First Principles Approach to Extracting Chemical Information from X-Ray Absorption Near-Edge Spectra of Ga-Containing Materials

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Abstract

The X-ray absorption near edge structure (XANES) can provide uniquely detailed information on the coordination environments of important Ga-containing materials with unknown structures, including catalytically-active materials. In this study, the Ga K-edge XANES was simulated using first principles-based methods for seven molecular Ga complexes, as well β -Ga₂O₃, in order to explore the chemical origins of the experimentally observed features. The theoretical spectra were computed using FEFF, CASTEP and StoBE, in order to assess the sensitivity of the results to the computational approach. While the XANES features depend on the Ga coordination environment, they are also sensitive to the electronegativity of the ligands and the symmetry at Ga. The white line position responds to changes in both the core state (due to differential screening) and the valence "p" states (arising from differences in ligand coordination).



Graphical TOC Entry

1. Introduction

Ga-based heterogeneous catalysts are interesting potential replacements for Pt- and Crbased catalysts in propane dehydrogenation,¹ and are the basis for a commercial process for propane dehydroaromatization.² Based on X-ray absorption spectra, IR spectra, and chemical titration evidence, as well as Density Functional Theory (DFT) modeling, the Ga sites proposed to be present in these materials under reaction conditions (typically, ca. 500 °C in the presence of H₂) have been variously described as mononuclear [Ga]⁺, [GaO]⁺, [GaH(OH)]⁺, [Ga(OH)₂]⁺, [GaH]²⁺, and [GaH₂]⁺,³⁻⁶ as well as dinuclear [Ga₂O₂]²⁺ and [Ga₂O₂H₂]²⁺,⁷⁻⁹ All are formally Ga(III), except for [Ga]⁺. Nevertheless, Ga(I) sites feature prominently in some proposed mechanisms for propane dehydrogenation,^{3, 5, 10, 11} where their participation is a topic of on-going debate due to controversies about spectroscopic assignments.

In principle, X-ray absorption spectroscopy at the Ga K-edge can provide rich information about the local coordination and electronic structures of Ga sites, even at elevated reaction temperatures and in the presence of reactants.^{3, 12-15} After calcination, supported Ga-containing catalysts have coordination numbers close to 4 with Ga-O distances ranging from 1.73 to 1.93 Å,^{4, ^{6, 13} according to best-fit models of the EXAFS data. Activation in flowing H₂ at 650 °C causes the amplitude of the Ga-O scattering path to decrease, implying that the active sites have a lower average coordination number, and/or that some of the oxygen-donor ligands have been replaced by hydrogen.^{4, 13}}

Although the Ga K-edge XANES is sensitive to structural and electronic changes, rigorous XANES assignments remain challenging. The energies and intensities of the $1s \rightarrow 4p$ transitions which occur in the region of the absorption edge reflect both the Ga oxidation state and its local bonding geometry. After calcination, supported Ga catalysts show broad white line (WL) maxima in the range 10375 - 10380 eV, similar to Ga(III) standard materials such as β -Ga₂O₃ and Ga(acac)₃.^{13, 16} Two peaks at the absorption edge have been assigned to tetrahedral and octahedral Ga(III) at photon energy values of 10375 and 10379 eV, respectively, and their intensities were used to estimate the fractions of each type of Ga site.¹⁷ During activation in flowing H₂, the WL sharpens and shifts to lower energies (10368 – 10373 eV).^{3, 4, 6, 10} In early studies, this shift was attributed to the reduction of Ga(III) to Ga(I),³ by comparison to the energy of a peak assigned to the Ga(I) component in mixed-valent Ga₂Cl₄ (i.e., Ga(I)Ga(III)Cl₄).¹⁰ However, this assignment was subsequently questioned in a XANES study of well-defined molecular Ga(III) complexes,¹⁶ whose edge positions cover the full range of energies assigned. In particular, the WL of trigonal planar Ga(CH₂SiMe₃)₃ appears at ca. 10371 eV, compared to 10370 eV for the Ga(I) site in Ga₂Cl₄. Thus, changes in the nature of the ligands and/or the Ga coordination number can cause the edge position to shift as much as a reduction in the formal oxidation state, with implications for the identification of active sites and reaction mechanisms in heterogeneous Ga catalysts.

