

For systems composed of independent molecules, it may not be chemically meaningful to consider internal coordinates that connect two molecules together. However, the previously described procedure will necessarily add bonds, angles, and dihedrals that span across molecules. This can have a negative impact on performance, as the optimizer will limit the extent to which these chemically irrelevant coordinates can change at each step. To resolve this problem, we also implement as an alternative the translational and rotation internal coordinates (TRIC) method of Wang and Song. <sup>53</sup> When using the TRIC method, our algorithm will stop adding bonds between atoms after a single iteration of the procedure described above. Each disconnected fragment of the molecular graph is then augmented with translational coordinates of the fragment barycenter and rotational coordinates following the TRIC procedure.

## B Geometry Optimizer Convergence Criterion

For consistency with ASE's built-in optimizer classes, Sella's optimization algorithm uses the maximum atomic force as the default convergence criterion. Using this criterion, convergence is achieved once the largest Cartesian force experienced by any atom in the molecule drops below a threshold value,

$$
\left\| \frac{d\mathcal{E}}{d\mathbf{x}_i} \right\|_2 \le f_{\text{max}} \quad \text{for all} \quad i = 1, \dots, n,
$$
 (36)

where  $x_i$  represents the position of atom i. When using constraints, it is necessary to first project out the components of the gradient in the constraint subspace. This is accomplished by projecting the gradient in the unconstrained space  $\tilde{g}$  back to Cartesian coordinates,

$$
\left(\frac{d\mathcal{E}}{d\mathbf{x}}\right)_{\text{proj}} = \mathbf{J}^T \mathbf{Q} \tilde{\mathbf{g}}.
$$
\n(37)

