

Title

Chemical Bonding Origin of Glass Formation in Metal-Organic Frameworks

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Abstract

Glass-forming metal organic frameworks (MOFs) have novel applications, but the origin of their peculiar melting behavior is unclear. Here, we report synchrotron X-ray diffraction electron densities of two zeolitic imidazolate frameworks (ZIFs), the glass-forming Zn-ZIF-zni and the isostructural thermally decomposing Co-ZIF-zni. Electron density analysis shows that the Zn–N bonds are more ionic than Co–N bond, which have distinct covalent features. Variable temperature Raman spectra reveal onset of significant imidazolate bond weakening in Co-ZIF-zni above 673 K. Melting can be controlled by tuning the metal–ligand and imidazole bonding strength as shown from nine solid solution structures of $\text{Co}_x\text{Zn}_{1-x}\text{-ZIF-zni}$ ($x = 0.3$ to 0.003) were synthesized, and a mere 4% Co-doping into Zn-ZIF-zni results in thermal decomposition instead of melting. These findings advocate significant chemical space for designing new MOF glasses either by engineering the chemical bonding with ligand functionalization or by avoiding defects at the metal site in the crystal lattice.

Main text

Metal organic framework (MOF) materials are usually preferred to be crystalline for potential applications to facilitate control over pore size or ligand functionalization^{1, 2}. However, amorphous MOFs³ are increasingly targeted in application e.g. for drug delivery,⁴ packing and delivery of enzymes,⁵ or hydro-

carbon adsorption⁶ due to abundance of active sites, lack of grain boundaries, isotropic mechanical response and better flexibility for shape-selective adsorptions. Pioneering work on amorphous MOFs has especially concerned the MOF family of zeolitic imidazolate frameworks (ZIFs)^{7, 8}. The network topology of ZIFs consists of tetrahedral metal centres (e.g. Zn(II), Co(II)) connected by imidazolate linkers, which replicates the building block of corner sharing SiO₄ tetrahedra in zeolites, Fig. 1. Contrary to the majority of MOFs that show thermal decomposition upon heating, several members of the ZIF family melt above 400 °C,⁹ and they form glasses when thermally quenched from the molten state. Glass-forming ZIFs have opened up new avenues of fundamental and applied research¹⁰. Pure ZIF glasses as well as fabricated glass membranes with intrinsic permanent porosity are used for gas mixture separations (H₂/CH₄, CO₂/N₂, and CO₂/CH₄, propane, propylene, and n-butane)^{11, 12} and CO₂ adsorption¹³. ZIF glasses are also being explored as candidates for optical luminescence materials,¹⁴ non-linear optical materials,¹⁵ or quasi-solid-state-electrolytes¹⁶. Notably, ZIF glasses offer better processability for bulk application than crystalline powders¹⁷.

The archetypical ZIF-4 material containing Zn²⁺ ions was the first glass-forming MOF¹⁸. Coudert *et al.* used first-principle molecular dynamics simulation and a combination of neutron and X-ray pair distribution function analysis to demonstrate that melting of Zn-ZIF-4 involves rapid fluctuations between three-coordinated and four-coordinated Zn-imidazolate units at picosecond time scale¹⁹. Interestingly, the isostructural Co-ZIF-4 material decomposes without melting under heat treatment. Zn-ZIF-4 and Co-ZIF-4 have similar thermal signatures when heated in argon atmosphere with solvent loss from the lattice around 220 °C, followed by amorphization above 300°C, and eventual recrystallization at about 370-400 °C to densest ZIF polymorph, the so called ZIF-zni²⁰. Melting of Zn-ZIF-zni is observed at 590 °C followed by decomposition above 600 °C (Fig. S1a). In contrast, Co-ZIF-zni decomposes at 530 °C (Fig. S1b). It is puzzling why these two isostructural ZIFs, Zn-ZIF-zni and Co-ZIF-zni, behave

fundamentally different upon heating. They differ only in the electronic configuration of their transition metal center with closed shell d^{10} for Zn(II) and open shell d^7 for Co(II). This suggests that the thermal stability of individual MOFs and hence their glass-forming ability is correlated to the nature of the metal–ligand bonding.

Various empirical design strategies based on mixed-ligand approach have been pursued to lower the melting point temperature of ZIF-62 (and derivatives), whose structures are isomorphous to ZIF-4.^{12, 21, 22} However, to date, the physicochemical origin of the melting behaviors remains poorly understood. When addressing the correlation between the melting/decomposition process and the metal–ligand bonding in Zn-ZIF/Co-ZIF, two questions appear important: i) How does the nature of Zn–ligand bond compare to the Co–ligand bond?, and ii) What is the impact of this difference on the structural dynamics and bond strength (stability) of imidazolate linkers in different ZIFs. Here we address the questions based on analysis of electron density (ED) distributions obtained from modelling of very accurate 25 K synchrotron X-ray diffraction data measured to a high resolution of $\sin\theta/\lambda = 1.35 \text{ \AA}^{-1}$ and 1.25 \AA^{-1} for Zn-ZIF-zni and Co-ZIF-zni respectively. ED analysis provides detailed information on chemical bonding and it has helped to form understanding of structure-property relationships in material research²³. However, ED studies of MOF materials remain a severe challenge both for experimental (X-ray diffraction) and computational (*ab initio*) methods. The possible presence of disordered solvents in the pores, framework disorder or large atomic displacements in the flexible linkers hinders accurate modelling of X-ray structure factors, while large open shell extended molecular systems with magnetic ordering are computationally costly and suffer convergence issues. In a few cases, use of intense high-energy synchrotron X-ray radiation in combination with area detector technology has resulted in experimental ED studies of MOFs^{24, 25, 26}.