To understand the fundamental causes of these dramatic differences, we undertook a study of variations in the Ga K-edge XANES for a series of six well-defined Ga(III) compounds, namely, $[Ga(OSiPh_3)_3]_2$, $[Ga(CH_2SiMe_3)(OSiPh_3)_2]_2$, $[Ga(CH_2SiMe_3)_2(OSiPh_3)]_2$, $Ga(CH_2SiMe_3)_3$, $Ga(CH_2SiMe_3)_3(OP(n-Bu)_3)$ and $Ga(t-Bu)_2(OC_6H_2-2,4,6,-t-Bu_3)$, as well as for β -Ga₂O₃, a reference material often used in XANES studies of Ga(III). Clearly, the widely differing appearances of the experimental spectra for the molecular complexes (shown in Figure 1 and originally reported by Getsoian et al.)¹⁶ are not a consequence solely of Ga oxidation state, which is formally (III) in all of the complexes, but also of their coordination environments. For this study, we performed simulations of the Ga K-edge XANES for each molecular complex using three firstprinciples methods: FEFF,^{18, 19} CASTEP,²⁰ and StoBe.²¹ The results reveal the origins of the differences in the electronic structures of the Ga(III) compounds, and have implications on the interpretation of XANES characterization measurements for Ga-based catalysts.

2. Methods

CASTEP Simulations. Quantum mechanical optimizations of all molecular structures (as depicted in Figures S1-S7) were performed with the CASTEP code,²² using the PBE functional. Ga K-edge X-ray absorption near-edge structure (XANES) spectra were simulated by generating a Ga pseudopotential with a 1s core hole, and performing a regular SCF cycle to generate the Kohn-Sham states (occupied and unoccupied) under the influence of the core hole.²⁰ The valence electron Kohn-Sham states were expanded using a series of plane waves, with a cutoff energy of at least 570 eV. Core electrons were treated using ultrasoft pseudopotentials, generated using allelectron atomic calculations prior to the beginning of standard self-consistent field cycles.²² Electronic exchange and correlation were computed within the generalized gradient approximation using the Perdew-Burke-Erzenhof (PBE) functional.²³ For calculations involving isolated molecules, the reciprocal space was sampled using only the Γ point. Prior to electronic structure analysis, molecular geometries were optimized until all atomic forces were less than 0.03 eV/Å. For calculations involving β -Ga₂O₃²⁴, the RPBE functional was used with an optimized volume of 1357 Å³ (a = 12.63 Å, b = 9.32 Å, c = 11.85 Å, $\alpha = 90.0^{\circ}$, $\beta = 103.7^{\circ}$, $\gamma = 90.0^{\circ}$). We note that the unit cell used here was a 1×2×3 unit cell to avoid artifacts in the computed XANES spectra due to the imposed periodic boundary conditions. The optimized geometries for the molecular complexes and the bulk oxide structure are provided as supplementary information in CIF format. A k-point grid (4×5×6) was used to calculate the XANES of β -Ga₂O₃. In order to enable

quantitative comparison of StoBE and FEFF results, the XANES spectra of some complexes and β -Ga₂O₃ were calculated using simplified and/or truncated geometries.

Absorption spectra were computed by evaluating the dipole transition matrix elements between the 1*s* core state and unoccupied Kohn-Sham states. The computed XANES were broadened using Gaussian functions with a full-width at half-maximum (FWHM) of 0.6 eV to simulate machine broadening as well as the default Ga lifetime broadening factors implemented in the Materials Studio package²⁵. For comparison with experiment, each simulated spectrum (which was computed on an energy scale referenced to the Fermi energy) was shifted to align its first maximum with the corresponding peak in the experimental spectrum. More details are given in the Supporting Information.

StoBe and FEFF Simulations. The same molecular structures of Ga(III) complexes (obtained from PBE optimizations using CASTEP) were used in both the StoBe 2011²¹ and FEFF 9.6^{18, 19} simulations. For β -Ga₂O₃, FEFF simulations were performed with clusters of radii 5.0 and 9.0 Å for the SCF and FMS, respectively, to obtain converged results. All other FEFF settings were defaults. For comparison to the published experimental data, the FEFF simulations were broadened using default parameters.

Since StoBe is a molecular code, its use to simulate the spectrum of crystalline β -Ga₂O₃ required the use of clusters (Figures S8 and S9). These were generated from the same experimental structure²⁶ in order to represent the tetrahedral and octahedral Ga sites consistently. For the StoBe simulations, the core-excited Ga atom was described using the all-electron IGLO F-III basis set. For other atoms, the basis sets and associated effective core potentials for Ga(+13) (51/41/41), Si(+4) (311/211/1), P(+5) (311/211/1), O(+6) (321/311/1), C(+4) (321/311/1), and H (311/1) (as distributed in StoBe) were used. Finally, StoBe simulations were performed using PBE, RPBE or