## References

- (1) Hermes, E. D.; Sargsyan, K.; Najm, H. N.; Zádor, J. Accelerated Saddle Point Refinement through Full Exploitation of Partial Hessian Diagonalization. J. Chem. Theory Comput. 2019, 15, 6536–6549.
- (2) Chill, S. T.; Stevenson, J.; Ruehle, V.; Shang, C.; Xiao, P.; Farrell, J. D.; Wales, D. J.; Henkelman, G. Benchmarks for Characterization of Minima, Transition States, and Pathways in Atomic, Molecular, and Condensed Matter Systems. J. Chem. Theory Comput. 2014, 10, 5476–5482.
- (3) Sella. 2022; https://github.com/zadorlab/sella, original-date: 2019-04- 09T15:16:08Z.
- (4) Larsen, A. H.; Mortensen, J. J.; Blomqvist, J.; Castelli, I. E.; Christensen, R.; Dułak, M.; Friis, J.; Groves, M. N.; Hammer, B.; Hargus, C.; Hermes, E. D.; Jennings, P. C.; Jensen, P. B.; Kermode, J.; Kitchin, J. R.; Kolsbjerg, E. L.; Kubal, J.; Kaasbjerg, K.; Lysgaard, S.; Maronsson, J. B.; Maxson, T.; Olsen, T.; Pastewka, L.; Peterson, A.; Rostgaard, C.; Schiøtz, J.; Schütt, O.; Strange, M.; Thygesen, K. S.; Vegge, T.; Vilhelmsen, L.; Walter, M.; Zeng, Z.; Jacobsen, K. W. The atomic simulation environment—a Python library for working with atoms. J. Phys.: Condens. Matter 2017, 29, 273002.
- (5) Schlegel, H. B. A comparison of geometry optimization with internal, cartesian, and mixed coordinates. Int. J. Quantum Chem.  $1992, 44, 243-252$ .
- (6) Pulay, P.; Fogarasi, G. Geometry optimization in redundant internal coordinates. J. Chem. Phys. 1992, 96, 2856–2860.
- (7) Baker, J.; Kessi, A.; Delley, B. The generation and use of delocalized internal coordinates in geometry optimization. J. Chem. Phys. 1996, 105, 192–212.
- (8) Peng, C.; Ayala, P. Y.; Schlegel, H. B.; Frisch, M. J. Using redundant internal coordinates to optimize equilibrium geometries and transition states. J. Comput. Chem. 1996, 17, 49–56.
- (9) Bakken, V.; Helgaker, T. The efficient optimization of molecular geometries using redundant internal coordinates. J. Chem. Phys. 2002, 117, 9160–9174, Publisher: American Institute of Physics.
- (10) Hermes, E. D.; Sargsyan, K.; Najm, H. N.; Zádor, J. Geometry optimization speedup through a geodesic approach to internal coordinates. J. Chem. Phys. 2021, 155, 094105, Publisher: American Institute of Physics.
- (11) Van de Vijver, R.; Zádor, J. KinBot: Automated stationary point search on potential energy surfaces. Comput. Phys. Commun. 2020, 248, 106947.
- (12) Wilson, R. B. A simplicial algorithm for concave programming. Ph. D. Dissertation, Graduate School of Business Administration 1963,
- (13) Nocedal, J.; Wright, S. J. In Numerical Optimization; Glynn, P., Robinson, S. M., Eds.; Springer Series in Operations Research; Springer-Verlag New York, Inc.: 175 Fifth Avenue, New York, NY 10010, USA, 1999.
- (14) Birkholz, A. B.; Schlegel, H. B. Exploration of some refinements to geometry optimization methods. Theor. Chem. Acc. 2016, 135, 84.
- (15) Wilson, E.; Decius, J.; Cross, P. Molecular Vibrations: The Theory of Infrared and Raman Vibrational Spectra; Dover Books on Chemistry Series; Dover Publications, 1980.
- (16) Moore, E. H. On the reciprocal of the general algebraic matrix. Bull. Am. Math. Soc. **1920**, 26, 394–395.
- (17) Penrose, R. A generalized inverse for matrices. Math. Proc. Cambridge Philos. Soc. 1955, 51, 406–413, Publisher: Cambridge University Press.
- (18) Sharada, S. M.; Bell, A. T.; Head-Gordon, M. A finite difference Davidson procedure to sidestep full ab initio hessian calculation: Application to characterization of stationary points and transition state searches. J. Chem. Phys. 2014,  $140$ , 164115, Publisher: American Institute of Physics.
- (19) Baker, J.; Bergeron, D. Constrained optimization in cartesian coordinates. J. Comput. Chem. 1993, 14, 1339–1346.
- (20) Baker, J. Constrained optimization in delocalized internal coordinates. J. Comput. Chem. 1997, 18, 1079–1095.
- (21) Olsen, J.; Jørgensen, P.; Simons, J. Passing the one-billion limit in full configuration-interaction (FCI) calculations. Chem. Phys. Lett. 1990, 169, 463– 472.
