Nitrile IR intensities characterize electric fields and hydrogen bonding in protic, aprotic, and protein environments

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Supporting Information Placeholder

ABSTRACT: Nitriles are widely used as vibrational probes; however, the interpretation of their IR frequencies is complicated by hydrogen bonding (H-bonding) in protic environments. We report a new vibrational Stark effect (VSE) that correlates the electric field projected on the -C≡N bond to the transition dipole moment and, by extension, the nitrile peak area or integrated intensity. This linear VSE applies to *both* H-bonding and non-H-bonding interactions. It can therefore be generally applied to determine electric fields in all environments. Additionally, it allows for semi-empirical extraction of the H-bonding contribution to the blueshift of the nitrile frequency. Nitriles were incorporated at H-bonding and non-H-bonding protein sites using amber suppression, and each nitrile variant was structurally characterized at high resolution. We exploited the combined information now available from variations in frequency and integrated intensity and demonstrate that nitriles are a generally useful probe for electric fields.

Noncovalent interactions play pivotal roles in chemistry, from microscopic solvation environments to electrode interfaces, homo- and -heterogenous catalysts, biological membranes, and enzyme active sites. Despite their importance, it can be difficult to experimentally quantify the many types of interactions present in the condensed phase. One quantitative metric for these interactions is the electric field, often determined through the vibrational Stark effect (VSE) on the frequency.[1]

The first vibrational Stark spectra in an *external* applied electric field were analyzed for simple nitriles (-C≡N).[2–4] Nitriles are often found on drugs,[5,6] and several methods have been developed to introduce nitrile probes at a range of positions in proteins,[7–10] nucleic acids,[11,12] biological membrane components,[13] as well as in non-biological settings.[14,15] Nitrile IR transitions are relatively strong and, unlike carbonyls (-C=O), occur in an uncluttered region of the IR spectrum. Unfortunately, the interpretation of nitrile frequency shifts using the language of the VSE is complicated by a well-known blueshift in H-bonding solvents,[10,16–18] undermining the utility of nitrile frequency shifts as direct quantitative probes of local electric fields, though changes in frequency can be useful at a qualitative level.[19]

The vibrational frequency of a well-isolated oscillator such as -C≡N or -C=O can be approximated to first order in the field by the linear (dipolar) VSE:

$$\bar{v}(\vec{F}) = \bar{v}_0 - \Delta\vec{\mu} \cdot \vec{F} \tag{1}$$

where $\Delta\vec{\mu}$ is the vibrational Stark tuning rate or difference dipole moment which determines the sensitivity of a probe frequency to the electric field, $\vec{F}$, projected by the environment onto the probe's difference dipole moment (typically parallel to the bond axis),[2] and $\bar{v}_0$ is the zero-field (vacuum) frequency. This frequency-field dependence can be calibrated by measuring vibrational solvatochromism and calculating the average fields using molecular dynamics (MD) simulations, as shown for a carbonyl group in Fig. 1A. This approach has been used with functionally relevant carbonyls to correlate electric fields at enzyme active sites with activation free energies.[20–22] While a similar linear frequency-field correlation also applies to nitriles for aprotic solvents, the aforementioned blueshift in protic solvents lies far from this correlation, shown for benzonitrile's -C≡N in Fig. 1B. Computational approaches which account for H-bonding frequency shifts are available but difficult to implement experimentally,[16,23] and experiments characterizing H-bonding[24,25] have thus far not been used to correct H-bond frequency shifts in order to obtain electric fields.