We have determined accurate EDs of Zn-ZIF-zni and Co-ZIF-zni based on multipole modelling²⁷ of 25 K single crystal synchrotron X-ray diffraction data supplemented with periodic DFT calculations. ZIF-zni, the highest density polymorph was selected for this study because its crystal structure is devoid of solvents, thus avoiding potential structural disorder. Similar to ZIF-4, other ZIFs such as ZIF-1, ZIF-3, ZIF-zeg, and ZIF-nog, also undergo thermal amorphization and recrystallization to ZIF-zni prior to melting/decomposition^{20, 28}. To complement the X-ray ED data we have collected multi-temperature (25-400 K) single crystal X-ray diffraction data for structural analysis as well as multi-temperature (300-773 K) Raman spectroscopy data to investigate the structural changes in the ZIF-zni frameworks during heating. Finally, the thermal behavior of a series of nine solid solutions of $\text{Co}_x\text{Zn}_{1-x}$ -ZIF-zni (x ranging from 0.3 to 0.003) was analyzed by simultaneous thermogravimetric analysis and differential scanning calorimetry (TGA and DSC) to further probe the role of metal–ligand bonding in regulating the melting/decomposition process in ZIFs.

The dense Zn and Co ZIF-zni MOFs (void volume is only ~4% of the unit cells) are isostructural and they belong to the tetragonal $I4_1cd$ space group. In the asymmetric unit of both crystal structures, two symmetry independent metal centres form a distorted tetrahedral geometry with four imidazolate rings in a 1:2 stoichiometry (Fig. 1 and Fig. S7). The N atoms of the imidazolate ligands are coordinated to the respective metal centres. The metal-nitrogen (M–N) bond lengths (mean value ~ 1.985 Å, Table S2) and $\angle\text{N–M–N}$ bonds angles (mean value ~109.6 °, Table S3) in both 25 K crystal structures are very similar, but differences in the electron configurations of the metal ions may affect the nature and strength of chemical bonding between the metals and the ligands.

Details of the multipole modelling of the X-ray structure factors to obtain the experimental EDs are provided in the method section and supporting information. The d-orbital populations of the Co atoms

derived from experimental ED fits well with their distorted tetrahedral geometry showing a preferential occupancy in d_{xy} orbitals compared to d_{yz} and d_{xz} (Table S4). The experimental static deformation densities reveal details about the metal–N bonds (Fig. 2). The deformation maps show aspherical ED features on Zn metal centres, i.e. a pronounced deviation from the expected spherical $3d^{10}$ density of a Zn atom. The lone pair on the N atoms along the N–Zn–N plane (blue contours in 2D map and blue lumps in 3D maps) are polarized toward the charge concentration regions of the Zn atoms. In contrast to the Zn–N bonding, there is a strong overlap between the N lone pairs and Co orbitals leading to ED accumulation in Co–N bonds. This suggests significant covalency in these bonds similar to the features observed in deformation densities of the C–N and C–C bonds in the imidazolate rings (Fig. S9).

To quantify the metal–ligand bonding we use topological analysis based on QTAIM theory,²⁹ and we compare EDs obtained from experiments as well as single point periodic DFT calculations in the experimental geometry (Table 1 and Table S5). Although the bond path lengths (R) are indistinguishable between the two structures, the Zn–N bonds have lower experimental ρ (ρ) values ($0.51\text{--}0.57\text{ e}\text{\AA}^{-3}$) at the bond critical points (BCPs) compared with the Co–N bonds ($0.60\text{--}0.72\text{ e}\text{\AA}^{-3}$). This corroborates the charge accumulation between Co and N observed in the static deformation maps (Fig. 2). The positive Laplacian values ($\nabla^2\rho$) at the BCPs are typical of metal–ligand bonding³⁰ but we do not observe a clear trend in the experimental values for the Zn–N and Co–N bonds presumably due to experimental noise (see supporting information, Fig. S2 and Fig. S4). The theoretical Laplacian values are consistently less positive in Co–N bonds than Zn–N bonds corroborating a more covalent character in the former. The local kinetic, potential and total energy densities evaluated at the BCPs characterize the metal–ligand interactions as a mixture of closed/shared shell interactions generally observed in transition metal coordination complexes^{30, 31}. The kinetic energy density per electron (G_b/ρ) is slightly higher in the Zn–N

bonds, which corroborates their more ionic nature. Similarly, the larger $|V_b|/G_b$ ratio and larger negative total energy density per electron (H_b/ρ) at the BCPs show that the Co–N bonds have more covalent character than the Zn–N bonds. It is interesting to note that differences in the topological parameters (ρ , G_b/ρ , $|V_b|/G_b$, H_b/ρ) between Zn-ZIF-zni and Co-ZIF-zni are more prominent in the experimental model as compared to the theoretical model, thereby underpinning the importance of experimental benchmark for testing theoretical studies of chemical bonding in extended structures³². In summary, the ED analysis establishes that the Co–N bonds are more covalent than the Zn–N bonds.