B88PD86 exchange-correlation functionals.^{27, 28} The StoBe simulations were broadened using an *ad hoc* self-energy with constant real and variable imaginary parts, as described previously.²⁹

X-ray Absorption Spectroscopy. The XAS spectrum of β -Ga₂O₃ was recorded at the Ga Kedge (10,367 eV) on beamline 2-2, Si(111) $\phi = 90$ °, at the Stanford Synchrotron Radiation Lightsource, which operates at 3.0 GeV with a current of 500 mA. β -Ga₂O₃ (99.99 %, Sigma-Aldrich) was diluted to ca. 2 wt% Ga with BN (99.5 %, Alfa Aesar). GaMe₃ (Strem, 99%) was diluted to ca. 2 wt% Ga with n-pentane (Acros Organics, 99%). Data were acquired in transmission mode, using N₂-filled ionization chambers as detectors. For energy calibration, a W foil (L_{III}-edge, 10,207 eV) was placed after the second ionization chamber following the incident beam direction. Experimental XAS data were processed and analyzed using the Demeter software package.³⁰ A linear function was subtracted from the pre-edge region, then the edge jump was normalized using the Athena program. The $\chi(k)$ data were isolated by subtracting a smooth, third-order polynomial approximating the absorption background of an isolated atom.



Figure 1: Comparison of published¹⁶ experimental XANES for six molecular Ga complexes, with simulations obtained using three theoretical approaches. The vertical lines indicate the position of the first maximum in the experimental XANES for each complex. Each CASTEP simulation was individually shifted to align with these experimental peaks. Spectra simulated using either FEFF and StoBE were shifted by a common amount, based on the shift required to match the experimental maximum for [Ga(OSiPh₃)₃]₂ (black vertical line).

3. Results and Discussion

The experimental and theoretical XANES spectra for six molecular Ga(III) compounds are compared in Figure 1. Since CASTEP does not provide an absolute energy scale, each CASTEP spectrum was adjusted so that its first maximum is aligned with the experimental value. The absolute positions of the CASTEP spectra can be obtained from first principles using the method of Pickard and coworkers³¹ (see Figure S10). In contrast, the FEFF and StoBe spectra require only an overall energy shift, therefore the maximum in the simulated spectrum of [Ga(OSiPh₃)₃]₂ was aligned with the experimental value, and all other simulations performed using the same method were shifted by the same amount. The overall agreement between theory and experiment is striking for both the StoBe and CASTEP simulations. In particular, StoBe provides very good estimates for the relative positions of the first maxima in each complex and reproduces the overall shape of the spectra. The latter is also true for CASTEP. The agreement is less good for the FEFF results. The large differences between the FEFF simulations and the experimental spectra arise from the muffin-tin approximation.³² It is particularly ill-suited for some of these molecular Ga complexes, due to the large local anisotropies in their electronic structures. Nevertheless, FEFF does provide much better agreement when the local environment is more isotropic, as discussed below for β -Ga₂O₃.

The most noticeable differences between theory and experiment occur 8-10 eV above the edge. In this region, the CASTEP simulations resemble the experimental results better, likely due to CASTEP's use of plane waves, *vs.* the local Gaussian-type orbitals (GTOs) used by StoBe, which capture long-range behavior more accurately than a localized basis set. In principle, the accuracy of the simulations may also be affected by the choice of exchange-correlation functional,

and the quality of the structural models. However, Figure 2 shows that the XANES simulations are largely insensitive to the choice of functional (see also Figures S11-S13). In addition, they are relatively insensitive to geometry simplifications distant from the Ga absorber: there was little change when the phenyl substituents located ca. 4 Å distant from Ga were replaced by methyls (Figure S14).



Figure 2: Comparison of the StoBe-predicted XANES for three representative molecular Ga complexes using the B88PD86, PBE and RPBE exchange-correlation functionals. The spectra were shifted as in Figure 1, based on the shift required to match the experimental maximum for [Ga(OSiPh₃)₃]₂.