While studying nitriles introduced as probes of electric fields in simple solvents and proteins, we observed substantial variation of the nitrile intensity, an effect noted before but not analyzed quantitatively.[15,26–28] Intriguingly, we found that the integrated intensity correlates with solvent field. As shown in the following, this new VSE circumvents the complication of the blueshift in H-bonding environments. The sensitivity of the transition dipole moment (TDM), $\vec{m}$, to an electric field can be expressed as

$$\vec{m}(\vec{F}) = \vec{m}_0 - \vec{F} \cdot \underline{A} - \vec{F} \cdot \underline{B} \cdot \vec{F} + \vartheta(F^3) + \dots \tag{2}$$

where $\underline{A}$ is the transition polarizability, $\underline{B}$ is the transition hyperpolarizability, and $\vec{m}_0$ is the zero-field TDM.[29] Paralleling Equation 1, $|\vec{m}|$ depends on $\underline{A}$ in a manner analogous to $\bar{v}$'s dependence on $\Delta\vec{\mu}$. The $|\vec{m}|$ can be obtained directly from the concentration normalized peak area by

$$|\vec{m}| = \sqrt{\frac{3\epsilon_0 hc \ln 10}{2\pi^2 N_A} \int \frac{\varepsilon(\bar{v})}{\bar{v}} d\bar{v}} \tag{3}$$

where $\varepsilon(\bar{v})$ is the extinction coefficient and $\int \varepsilon(\bar{v})d\bar{v}$ is the integrated peak area, and ϵ_0, h, c, N_A, and $\bar{v}$ are the vacuum permittivity, Plank's constant, speed of light, Avogadro's number, and wavenumber frequency, respectively.[30]
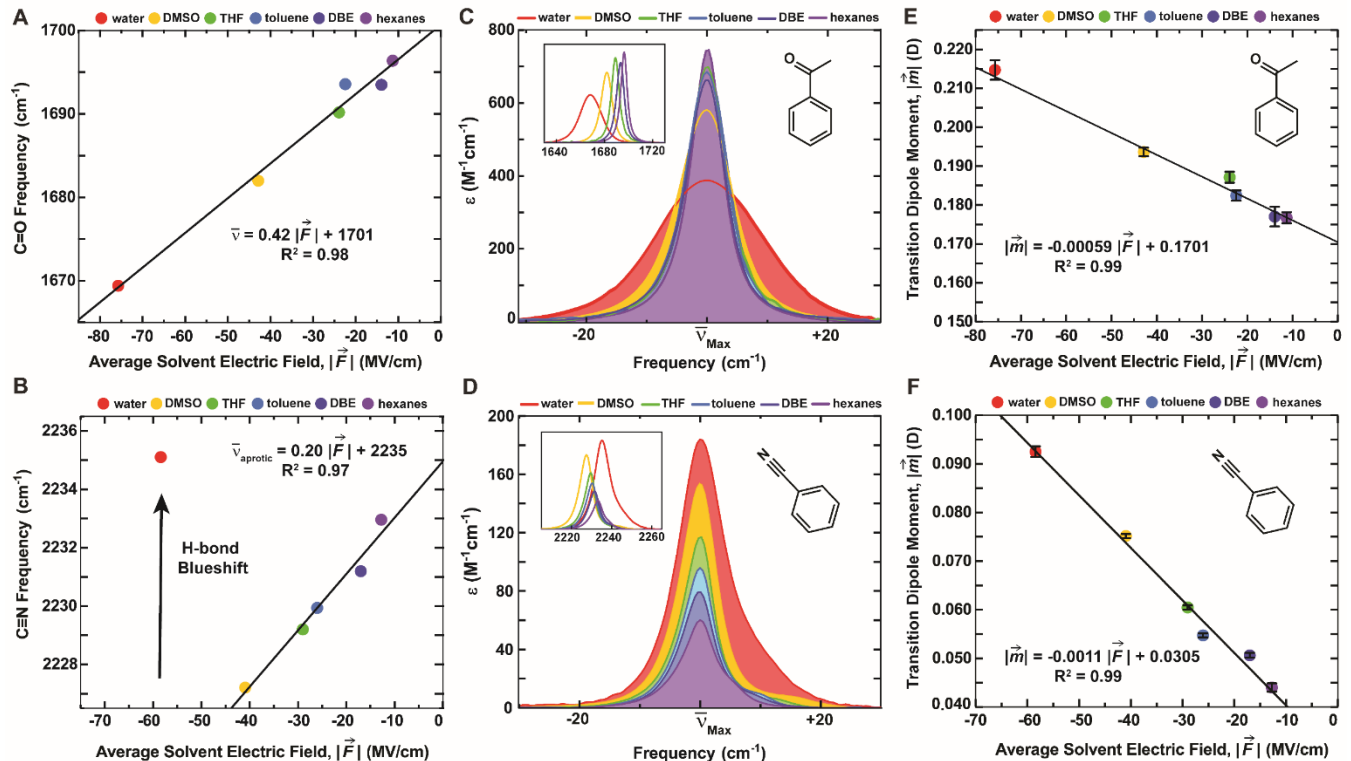
Figure 1. Peak frequency and peak area both change as a function of solvent electric field. (A, B) Plots of peak frequencies *vs* average solvent electric field calculated with polarizable MD for (A) ACP and (B) BZN. In (B) an arrow highlights the blueshift in -C≡N frequency due to H-bonding in water. (C, D) Overlays of IR spectra of (C) ACP and (D) BZN where the IR peak position ($\bar{\nu}_{max}$) has been centered in the plot and peaks are shaded to highlight peak area changes. Peak area increases with solvent polarity. Insets show non-centered spectra. (E, F) Plots of the TDM *vs* the calculated average solvent electric field for both ACP (E) and BZN (F).