Further, we calculated the total density of states (DOS) and projected density of states (PDOS) of Zn-ZIF-zni and Co-ZIF-zni as shown in Fig. S11 and Fig. S12. In Zn-ZIF-zni, analysis of the individual atomic orbitals PDOS reveals that the valence bands (VB) near the Fermi level are largely consist of ligand atoms with negligible participation from the Zn atoms (Fig. S11). On the contrary, the Co 3d orbitals below the Fermi level forms narrow and broad bands in Co-ZIF-zni (Fig. S12). The broad bonding bands due to clear mixing of 3d states of Co atoms and 2p states of C, N atoms in the VB region account for the partial covalent character of Co–N bonds and also hint at the possibility of π -backbonding interactions between Co and imidazolate ligands.

The difference in chemical bonding impact the thermal vibrations of the metal and N atoms in the respective ZIFs. The evolution of atomic displacement parameters (ADP) with increasing temperature provide insights into the structural dynamics³³ and therefore information relevant to the thermal stabilities of Zn/Co-ZIF-zni. We have collected synchrotron single crystal X-ray diffraction data for structural analysis for Zn-ZIF-zni and Co-ZIF-zni at temperatures of 25K, 100K, 200K, 300 K and 400 K. The mean equivalent isotropic ADPs of the metals ($\langle U_{eq}^M \rangle$), N atoms ($\langle U_{eq}^N \rangle$) and C atoms ($\langle U_{eq}^C \rangle$)

are compared for all temperatures (Fig. S13). The Zn atoms have larger ADPs than the Co atoms reflecting the stiffer covalent Co–N bonds, which restrict the movement of the Co atoms. Not surprisingly, the ADPs of the N and C atoms in the imidazolate rings are higher than for the metals (lighter mass). The equivalent isotropic ADPs (U_{eq}) of ring C atoms (Fig. S14) and the rate of change of the ADPs with temperature (Fig. S15) indicate that the C1 and C2 atoms of ring 1 and C4 and C5 atoms of ring 2 display very pronounced out-of-plane libration similar to benzene ring in MOF-5.³³ This reflects significant local distortion of the imidazolate ring with increasing temperature, where the upper part of the ring (M1–N–C–N–M2) vibrates less out-of-plane than the lower part (–C–C–).

To probe the distortion of the frameworks at a molecular level, the ratio of U_{eq} values of ring C atom and its nearest neighbour metal atom (Zn, Co) is calculated as $R_d = U_{eq}(\text{Zn, Co}) / U_{eq}(\text{C})$. The nitrogen atoms are excluded from this parameter because they are bonded to both C and M atoms and hence can introduce biases to this analysis. R_d is an estimate of difference in thermal vibration between the metal atoms and the imidazolate rings, and it is an indirect measure of the framework distortion, see Fig. 3. Beyond 200 K, the ADPs of the “lower ring” C1 and C2 increases strongly compared with Zn and Co atoms, whereas the ADP of the C3 atom bonded to two nitrogens largely follow the trend of the metal ($dU_{eq}(\text{C3})/dT \sim 1.9 \times 10^{-4} \text{ \AA}^2\text{K}^{-1}$, $dU_{eq}(\text{Zn, Co})/dT \sim 1.4 \times 10^{-4} \text{ \AA}^2\text{K}^{-1}$, Fig. S16, Fig. S17). Comparison of R_d values across all temperatures reveal larger difference in ADPs between the tightly bound Co and the flexible C atoms of ring 1 than for Zn and a similar trend is seen for the Co atom and the C atoms of the other imidazolate rings (Fig. S18, S19 and Fig. S20). Co-ZIF-zni clearly exhibits greater framework distortion during heating to 400 K and therefore have lower thermal stability than Zn-ZIF-zni⁹.

It was not possible to study the structural changes beyond 400 K using single crystal X-ray crystallography due to deterioration in data quality. Instead, multi-temperature Raman spectroscopy was employed to indirectly follow structural changes. Sun *et al.* used temperature dependent Fourier transform infrared (FTIR) spectroscopy to study the thermolysis mechanism of ZIF-8 (Zn) and ZIF-67 (Co) MOFs³⁴. Here Raman spectra were recorded from 300 K to 773 K/823 K under nitrogen flow in the range of 1100-1600 cm^{-1} since signature peaks of Zn–N/Co–N bonds were too weak to detect in the lower wavenumber region. The major peaks are assigned³⁵ to the following modes in imidazolate ring- symmetric stretch of longer C–N bonds (1145-1177 cm^{-1}), symmetric stretch of shorter C–N bonds in N–C–N linkage (1255-1284 cm^{-1}) and C=C bond stretching (1490-1495 cm^{-1}) (Fig. S21). The Raman peaks broaden with increasing temperature and the peak positions shift to lower wavenumbers because of lattice expansion. Beyond 700 K, the peak intensities in Co-ZIF-zni drastically reduce and identification of the weaker C=C Raman band becomes difficult.