- (22) Saad, Y.; Lehoucq, R.; Sorensen, D. Templates for the Solution of Algebraic Eigenvalue Problems; Software, Environments and Tools; Society for Industrial and Applied Mathematics, 2000; pp 37–44.
- (23) Davidson, E. R. The iterative calculation of a few of the lowest eigenvalues and corresponding eigenvectors of large real-symmetric matrices. J. Comput. Phys. 1975, 17, 87–94.
- (24) Morgan, R.; Scott, D. Generalizations of Davidson's Method for Computing Eigenvalues of Sparse Symmetric Matrices. Siam J. Sci. Stat. Comput. 1986, 7, 817–825.
- (25) G. Sleijpen, G.; Van der Vorst, H. A Jacobi–Davidson Iteration Method for Linear Eigenvalue Problems. SIAM J. Matrix Anal. Appl. 1996, 17, 401–425.
- (26) Sleijpen, G. L. G.; Wubs, F. W. Effective preconditioning techniques for eigenvalue problems. Preprint 1117 1999,
- (27) Stathopoulos, A. Nearly Optimal Preconditioned Methods for Hermitian Eigenproblems under Limited Memory. Part I: Seeking One Eigenvalue. Siam J. Sci. Comput. 2007, 29, 481–514.
- (28) Schnabel, R. B. Quasi-Newton Methods Using Multiple Secant Equations ; CU-CS-247-83. Computer Science Technical Reports 1983, 224, 1–40.
- (29) Anglada, J. M.; Bofill, J. M. How good is a Broyden–Fletcher–Goldfarb–Shannolike update Hessian formula to locate transition structures? Specific reformulation of Broyden–Fletcher–Goldfarb–Shanno for optimizing saddle points. J. Comput. Chem. 1998, 19, 349–362.
- (30) Bofill, J. M. Remarks on the updated Hessian matrix methods. Int. J. Quantum Chem. 2003, 94, 324–332.
- (31) Gower, R. M.; Gondzio, J. Action constrained quasi-Newton methods.  $arXiv:1412.8045$  [cs, math] 2014, arXiv: 1412.8045.
- (32) Fischer, T. H.; Almlof, J. General methods for geometry and wave function optimization. J. Phys. Chem. 1992, 96, 9768–9774.
- (33) Ryckaert, J.-P.; Ciccotti, G.; Berendsen, H. J. C. Numerical integration of the cartesian equations of motion of a system with constraints: molecular dynamics of n-alkanes. J. Comput. Phys. 1977, 23, 327–341.
- (34) Andersen, H. C. Rattle: A "velocity" version of the shake algorithm for molecular dynamics calculations. J. Comput. Phys. 1983, 52, 24–34.
- (35) Vasantharajan, S.; Biegler, L. T. Large-scale decomposition for successive quadratic programming. Comput. Chem. Eng. 1988, 12, 1087–1101.
- (36) Banerjee, A.; Adams, N.; Simons, J.; Shepard, R. Search for stationary points on surfaces. J. Phys. Chem. 1985, 89, 52–57.
- (37) Anglada, J. M.; Bofill, J. M. A reduced-restricted-quasi-Newton–Raphson method for locating and optimizing energy crossing points between two potential energy surfaces. J. Comput. Chem. 1997, 18, 992–1003.
- (38) Besalú, E.; Bofill, J. M. On the automatic restricted-step rational-functionoptimization method. Theor. Chem. Acc. 1998, 100, 265–274.
- (39) Valiev, M.; Bylaska, E. J.; Govind, N.; Kowalski, K.; Straatsma, T. P.; Van Dam, H. J. J.; Wang, D.; Nieplocha, J.; Apra, E.; Windus, T. L.; de Jong, W. A. NWChem: A comprehensive and scalable open-source solution for large scale molecular simulations. Comput. Phys. Commun. 2010, 181, 1477– 1489.
- (40) Aprà, E.; Bylaska, E. J.; de Jong, W. A.; Govind, N.; Kowalski, K.; Straatsma, T. P.; Valiev, M.; van Dam, H. J. J.; Alexeev, Y.; Anchell, J.; Anisimov, V.; Aquino, F. W.; Atta-Fynn, R.; Autschbach, J.; Bauman, N. P.; Becca, J. C.; Bernholdt, D. E.; Bhaskaran-Nair, K.; Bogatko, S.; Borowski, P.; Boschen, J.; Brabec, J.; Bruner, A.; Cauët, E.; Chen, Y.; Chuev, G. N.; Cramer, C. J.; Daily, J.; Deegan, M. J. O.; Dunning, T. H.; Dupuis, M.; Dyall, K. G.; Fann, G. I.; Fischer, S. A.; Fonari, A.; Früchtl, H.; Gagliardi, L.; Garza, J.; Gawande, N.; Ghosh, S.; Glaesemann, K.; Götz, A. W.; Hammond, J.; Helms, V.; Hermes, E. D.; Hirao, K.; Hirata, S.; Jacquelin, M.; Jensen, L.; Johnson, B. G.; Jónsson, H.; Kendall, R. A.; Klemm, M.; Kobayashi, R.; Konkov, V.; Krishnamoorthy, S.; Krishnan, M.; Lin, Z.; Lins, R. D.; Littlefield, R. J.; Logsdail, A. J.; Lopata, K.; Ma, W.; Marenich, A. V.; Martin del Campo, J.; Mejia-Rodriguez, D.; Moore, J. E.; Mullin, J. M.; Nakajima, T.; Nascimento, D. R.;