Table 1: Relative white line energies (WLE) predicted by StoBe for various molecular Ga(III) complexes and β -Ga₂O₃ with differing local Ga coordination environments, and the contributions to WLE of the core ionization energy (CIE) and valence state energy (VSE)^a

Molecular Ga complex	Coordination	WLE	CIE	VSE	Valence	d(Ga-L) ^c
	environment				occupancy ^b	
β -Ga ₂ O ₃ (oct)	O ₆ -Ga	3.01	2.50	0.51	1.21	1.93
β -Ga ₂ O ₃ (tet)	O ₄ -Ga	0.58	2.71	-2.13	1.20	1.83
[Ga(OSiPh ₃) ₃] ₂	O4-Ga	-2.49	0.00	-2.49	1.07	1.79
[Ga(CH ₂ SiMe ₃)(OSiPh ₃) ₂] ₂	O ₃ -Ga-C	-3.46	-0.82	-2.64	1.14	1.84
[Ga(CH ₂ SiMe ₃) ₂ (OSiPh ₃)] ₂	O ₂ -Ga-C ₂	-4.25	-1.47	-2.78	1.20	1.98
$Ga(CH_2SiMe_3)_3(OP(n-Bu)_3)$	O-Ga-C ₃	-5.08	-2.75	-2.33	1.30	2.01
Ga(<i>t</i> -Bu) ₂ (OC ₆ H ₂ -2,4,6- <i>t</i> -Bu ₃)	O-Ga-C ₂	-5.53	-1.31	-4.22	1.18	1.85
Ga(CH ₂ SiMe ₃) ₃	Ga-C ₃	-5.85	-1.57	-4.28	1.28	1.99
GaMe ₃	Ga-C ₃	-5.98	-1.05	-4.93	1.40	1.98

^a All energies are in eV, and are reported relative to the CIE of [Ga(OSiPh₃)₃]₂. ^b Total Ga valence (*sp*) occupancy, from Natural Atomic Orbital analysis. ^cMinimum distance (in Å) between Ga and its near-neighbor atoms.

Table 1 shows the WL energy (WLE) predicted by StoBe for each complex, as well as values for the core ionization energy (CIE) and the energy of the first excited valence state (VSE), where WLE = CIE + VSE. All energies are given relative to the highest CIE, found for $[Ga(OSiPh_3)_3]_2$. The 3-coordinate complexes $Ga(CH_2SiMe_3)_3$ and $Ga(t-Bu)_2(OC_6H_2-2,4,6,-t-Bu_3,$ as well as GaMe₃ (see Figure S15), have lower WLEs than any of the 4-coordinate complexes

[Ga(OSiPh₃)₃]₂, [Ga(CH₂SiMe₃)(OSiPh₃)₂]₂, [Ga(CH₂SiMe₃)₂(OSiPh₃)]₂ and Ga(CH₂SiMe₃)₃(OP(*n*-Bu)₃). A similar dependence of the WLE on coordination number was reported for the K-edges of Cu(II), Zn(II), Al(III), and Ti(IV) complexes.³³⁻³⁷ In the series of Ga compounds, the shift in WLE is associated mainly with a change in the position of the first intense transition, with the 3-coordinate compounds showing larger negative average VSE values (-4.5 \pm 0.4 eV) compared to the 4-coordinate compounds (-2.5 \pm 0.2 eV).

To understand the origin of the coordination number dependence of the WLE, molecular orbitals associated with the first intense transition were computed, and are shown in Figure 3 for three representative cases: 3-coordinate Ga(CH₂SiMe₃)₃, 4-coordinate [Ga(OSiPh₃)₃]₂, and intermediate Ga(CH₂SiMe₃)₃(OP(*n*-Bu)₃), which is a 4-coordinate complex but with a Ga-C₃ moiety that is nearly planar. Since the XANES is closely associated with the final state wavefunction,³⁸ the orbitals calculated by either CASTEP or StoBe for the bound excited valence states associated with the WL are quite similar, despite being calculated using plane waves and GTOs, respectively. The main difference is that the CASTEP orbitals appear to be slightly more delocalized. Orbitals associated with higher energy transitions also have greater spatial extension (Figures S16-S18). However, WLEs are unlikely to be greatly affected by delocalization of bound states, since the dipole matrix elements are determined mainly by contributions from the local environment around the Ga absorber.



Figure 3: Comparison of CASTEP optimized geometries for three representative Ga complexes and the molecular orbitals calculated by CASTEP and StoBe for the excited states associated with their white lines. Color scheme for geometries in top panels: Ga (pink), Si (orange), C (grey), H (white), O (red), P (green). Color scheme for orbitals: Ga (pink), Si (light grey), C (grey), H (white), O (red), P (orange). The orbitals are represented at an isosurface value of 0.02 Bohr^{-3/2}.