The symmetric stretching modes [$\nu_s(\text{C1N1}, \text{C2N2})$, $\nu_s(\text{C3N1}, \text{C3N2})$] in the imidazolate rings of Co-ZIF-zni have lower frequencies than in Zn-ZIF-zni at all temperatures indicating weaker bonds in the former MOF (Fig. 3, ring 1 is a representative example for all imidazolate linkers). This observation emphasizes the implicit role of π -backbonding interaction between Co and ligands for weakening the imidazolate bonds of Co-ZIF-zni. The two characteristic C–N stretching modes in Zn-ZIF-zni exhibit a smooth red shift from 1176 cm^{-1} to 1163 cm^{-1} and from 1284 cm^{-1} to 1273 cm^{-1} during heating. In contrast, the shift of peak positions (1167 cm^{-1} to 1144 cm^{-1} , 1277 cm^{-1} to 1255 cm^{-1}) is discontinuous in Co-ZIF-zni, particularly the $\nu_s(\text{C3N1}, \text{C3N2})$ modes show inflection beyond 673 K. This suggests strong ring deformation prior to imidazolate bond dissociation in Co-ZIF-zni as the temperature approaches decomposition point (803 K, Fig. S1b). In Zn-ZIF-zni, the C–N Raman peaks are observed up to 823 K

(Fig. S19), i.e. close to melting point (~ 860 K), suggesting that the imidazolate rings remain intact in the melted form due to relatively stronger imidazolate bonds.

In order to further probe the role of the metal–N bonds on the thermal behavior of these ZIF compounds, we prepared a series of nine solid solution (SS) materials, $\text{Co}_x\text{Zn}_{1-x}\text{-ZIF-zni}$, where the Co mole fraction (x) ranges from 0.3 to as low as 0.003. The final stoichiometry for each SS was determined from Inductively Coupled Plasma Optical Emission Spectrometry (ICP-OES) analysis (Table S11). TGA coupled with DSC measurements were done by heating the SS samples to 750°C in argon to analyze their thermal behavior. The phase diagram for all SS materials is presented in Fig. 4 based on DSC-TGA measurements. T_m is the offset temperature of the endothermic peak in DSC profile, while T_d is the onset decomposition temperature deduced from TGA graphs (Fig. S22).

We observe a systematic dependence of thermal stability of $\text{Co}_x\text{Zn}_{1-x}\text{-ZIF-zni}$ SS phases on percentage of Co doping in the crystal structure (Fig. 4). In the mole fraction range of 0.3-0.04 Co, T_d is lower than T_m ($\sim 7\text{-}21^\circ\text{C}$) suggesting decomposition for these SS materials. Thus, even a very low quantity of Co dopant ($\sim 4\%$) fundamentally alters the thermal stability with respect to the parent Zn-ZIF-zni structure. At ultralow mole fractions of Co ($x = 0.01$ to 0.003) in $\text{Co}_x\text{Zn}_{1-x}\text{-ZIF-zni}$, sharp endothermic peaks of melting are observed as the corresponding T_m values are lower than T_d values by $\sim 11^\circ\text{C}$. The melting behavior is observed macroscopically where the melted SS compounds lose their needle-like morphology compared with decomposed SS samples (Fig. 4). It is intriguing that a few percent of Co atoms in Zn-ZIF-zni is able to disrupt the lattice and cause decomposition. Hence, in a cluster of $3 \times 3 \times 3$ Zn atoms in the Zn-ZIF-zni lattice mimicking 4% doped Co SS phase, substitution of one Zn by Co causes chemical disruption and decomposition (Fig. S23). This implies that the “correlation length” of the local chemical bonding consisting of at least a trimer of metal atoms extends for about $8\text{-}11 \text{ \AA}$ in Zn-ZIF-

zni. To rephrase it, a theoretical cluster calculation would not represent the Zn-ZIF-zni chemical bonding if not containing at least $4 \times 4 \times 4$ Zn atoms.

In summary, using synchrotron X-ray diffraction we have determined the EDs of a meltable MOF, Zn-ZIF-zni and the isostructural Co-ZIF-zni, which decomposes upon heating. Static ED deformation maps and topological analysis reveal that for Co-ZIF-zni, the Co–N bonds have distinct covalent features, whereas the Zn–N bonds in Zn-ZIF-zni are more ionic. This is furthermore corroborated by band structure calculations, which shows covalent mixing between the valence states of C, N atoms and the 3d orbitals of Co atoms in Co-ZIF-zni. Temperature dependent Raman spectroscopy shows that the imidazolate bonds in Zn-ZIF-zni are relatively more stable than in Co-ZIF-zni, which show an abrupt bond weakening above 673 K. The inherently weaker imidazolate bonds in Co-ZIF-zni could be related to subtle π -backbonding between Co and imidazolate ligands. Multi-temperature X-ray diffraction data corroborate that the thermally induced framework distortion is smaller in Zn-ZIF-zni compared with Co-ZIF-zni. The thermal decomposition versus melting behavior was studied in a series of Co-substituted Zn-ZIF-zni solid solution phases. Even a very low Co concentration of $\sim 4\%$ leads to disruption of the Zn-ZIF crystal lattice causing a fundamental change in its thermal behavior. It appears that the relatively more ionic Zn–N bond and more stable imidazolate linkers facilitate melting in Zn-ZIF-zni. In contrast, the more directional covalent Co–N bonding, and resultant weakened imidazolate bonding, leads to decomposition in Co-ZIF-zni. This shows that the nature of the metal–ligand bonds as well as the strengths of internal bonds in the linker molecule, play a key role in controlling the delicate balance between melting and decomposition process in this class of ZIF compounds. There appears to be considerable chemical space for example through ligand functionalization or choice of metal to design new MOF glasses.