Nichols, J. A.; Nichols, P. J.; Nieplocha, J.; Otero-de-la Roza, A.; Palmer, B.; Panyala, A.; Pirojsirikul, T.; Peng, B.; Peverati, R.; Pittner, J.; Pollack, L.; Richard, R. M.; Sadayappan, P.; Schatz, G. C.; Shelton, W. A.; Silverstein, D. W.; Smith, D. M. A.; Soares, T. A.; Song, D.; Swart, M.; Taylor, H. L.; Thomas, G. S.; Tipparaju, V.; Truhlar, D. G.; Tsemekhman, K.; Van Voorhis, T.; Vázquez-Mayagoitia, Á.; Verma, P.; Villa, O.; Vishnu, A.; Vogiatzis, K. D.; Wang, D.; Weare, J. H.; Williamson, M. J.; Windus, T. L.; Woliński, K.; Wong, A. T.; Wu, Q.; Yang, C.; Yu, Q.; Zacharias, M.; Zhang, Z.; Zhao, Y.; Harrison, R. J. NWChem: Past, present, and future. J. Chem. Phys. 2020, 152, 184102, Publisher: American Institute of Physics.

(41) Epifanovsky, E.; Gilbert, A. T. B.; Feng, X.; Lee, J.; Mao, Y.; Mardirossian, N.; Pokhilko, P.; White, A. F.; Coons, M. P.; Dempwolff, A. L.; Gan, Z.; Hait, D.; Horn, P. R.; Jacobson, L. D.; Kaliman, I.; Kussmann, J.; Lange, A. W.; Lao, K. U.; Levine, D. S.; Liu, J.; McKenzie, S. C.; Morrison, A. F.; Nanda, K. D.; Plasser, F.; Rehn, D. R.; Vidal, M. L.; You, Z.-Q.; Zhu, Y.; Alam, B.; Albrecht, B. J.; Aldossary, A.; Alguire, E.; Andersen, J. H.; Athavale, V.; Barton, D.; Begam, K.; Behn, A.; Bellonzi, N.; Bernard, Y. A.; Berquist, E. J.; Burton, H. G. A.; Carreras, A.; Carter-Fenk, K.; Chakraborty, R.; Chien, A. D.; Closser, K. D.; Cofer-Shabica, V.; Dasgupta, S.; de Wergifosse, M.; Deng, J.; Diedenhofen, M.; Do, H.; Ehlert, S.; Fang, P.-T.; Fatehi, S.; Feng, Q.; Friedhoff, T.; Gayvert, J.; Ge, Q.; Gidofalvi, G.; Goldey, M.; Gomes, J.; González-Espinoza, C. E.; Gulania, S.; Gunina, A. O.; Hanson-Heine, M. W. D.; Harbach, P. H. P.; Hauser, A.; Herbst, M. F.; Hernández Vera, M.; Hodecker, M.; Holden, Z. C.; Houck, S.; Huang, X.; Hui, K.; Huynh, B. C.; Ivanov, M.; Jász, A.; Ji, H.; Jiang, H.; Kaduk, B.; Kähler, S.; Khistyaev, K.; Kim, J.; Kis, G.; Klunzinger, P.; Koczor-Benda, Z.; Koh, J. H.; Kosenkov, D.; Koulias, L.; Kowalczyk, T.; Krauter, C. M.; Kue, K.; Kunitsa, A.; Kus, T.; Ladjánszki, I.; Landau, A.; Lawler, K. V.; Lefrancois, D.; Lehtola, S.; Li, R. R.; Li, Y.-P.; Liang, J.; Liebenthal, M.; Lin, H.-H.; Lin, Y.-S.; Liu, F.; Liu, K.-Y.; Loipersberger, M.; Luenser, A.; Manjanath, A.; Manohar, P.; Mansoor, E.; Manzer, S. F.; Mao, S.-P.; Marenich, A. V.; Markovich, T.; Mason, S.; Maurer, S. A.; McLaughlin, P. F.; Menger, M. F. S. J.; Mewes, J.-M.; Mewes, S. A.; Morgante, P.; Mullinax, J. W.; Oosterbaan, K. J.; Paran, G.; Paul, A. C.; Paul, S. K.; Pavošević, F.; Pei, Z.; Prager, S.; Proynov, E. I.; Rák, A.; Ramos-Cordoba, E.; Rana, B.; Rask, A. E.; Rettig, A.; Richard, R. M.; Rob, F.; Rossomme, E.; Scheele, T.; Scheurer, M.; Schneider, M.; Sergueev, N.; Sharada, S. M.; Skomorowski, W.; Small, D. W.; Stein, C. J.; Su, Y.-C.; Sundstrom, E. J.; Tao, Z.; Thirman, J.; Tornai, G. J.; Tsuchimochi, T.; Tubman, N. M.; Veccham, S. P.; Vydrov, O.; Wenzel, J.; Witte, J.; Yamada, A.; Yao, K.; Yeganeh, S.; Yost, S. R.; Zech, A.; Zhang, I. Y.; Zhang, X.; Zhang, Y.; Zuev, D.; Aspuru-Guzik, A.; Bell, A. T.; Besley, N. A.; Bravaya, K. B.; Brooks, B. R.; Casanova, D.; Chai, J.-D.; Coriani, S.; Cramer, C. J.; Cserey, G.; DePrince, A. E.; DiStasio, R. A.; Dreuw, A.; Dunietz, B. D.; Furlani, T. R.; Goddard, W. A.; Hammes-Schiffer, S.; Head-Gordon, T.; Hehre, W. J.; Hsu, C.-P.; Jagau, T.-C.; Jung, Y.; Klamt, A.; Kong, J.; Lambrecht, D. S.; Liang, W.; Mayhall, N. J.; McCurdy, C. W.; Neaton, J. B.; Ochsenfeld, C.; Parkhill, J. A.; Peverati, R.; Rassolov, V. A.; Shao, Y.; Slipchenko, L. V.; Stauch, T.; Steele, R. P.; Subotnik, J. E.; Thom, A. J. W.; Tkatchenko, A.; Truhlar, D. G.; Van Voorhis, T.; Wesolowski, T. A.; Whaley, K. B.; Woodcock, H. L.; Zimmerman, P. M.; Faraji, S.; Gill, P. M. W.; Head-Gordon, M.; Herbert, J. M.; Krylov, A. I. Software for the frontiers of quantum chemistry: An overview of developments in the Q-Chem 5 package. J. Chem. Phys. 2021, 155, 084801, Publisher: American Institute of Physics.