FTIR spectra were measured for acetophenone (ACP, -C=O band) and benzonitrile (BZN, -C≡N band) in a variety of aprotic and protic solvents of increasing polarity (Figure 1C-D). The relative increase in peak area from hexanes to water is only 48% for ACP's -C=O but is 342% for BZN's -C≡N (Figure S4). MD simulations using the polarizable AMOEBA09 force field[31] in Tinker 8.7[32] were used to calculate the average solvent electric field, $|\vec{F}|$, projected onto the -C=O and -C≡N probes in the same solvents, as described previously[33–36]. We demonstrate a linear trend when correlating $|\vec{m}|$, obtained from the peak areas using Equation 3, with the average $|\vec{F}|$ exerted on ACP's -C=O (Figure 1E). Importantly, for BZN's -C≡N, a linear relationship is also observed but with a steeper slope (Figure 1F), indicating a larger sensitivity of TDM to field relative to ACP's -C=O (see SI Section S6). In the context of the VSE on $|\vec{m}|$(Eq. 2), the linear correlation between $|\vec{m}|$ and the average $|\vec{F}|$ demonstrates that higher order terms beyond $\underline{A}$ do not meaningfully contribute to tuning $|\vec{m}|$, as expected.[3,37] Moreover, the linear relation between $|\vec{m}|$ and the average $|\vec{F}|$ is strictly monotonic for all solvents and well captured by the VSE, in contrast to the complicated picture for the -C≡N frequency dependence observed in Fig 1B.[24,35] Taken together, by measuring the effect of solvent electric field on peak area, we are offered a direct, orthogonal calibration for measuring electric fields at nitriles which applies regardless of nitrile H-bonding status.

To demonstrate the utility of this new observable, we labeled photoactive yellow protein (PYP) with *o*-cyanophenylalanine (oCNF; Figure 2A-B) introduced site-specifically at multiple sites by amber suppression.[38] PYP is a 14 kDa protein with an endogenous chromophore which can be used to measure protein concentration. Four native phenylalanine sites were targeted for nitrile probe incorporation (Figures 2A and S11) such that the nitrile is placed in electrostatically distinct non-H-bonding and H-bonding environments. High resolution X-ray structures (<1.2 Å; Figure 2C–F) show that the nitrile probes are in a single orientation and occupy distinct environments: F28oCNF possesses one H-bond donor (3.2 Å heavy atom distance), F92oCNF has two donors (2.9 and 3.1 Å), and F62oCNF and F96oCNF are both in nonpolar environments (only carbons within 3.5 Å of the nitrile nitrogen).