Methods

Sample synthesis. Single crystals of Zn-ZIF-zni, Co-ZIF-zni and solid solution $\text{Co}_x\text{Zn}_{1-x}\text{-ZIF-zni}$ ($x=0.4\text{-}0.003$) were synthesized by autoclave synthesis.

0.5 mmol (148.75 mg) of $\text{Zn}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$ and 4 mmol (272.31 mg) of Imidazole was dissolved in 9ml DMF and 6ml butanol. The resulting solution was transferred to a glass vial kept inside a 25 ml Teflon coated steel autoclave. The autoclave was closed and heated at 150°C for 5 days. White needle shaped crystals were obtained and washed with ethanol 3 times.

0.5 mmol (124.52 mg) of $\text{Co}(\text{OAc})_2 \cdot 4\text{H}_2\text{O}$ and 1 mmol (68.08 mg) of Imidazole was added to 5 ml of quinoline and 5 ml of isoamyl alcohol and stirred for 12 hours at room temperature. The resulting mixture was transferred to a glass vial kept inside a 25 ml Teflon coated steel autoclave. The autoclave was closed and heated at 140°C for 3 days. Purple blocked shaped crystals were obtained and washed with ethanol 3 times.

Nine solid solutions $\text{Co}_x\text{Zn}_{1-x}\text{-ZIF-zni}$ single crystals were synthesized following the procedure for the preparation of pristine Zn-ZIF-zni. The precursor metal salts- $\text{Zn}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$ and $\text{Co}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$ were taken in the following stoichiometric ratios-3:1, 4:1, 6:1, 8:1, 12:1, 20:1, 50:1, 100:1, 200:1.

Single-crystal synchrotron X-ray diffraction. Single crystals synchrotron X-ray diffraction data were measured on Zn-ZIF-zni and Co-ZIF-zni at beamline BL02B1 at SPring-8, Japan. A PILATUS3 X CdTe 1M photon counting detector was used to collect the diffracted intensities at $\lambda = 0.24823(5) \text{ \AA}$ and at $T=25 \text{ K}$ using Helium cryostat. The crystal to detector distance was kept at 130 mm. For both crystals, 1800 frames were collected at four different ϕ angles: 0° , -15° , -30° and -45° with ω scans (0° to 180°) and fixed $2\theta=0^\circ$. Scan width for ω scan was 0.1° per frame with an exposure time of 0.8 seconds per 0.1°

scan. The frames were integrated using SAINT in which the integration box sizes were optimized manually by careful inspection of each run. The data were scaled in the program SADABS,³⁶ which was also used to carry out an absorption and oblique correction. Finally, data averaging was performed with the program SORTAV.³⁷ A completeness of more than 99.5 % was obtained up to 1.35 Å⁻¹ and 1.25 Å⁻¹ for Zn-ZIF-zni and Co-ZIF-zni, respectively, and data redundancies of ~22-25. Two low order reflections (200) and (400) were omitted from the refinement. The (200) reflection was hidden behind the beamstop while the (400) reflection was overexposed. The crystal structures were solved with the SHELXT program³⁸ and refined by SHELXL³⁹ included in WINGX package. The hydrogen atom was located in the difference Fourier map and its position and isotropic ADPs were refined.

Multi-Temperature crystallographic data. Single crystal synchrotron X-ray diffraction data were collected at 100 K, 200 K, 300 K and 400 K for Zn-ZIF-zni and Co-ZIF-zni using the same crystals as for the 25 K high-resolution ED data sets. The crystal to detector distance was kept at 130 mm. For all data sets, 1800 frames were collected at two different ϕ angles: 0° and -45° with ω scans (0° to 180°) and fixed $2\theta = 0^\circ$. The scan width for ω scan was 0.1° per frame with an exposure time of 0.8 seconds. The frames were integrated using SAINT and the data were scaled using SADABS which was also used to carry out an absorption and oblique correction. Data averaging was performed with the program SORTAV. The hydrogen atoms were geometrically fixed using riding model constraints. The higher thermal motions (manifested as large ADPs) of the imidazolate ring 1 and ring 2 atoms in 400 K crystal structures of both Zn-ZIF-zni and Co-ZIF-zni were modeled in two different ways: i) an ordered model, in which the ring atoms are represented as elongated ellipsoids (Fig. S8), and ii) a disordered model, in which the ring atoms are split into two components occupying different crystallographic sites.

Multipole modelling and electron density analysis. The ED was obtained from the experimental structure factors using the Hansen-Coppens multipole formalism²⁷ in the program package XD2015.⁴⁰ The refinements were carried out with scattering factors derived from Slater-type orbital (STO) atomic relativistic wavefunctions found in the SCM data bank of XD2015. In the refinement strategy, experimental weights and a $I/\sigma(I) > 3$ were used against F^2 . Further details about the multipole refinement is provided in the supporting information. The quantitative analysis of the electron density topology was performed using the XDPROP program in XD2015 with the kinetic energy density G estimated using the Abramov functional.⁴¹ The potential energy density V and the total energy density H were calculated using the local virial theorem.