(42) Turney, J. M.; Simmonett, A. C.; Parrish, R. M.; Hohenstein, E. G.; Evangelista, F. A.; Fermann, J. T.; Mintz, B. J.; Burns, L. A.; Wilke, J. J.; Abrams, M. L.; Russ, N. J.; Leininger, M. L.; Janssen, C. L.; Seidl, E. T.; Allen, W. D.; Schaefer, H. F.; King, R. A.; Valeev, E. F.; Sherrill, C. D.; Crawford, T. D. Psi4: an open-source ab initio electronic structure program. WIREs Comput. Mol. Sci. 2012,  $2, 556-565$ , eprint:

https://onlinelibrary.wiley.com/doi/pdf/10.1002/wcms.93.

- (43) Hermes, E. Sella. 2022; https://doi.org/10.5281/zenodo.6459455.
- (44) Ceriotti, M.; More, J.; Manolopoulos, D. E. i-PI: A Python interface for ab initio path integral molecular dynamics simulations. Comput. Phys. Commun. 2014, 185, 1019–1026.
- (45) Giannozzi, P.; Baroni, S.; Bonini, N.; Calandra, M.; Car, R.; Cavazzoni, C.; Ceresoli, D.; Chiarotti, G. L.; Cococcioni, M.; Dabo, I.; Corso, A. D.; Gironcoli, S. d.; Fabris, S.; Fratesi, G.; Gebauer, R.; Gerstmann, U.; Gougoussis, C.; Kokalj, A.; Lazzeri, M.; Martin-Samos, L.; Marzari, N.; Mauri, F.; Mazzarello, R.; Paolini, S.; Pasquarello, A.; Paulatto, L.; Sbraccia, C.; Scandolo, S.; Sclauzero, G.; Seitsonen, A. P.; Smogunov, A.; Umari, P.; Wentzcovitch, R. M. QUANTUM ESPRESSO: a modular and open-source software project for quantum simulations of materials. J. Phys.: Condens. Matter 2009, 21, 395502.
- (46) Giannozzi, P.; Andreussi, O.; Brumme, T.; Bunau, O.; Nardelli, M. B.; Calandra, M.; Car, R.; Cavazzoni, C.; Ceresoli, D.; Cococcioni, M.; Colonna, N.; Carnimeo, I.; Corso, A. D.; Gironcoli, S. d.; Delugas, P.; DiStasio, R. A.; Ferretti, A.; Floris, A.; Fratesi, G.; Fugallo, G.; Gebauer, R.; Gerstmann, U.; Giustino, F.; Gorni, T.; Jia, J.; Kawamura, M.; Ko, H.-Y.; Kokalj, A.; Küçükbenli, E.; Lazzeri, M.; Marsili, M.; Marzari, N.; Mauri, F.; Nguyen, N. L.; Nguyen, H.-V.; Otero-de-la Roza, A.; Paulatto, L.; Poncé, S.; Rocca, D.; Sabatini, R.; Santra, B.; Schlipf, M.; Seitsonen, A. P.; Smogunov, A.; Timrov, I.; Thonhauser, T.; Umari, P.; Vast, N.; Wu, X.; Baroni, S. Advanced capabilities for materials modelling with Quantum ESPRESSO. J. Phys.: Condens. Matter 2017, 29, 465901.
- (47) Soler, J. M.; Artacho, E.; Gale, J. D.; García, A.; Junquera, J.; Ordejón, P.; Sánchez-Portal, D. The SIESTA method for ab initio order-N materials simulation. J. Phys.: Condens. Matter 2002, 14, 2745–2779, Publisher: IOP Publishing.