The IR spectra of these PYP variants reveal significant variation of *both* the -C≡N peak area and frequency (Figure 3A). A calibration curve was made for *o*-tolunitrile (oTN) as a model compound for oCNF (Figure S7C). The peak areas for the nitrile in F96oCNF and F62oCNF are similar to oTN's peak area in hexanes, while the areas for F28oCNF and F92oCNF are considerably larger and indicate water-like environments, consistent with expectations from X-ray structures. Translating the peak areas to TDMs gives values for $|\vec{m}|$ ranging from 0.050 ± 0.002 D in F96oCNF to 0.096 ± 0.002 D in F92oCNF (Table 1). Using the oTN TDM calibration (Figure S7C), we extract the average electric fields, which range from -13 ± 2 MV/cm to -61 ± 2 MV/cm. Note that the electric field projected on the nitrile of F96oCNF is ~4 times smaller than that on F28oCNF's nitrile even though the peak frequencies are quite similar, illustrating that TDM tuning allows for identification of non-equivalent fields between species possessing nearly equivalent frequencies.
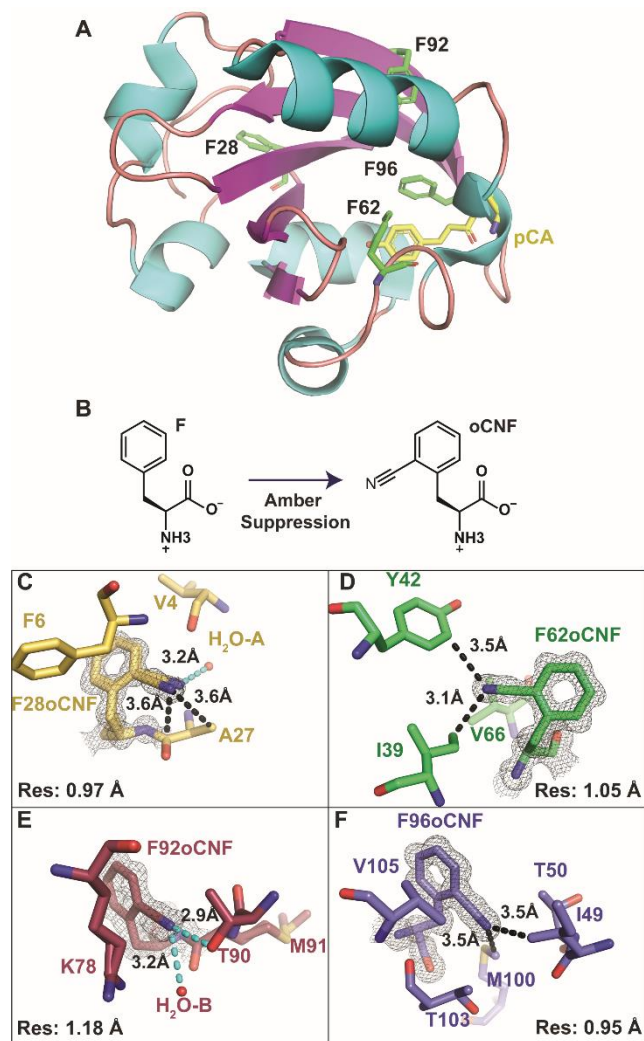
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Figure 2. Overview of nitrile incorporation sites in PYP and their structural characterization. (A) The four phenylalanine sites targeted for nitrile probe incorporation are displayed in green, and the PYP chromophore, *p*-coumaric acid (pCA), is shown in yellow. (B) Native PYP phenylalanines (F) were replaced with genetically encoded *o*-cyanophenylalanine (oCNF) via amber stop codon suppression in four separate variants. Crystallographic structures of the nitrile site are displayed for F28oCNF (C), F62oCNF (D), F92oCNF (E), and F96oCNF (F) PYPs. Electron density maps for the oCNF in each structure are displayed ($2mF_o$-DF_c contoured at 1σ) along with resolution of the crystal structure (bottom corner) and the distance between the nitrile nitrogen and the nearest two atoms (H-bond donor distances in cyan). For F28oCNF, two highly similar oCNF conformers are present; the distance of H_2O-A to the closer nitrile conformer is shown.