Theoretical calculations. Positional parameters obtained from the multipolar refinement were used for the density functional theory calculations using the Perdew, Burke, and Ernzerhof (PBE) GGA functional⁴² with TZVP⁴³ basis set for Zn, Co, C, N, H atoms included in CRYSTAL17 package.⁴⁴ To avoid issues related to quasi-linear dependence, the coefficients of the outermost s and p functions were recalculated and the f functions were removed. The shrinking factors (IS1, IS2, and IS3) and the reciprocal lattice vectors were set to 4 (with 13 k-points in the irreducible Brillouin zone). The bi-electronic Coulomb and exchange series values for the truncation parameter were set as ITOL1–ITOL4 = 7 and ITOL5 = 14, respectively, for the calculations. Both antiferromagnetic (AFM) and ferromagnetic (FM) spin ordering was considered for SCF calculation of Co-ZIF-zni. Our results revealed that the AFM state is energetically more favorable than the FM state by 230 meV for Co-ZIF-zni and therefore all the calculations were performed on AFM phase of Co-ZIF-zni. In case of Zn-ZIF-zni, the level shifter was set to 0.6 Hartree/cycle with 30 % of Fock/KS matrices mixing. In Co-ZIF-zni, the level shifter was set to 0.7 Hartree/cycle with 50 % of Fock/KS matrices mixing. An SCF convergence limit of the order of 10^{-7}

Hartree was used. Upon convergence, the periodic electron densities as obtained from the wave function was subjected to the topological analysis by Bader's Quantum Theory of Atoms in Molecules (QTAIM)²⁹ using the TOPOND program, which is directly interfaced with the CRYSTAL17 code. For computations of total density of states and individual atomic orbital projected density of states of the metal atoms and C, N atoms, the shrinking factors (IS1, IS2, and IS3) and the reciprocal lattice vectors were set to 14 (with 244 k-points in the irreducible Brillouin zone).

Raman spectroscopy. Multi-temperature Raman spectra were recorded on a Renishaw inVia Reflex Micro Raman spectrometer. In case of Zn-ZIF-zni, the single crystal was exposed to 785 nm 50 mW He laser operated at 10% power and 1200 lines/mm grating. The exposure time was 1 sec and the number of accumulations was 30/40. For Co-ZIF-zni, the single crystal was exposed to 457 nm 9.5 mW Ar-ion laser operated at 50% power and 2400 lines/mm grating. The exposure times were 1-2 sec and the number of accumulations was 30. The temperature dependent Raman studies were done using a Linkam THMS 600 heating-cooling stage. Temperature was monitored through the temperature controller attached to the Linkam stage. All the respective spectra were collected under nitrogen flow. The Raman spectra obtained were smoothened using the Savitsky-Golay algorithm with a polynomial order of two and baseline correction was done to remove the background.

Differential Scanning Calorimetry and Thermogravimetric analysis. For these experiments, a NETZSCH DSC 404 F1 Pegasus instrument was used. All the experiments with a heating rate of 10 K/min were conducted in an argon atmosphere. Instrumental contributions were subtracted by measuring a correction file on an empty, clean sample crucible using an identical heating regime measured immediately before each sample.

Inductively Coupled Plasma Optical Emission Spectrometry (ICP-OES). Samples of $\text{Co}_x\text{Zn}_{1-x}$ -ZIF-zni solid solution were dissolved with PlasmaPure Nitric acid (HNO_3) and diluted to 1% acid concentration with MilliQ-water. The measurements were carried out on a Spectro ARCOS ICP-OES equipped with a Burgener Nebulizer and Cyclonic Spray Chamber with an ASX-520 Auto sampler. For quantification of the elements, standard curves series were measured⁴⁵. The series consisted of concentrations for each element: 0, 0.01, 0.05, 0.1, 0.5, 1, 5, 10, 20, 50 ppm.

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Table 1. Topological analysis at bond critical points (BCPs) of metal–N bonds from experimental multipole model (first line) and periodic DFT calculations (second line in italics). R is the bond path length, which is a gradient trajectory length originating at the BCP and terminating at the nuclei (Fig. S10). ρ and $\nabla^2\rho$ are the electron density and its Laplacian, respectively. G_b , V_b and H_b are the kinetic, potential and total energy density ($H_b = G_b + V_b$) at the BCP, respectively. H_b/ρ and G_b/ρ are in atomic units (a.u.). ε is the bond ellipticity [$\varepsilon=(\lambda_1/\lambda_2)-1$] defined as the ratio of the eigenvalues λ_1 and λ_2 of the Hessian matrix that corresponds to the negative curvature of electron density along two perpendicular directions to the bond path.