- (48) García, A.; Papior, N.; Akhtar, A.; Artacho, E.; Blum, V.; Bosoni, E.; Brandimarte, P.; Brandbyge, M.; Cerdá, J. I.; Corsetti, F.; Cuadrado, R.; Dikan, V.; Ferrer, J.; Gale, J.; García-Fernández, P.; García-Suárez, V. M.; García, S.; Huhs, G.; Illera, S.; Korytár, R.; Koval, P.; Lebedeva, I.; Lin, L.; López-Tarifa, P.; Mayo, S. G.; Mohr, S.; Ordejón, P.; Postnikov, A.; Pouillon, Y.; Pruneda, M.; Robles, R.; Sánchez-Portal, D.; Soler, J. M.; Ullah, R.; Yu, V. W.-z.; Junquera, J. Siesta: Recent developments and applications. J. Chem. Phys. 2020, 152, 204108, Publisher: American Institute of Physics.
- (49) Gonze, X.; Amadon, B.; Antonius, G.; Arnardi, F.; Baguet, L.; Beuken, J.-M.; Bieder, J.; Bottin, F.; Bouchet, J.; Bousquet, E.; Brouwer, N.; Bruneval, F.; Brunin, G.; Cavignac, T.; Charraud, J.-B.; Chen, W.; Côté, M.; Cottenier, S.; Denier, J.; Geneste, G.; Ghosez, P.; Giantomassi, M.; Gillet, Y.; Gingras, O.; Hamann, D. R.; Hautier, G.; He, X.; Helbig, N.; Holzwarth, N.; Jia, Y.; Jollet, F.; Lafargue-Dit-Hauret, W.; Lejaeghere, K.; Marques, M. A. L.; Martin, A.; Martins, C.; Miranda, H. P. C.; Naccarato, F.; Persson, K.; Petretto, G.; Planes, V.; Pouillon, Y.; Prokhorenko, S.; Ricci, F.; Rignanese, G.-M.; Romero, A. H.; Schmitt, M. M.; Torrent, M.; van Setten, M. J.; Van Troeye, B.; Verstraete, M. J.; Zérah, G.; Zwanziger, J. W. The Abinit project: Impact, environment and recent developments. Comput. Phys. Commun. 2020, 248, 107042.
- (50) Romero, A. H.; Allan, D. C.; Amadon, B.; Antonius, G.; Applencourt, T.; Baguet, L.; Bieder, J.; Bottin, F.; Bouchet, J.; Bousquet, E.; Bruneval, F.; Brunin, G.; Caliste, D.; Côté, M.; Denier, J.; Dreyer, C.; Ghosez, P.; Giantomassi, M.; Gillet, Y.; Gingras, O.; Hamann, D. R.; Hautier, G.; Jollet, F.; Jomard, G.; Martin, A.; Miranda, H. P. C.; Naccarato, F.; Petretto, G.; Pike, N. A.; Planes, V.; Prokhorenko, S.; Rangel, T.; Ricci, F.; Rignanese, G.-M.; Royo, M.; Stengel, M.; Torrent, M.; van Setten, M. J.; Van Troeye, B.; Verstraete, M. J.; Wiktor, J.; Zwanziger, J. W.; Gonze, X. ABINIT: Overview and focus on selected

capabilities. J. Chem. Phys. 2020, 152, 124102, Publisher: American Institute of Physics.

- (51) Blum, V.; Gehrke, R.; Hanke, F.; Havu, P.; Havu, V.; Ren, X.; Reuter, K.; Scheffler, M. Ab initio molecular simulations with numeric atom-centered orbitals. Comput. Phys. Commun. 2009, 180, 2175–2196.
- (52) Cordero, B.; Gómez, V.; Platero-Prats, A. E.; Revés, M.; Echeverría, J.; Cremades, E.; Barragán, F.; Alvarez, S. Covalent radii revisited. Dalton Trans. 2008, 2832–2838, Publisher: The Royal Society of Chemistry.
- (53) Wang, L.-P.; Song, C. Geometry optimization made simple with translation and rotation coordinates. J. Chem. Phys. 2016, 144, 214108, Publisher: American Institute of Physics.

## TOC Graphic



A Venn diagram illustrating the overlap of three major design features of Sella: saddle point optimization, geodesic internal coordinates, and full automatability of the code.