Various additive models have been developed to describe the behavior of the -C≡N stretching frequency in the presence of H-bonds (electrostatic multipole expansion;[23] separation into forces of various origin;[16] empirical models,[18] see discussion in SI Section S17). We adopt a simple semi-empirical, experimentally useful model and partition the observed nitrile frequency ($\bar{\nu}_{obs}$) relative to vacuum ($\bar{\nu}_0$) into additive contributions from the shift due to the dipolar VSE ($\Delta\bar{\nu}_{non-HB}$; obtained from solvatochromism in aprotic solvents) and from the frequency shift due to H-bonding not captured by the dipolar VSE ($\Delta\bar{\nu}_{HB}$):

$$\bar{\nu}_{obs} = \Delta\bar{\nu}_{HB} + \Delta\bar{\nu}_{non-HB} + \bar{\nu}_0 \qquad (4)$$

as illustrated in Figure 3C. Using the aprotic oCNF nitrile frequency-field calibration ($|\Delta\vec{\mu}| = 0.19 \frac{cm^{-1}}{MV/cm}, \bar{\nu}_0 = 2231.4$ cm-1; Figure S8B) and fields inferred from the TDM-field calibration (Figure 3B), $\Delta\bar{\nu}_{non-HB}$ was determined (Figure 3C). Per Equation 4, $\Delta\bar{\nu}_{HB}$ is the difference between $\bar{\nu}_{obs}$ and $\Delta\bar{\nu}_{non-HB} + \bar{\nu}_0$ (Figure 3C). Using this approach, we observe a large range in $\Delta\bar{\nu}_{HB}$, with the most H-bonding variants, F28oCNF and F92oCNF, possessing $\Delta\bar{\nu}_{HB}$ of 7 and 21 cm-1, respectively (Table 1). This is an intriguing observation as $\Delta\bar{\nu}_{HB}$ has been shown to be sensitive to H-bond geometry (i.e., H-bond angles and distances)[23,39] and the geometry determines H-bond strength.[40]

To further refine our analysis, we calculated the dependence of $\Delta\bar{\nu}_{HB}$ on the heavy atom H-bond angle (O···N≡C, α in Figure 4, proton positions are not resolved in the X-ray structures) for comparison with experimental results. Geometry optimization and vibrational frequency analysis [b3lyp/6-311++g(d,p)] were performed on oTN and a single water where α was varied from 70°–175° (see SI Section S16). H-bond distance was similar (3.0–3.2 Å over 90–175°) following geometry optimization, and DFT-calculated frequencies demonstrate an angle-dependence to the H-bond shift for oTN (Figure S19D) with the most significant blueshifts near 180°, similar to Choi *et al.*[39] Given the H-bond angles observed between F92oCNF and T90 (174°) and between F28oCNF and H_2O-A (110°), these DFT calculations predict $\Delta\bar{\nu}_{HB}$ of ~4 and ~22 cm-1 for F28oCNF and F92oCNF, respectively. While H-bond distance and number of donors[16] also contribute to $\Delta\bar{\nu}_{HB}$, we observe good agreement between predicted $\Delta\bar{\nu}_{HB}$ and observed values (7 cm-1, 21 cm-1). These results indicate the potential of $\Delta\bar{\nu}_{HB}$ to report on H-bonding geometries, though further benchmarking is necessary.