Bonds	R (Å)	ρ (eÅ ⁻³)	$\nabla^2\rho$ (eÅ ⁻⁵)	ε	G_b/ρ	$ V_b /G_b$	H_b/ρ
Zn1–N4	1.991	0.533	9.526	0.08	1.36	1.08	-0.11
	<i>1.991</i>	<i>0.606</i>	<i>8.812</i>	<i>0.06</i>	<i>1.26</i>	<i>1.19</i>	<i>-0.24</i>
Co1–N4	1.994	0.678	9.808	0.11	1.30	1.22	-0.28
	<i>1.993</i>	<i>0.619</i>	<i>8.362</i>	<i>0.07</i>	<i>1.22</i>	<i>1.22</i>	<i>-0.27</i>
Zn1–N6	1.987	0.512	9.808	0.09	1.41	1.05	-0.07
	<i>1.986</i>	<i>0.612</i>	<i>8.931</i>	<i>0.05</i>	<i>1.26</i>	<i>1.19</i>	<i>-0.24</i>
Co1–N6	1.987	0.714	10.493	0.09	1.33	1.22	-0.30
	<i>1.987</i>	<i>0.626</i>	<i>8.595</i>	<i>0.03</i>	<i>1.23</i>	<i>1.22</i>	<i>-0.27</i>
Zn2–N3	1.995	0.528	9.961	0.08	1.41	1.06	-0.08
	<i>1.995</i>	<i>0.602</i>	<i>8.746</i>	<i>0.06</i>	<i>1.25</i>	<i>1.19</i>	<i>-0.23</i>
Co2–N3	1.990	0.719	9.669	0.09	1.27	1.26	-0.33
	<i>1.989</i>	<i>0.624</i>	<i>8.393</i>	<i>0.08</i>	<i>1.22</i>	<i>1.22</i>	<i>-0.27</i>
Zn2–N7	1.979	0.564	10.471	0.09	1.42	1.08	-0.12
	<i>1.979</i>	<i>0.622</i>	<i>9.184</i>	<i>0.05</i>	<i>1.28</i>	<i>1.19</i>	<i>-0.24</i>
Co2–N7	1.987	0.673	10.652	0.09	1.36	1.18	-0.25
	<i>1.981</i>	<i>0.635</i>	<i>8.732</i>	<i>0.05</i>	<i>1.25</i>	<i>1.22</i>	<i>-0.27</i>

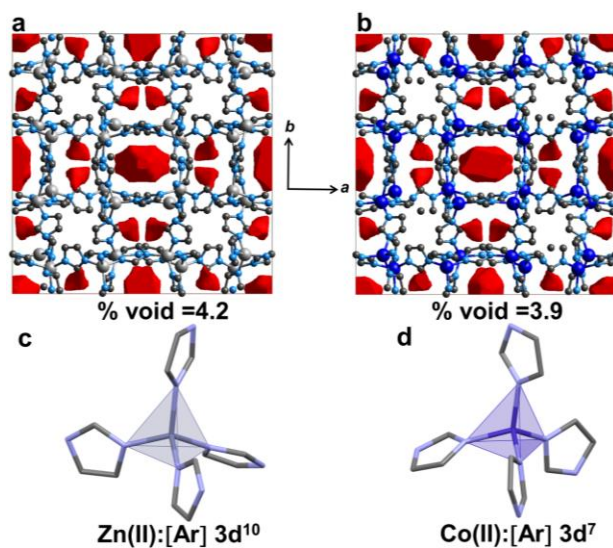


Fig. 1 | Structure of ZIF-zni MOFs. **a-b,** (Top) Unit cell packing with voids⁴⁶ (red surfaces) viewed along the *c* axis of unit cell in (a) Zn-ZIF-zni and (b) Co-ZIF-zni. **c-d,** Building unit of ZIF-zni with (c) Zn (II) and (d) Co (II) in a tetrahedral coordination with four crystallographic independent imidazolate rings.

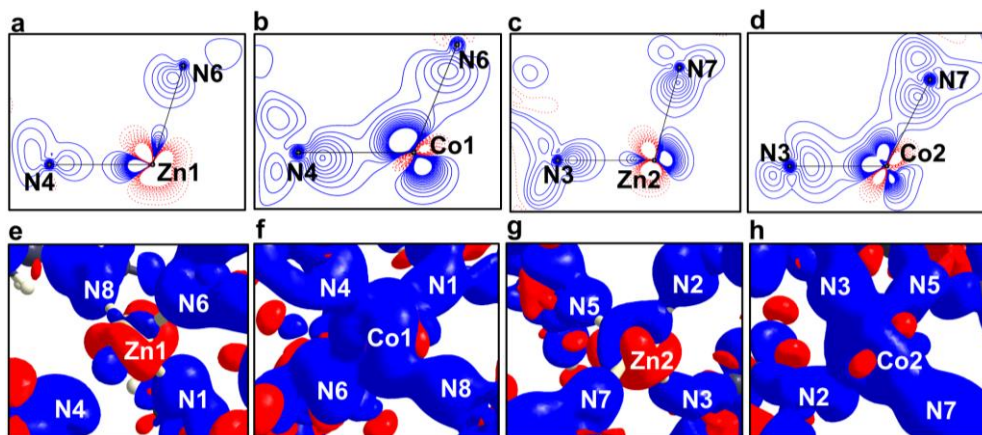


Fig. 2 | Static deformation maps from multipole modelling of experimental structure factors. a-d, 2D static deformation density in the plane of N–M–N for (a, c) Zn-ZIF-zni, (b, d) Co-ZIF-zni. Blue solid lines and red dots represents positive and negative contours, respectively, and the contour level is $\pm 0.1 \text{ e}\text{\AA}^{-3}$. e-h, 3D static deformation density of MN₄ tetrahedra for (e, g) Zn-ZIF-zni, (f, h) Co-ZIF-zni. Blue regions indicates charge concentration region while red regions indicate charge depleted regions and the contour level is $\pm 0.1 \text{ e}\text{\AA}^{-3}$.