In summary, combined analysis of frequency and integrated peak intensity tuning with nitrile probes can be used to extract environmental electric fields and characterize H-bonding. TDM analysis enables quantification of significantly different fields for nitriles whose frequencies are nearly identical. We have utilized a simple phenomenological approach in this communication motivating a more comprehensive theoretical description of TDM and frequency tuning. It may be useful to re-evaluate intensity changes seen in other studies.[15,27] Moreover, nitrile intensities can be used to determine electric fields and H-bonding in the extensive array of systems in which nitriles have been incorporated.[7,9–11,13,14,41]
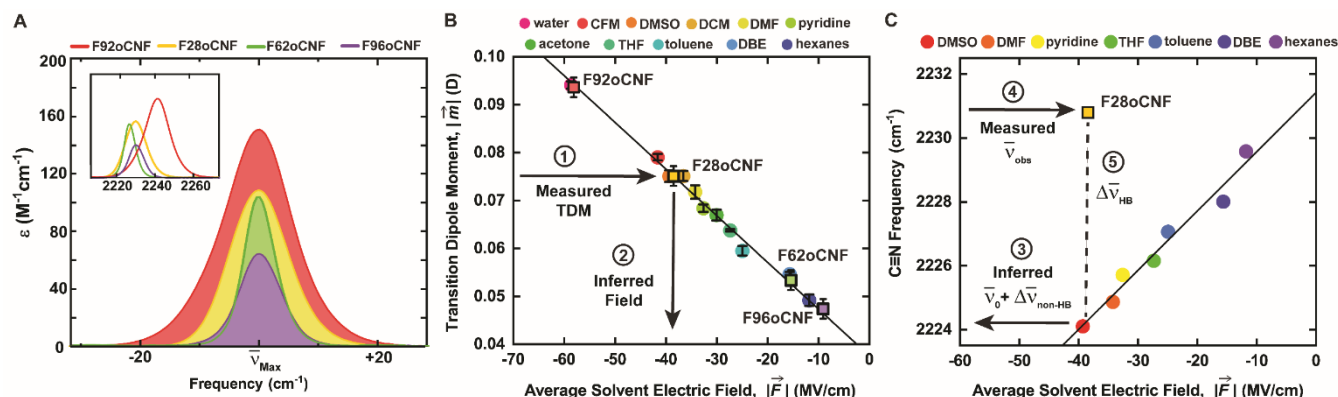
Figure 3. IR spectra of the -C≡N stretch from incorporated oCNF residues at four sites in PYP (c.f. Figure 2) and extracted parameters. (A) Shaded IR overlay of the four labeled sites highlights the significant change in peak area from protein nitrile environments. Inset contains non-centered spectra. (B) Using the measured TDM (1), the electric field for proteins can be inferred from the oTN TDM *vs* field calibration curve (2) (Figure S7C) as demonstrated for F28oCNF. (C) The aprotic solvent field *vs* frequency calibration curve for oTN. Using the field inferred from the TDM (2), one can infer $\bar{v}_0 + \Delta\bar{v}_{non-HB}$ (3), and by measuring $\bar{v}_{obs}$ (4), the blueshift ($\Delta\bar{v}_{HB}$) can be deduced (5)(Table 1).

Table 1. Extracted IR peak parameters for the oCNF nitrile incorporated at different PYP sites

Environ- ment	$\bar{v}_{obs}$ (cm^{-1})[a]	FWHM (cm^{-1})[a]	$|\vec{m}|$ (D)	$|\vec{F}|$ (MV/cm)	$\Delta\bar{v}_{non-HB}$ (cm^{-1})[b]	$\Delta\bar{v}_{HB}$ (cm^{-1})[b]
Hexanes	2229.6	5.1	0.049 ± 0.001	-12 ± 1	-2	0
Water	2231.7	9.5	0.094 ± 0.001	-59 ± 1	-11	11
F28oCNF	2230.8	12.7	0.075 ± 0.002	-38 ± 2	-6	7
F62oCNF	2228.1	6.5	0.053 ± 0.002	-15 ± 2	-3	0
F92oCNF	2241.3	14.4	0.094 ± 0.002	-58 ± 2	-11	21
F96oCNF	2231.2	8.4	0.047 ± 0.002	-9 ± 2	-2	2

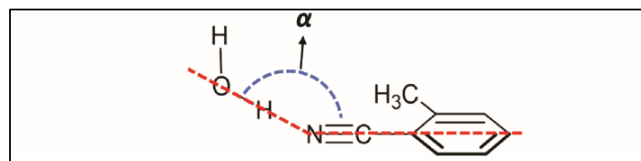[a] Peak positions ($\bar{v}_{obs}$) and FWHM all have an error of <0.1 cm^{-1}.

Figure 4. The heavy atom H-bond angle (O···N≡C; α) between a water and the nitrile of oTN, angle shown in blue.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge on the ACS Publications website.

Experimental details on sample preparation and IR spectroscopy, IR spectra and TDM calibration curves for oTN, fixed-charge TDM calibration curves for ACP and BZN, Stark spectroscopy data and discussion, UV-vis spectra of PYP variants, mass spectra of PYP variants, crystallography procedure and structural details (PDF)

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