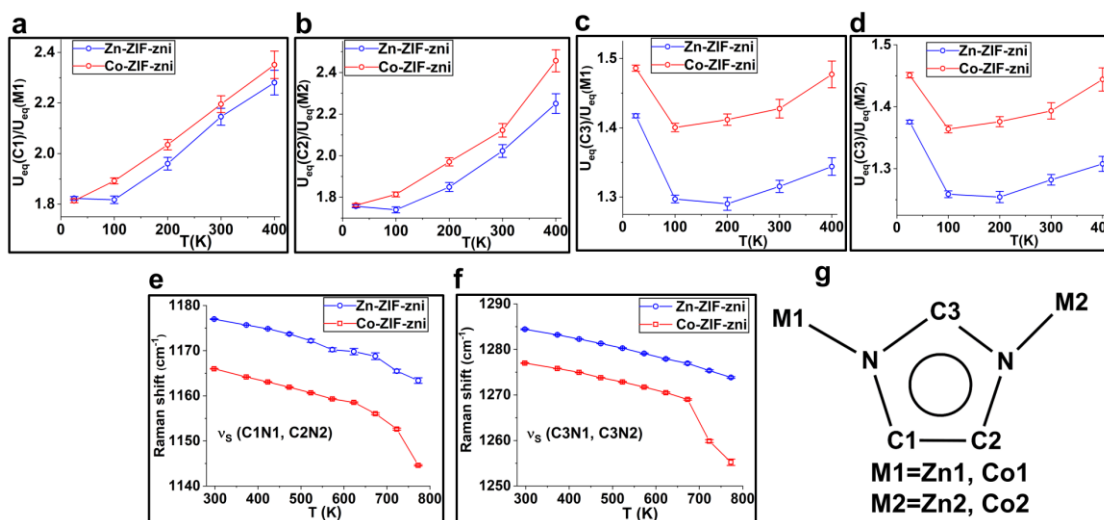


Fig. 3 | Structural variation of Zn-ZIF-zni and Co-ZIF-zni with temperature. **a-d** Comparison of temperature-dependent variation in R_d (ratio of U_{eq} values of C atoms of ring 1 (g) and its nearest neighbour Zn/Co) based on variable temperature synchrotron X-ray diffraction data. **e-f**. Comparison of temperature-dependent variation in Raman peak shifting corresponding to (e) symmetric stretch of longer C–N bonds ($1145\text{--}1177\text{ cm}^{-1}$) and (f) symmetric stretch of shorter C–N bonds in N–C–N linkage ($1255\text{--}1284\text{ cm}^{-1}$) of imidazolate rings. **g**, Molecular diagram of ring 1 bonded with respective metals as a representative example for all imidazolate linkers.

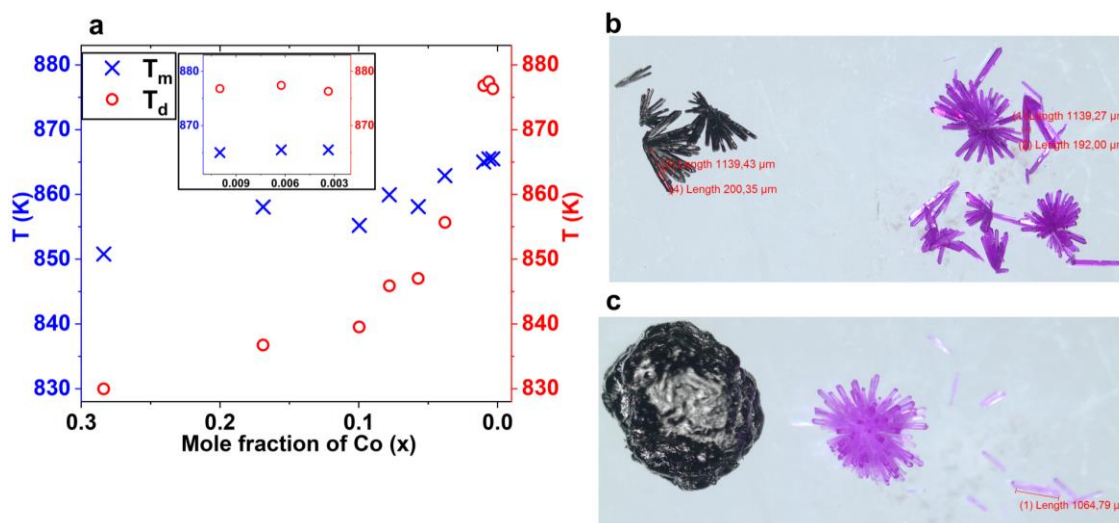


Fig. 4 | Thermal behavior of Co-doped Zn-ZIF-zni crystals. **a**, Phase diagram of $\text{Co}_x\text{Zn}_{1-x}\text{-ZIF-zni}$ ($x=0.3$ to 0.003) solid solutions. T_m is the offset temperature of the calorimetric melting peak in DSC plots while T_d is the extrapolated onset decomposition temperature on the TGA curves (Fig. S22) obtained from simultaneous TGS/DSC experiments. The respective T_m and T_d for the ultralow mole fractions of Co ($x = 0.01$ to 0.003) in $\text{Co}_x\text{Zn}_{1-x}\text{-ZIF-zni}$ is shown in inset. **b-c**, Optical images of two solid solution sample (b) $\text{Co}_{0.04}\text{Zn}_{0.96}\text{-ZIF-zni}$, (c) $\text{Co}_{0.01}\text{Zn}_{0.99}\text{-ZIF-zni}$ before heating (right-purple crystals) and after heating (left-black product) experiments in argon atmosphere upto 750°C . The scale of individual images is 1mm.