Given the dipole selection rules for K-edges, it is not surprising that all the calculated orbitals have considerable Ga *p*-character. The orbital at the WLE associated with 3-coordinate $Ga(CH_2SiMe_3)_3$ is almost "pure" *p*, due to the trigonal planar nature of the Ga-C₃ arrangement. In contrast, the corresponding orbital associated with 4-coordinate, non-planar [Ga(OSiPh_3)_3]_2

includes an antibonding component that reflects coupling to ligand-based orbitals. This difference explains much of the coordination-number dependent VSE shift of ca. 2.0 eV in the associated transitions: relative to the energy of the *p* orbital in the 3-coordinate complex, anti-bonding coupling to the ligand orbitals raises the state energy in the 4-coordinate complex, resulting in an increased transition energy that is closer to the ionization energy. The VSE in $Ga(CH_2SiMe_3)_3(OP(n-Bu)_3)$ is intermediate since, similar to 3-coordinate $Ga(CH_2SiMe_3)_3$, its orbital has well-defined *p*-character, yet the presence of some antibonding coupling involving the $OP(n-Bu)_3$ ligand still increases the transition energy as in 4-coordinate [Ga(OSiPh_3)_3]_2.

As Table 1 shows, the minimum distance and its nearest neighbor atoms varies with the electronegativity of the ligand, from 1.79 Å for O₄-Ga to 2.01 Å for O-Ga-C₃, with a corresponding shift of 2.6 eV in the WLEs. This effect is manifested in the CIE, where the difference between the highest and lowest CIEs (for [Ga(OSiPh₃)₃]₂ and Ga(CH₂SiMe₃)₃(OP(*n*-Bu)₃), respectively) is 2.8 eV (Table 1). The effect is therefore larger than the 2.0 eV change in VSE which results from the increase in coordination number from 3 to 4, and more gradual. The change in CIE correlates with the overall valence occupation of Ga, estimated using Natural Atomic Orbital analysis,³⁹ which is in turn related to the average nearest-neighbor distance. Longer bond distances result in more localized Ga valence states and larger valence occupations, which screen the core states more efficiently. However, the CIE also depends, to a lesser degree, on the second coordination sphere of the complex, since both $Ga(CH_2SiMe_3)_3$ and $GaMe_3$ have the same $Ga-C_3$ nearest neighbor environment, while their CIE values differ by -0.52 eV (with the CIE value for Ga(CH₂SiMe₃)₃) being more negative than for GaMe₃). The CIE also likely reflects the difference between a methyl and a neosilyl due to the hyperconjugation involving the Si atom, which alters the electron density at the carbon donor atom.

We also explored the XANES of β -Ga₂O₃, a common reference material used in spectroscopy of Ga-containing heterogeneous catalysts. This oxide contains both tetrahedral and octahedral Ga sites. Figure 4 shows the experimental spectrum as well as FEFF, CASTEP and StoBe simulations of the β -Ga₂O₃ XANES, including the contributions from each Ga site and their average (where the latter is used for quantitative comparison between theory and experiment). The WL appears at an even higher energy than for 4-coordinate [Ga(OSiPh₃)₃]₂, the most closelyrelated of the Ga(III) molecular systems (10380 vs. 10376 eV, respectively). The agreement between the calculated spectra and the experimental data is particularly good for the FEFF and StoBe calculations.

Surprisingly, all three theoretical methods find that the main peak associated with the 6coordinate site in β -Ga₂O₃ appears at a lower energy than that for the 4-coordinate site. This order seems counterintuitive, given that higher coordination numbers in the molecular Ga complexes described above are correlated with higher WLEs. However, both FEFF and StoBe predict a lower energy shoulder for the 4-coordinate site, which is absent in the contribution from the 6-coordinate site. The shoulder is clearly visible in the experimental spectrum at ca. 10377 eV, just before the WL. Thus, the lowest energy feature for the 4-coordinate site does indeed precede the main peak for the 6-coordinate site. To further deconvolute the XANES dependence on the Ga coordination number, we compare the VSE value of [Ga(OSiPh₃)₃]₂, -2.49 eV, with that of the 4-coordinate site of β -Ga₂O₃, -2.13 eV (Table 1). The small difference is likely due to symmetry: Ga in [Ga(OSiPh₃)₃]₂ is distorted from tetrahedrality relative to the nearly perfect tetrahedral environment in the oxide. Compared to [Ga(OSiPh₃)₃]₂, the higher WLE value for the 4-coordinate site in β -Ga₂O₃ (resulting in the main peak) is due to its positive CIE value. We hypothesize that it reflects the withdrawing effect of oxygen in higher coordination shells on the Ga electron density, as well as the difference in symmetry. Finally, we note that the VSE value for the 6-coordinate site in β -Ga₂O₃ correlates well with its coordination environment, since its positive value of 0.51 eV is larger than for any other complex in Table 1, all of which have lower coordination numbers.



Figure 4: Comparison of the experimental Ga K-edge XANES of β -Ga₂O₃ with spectra simulated using FEFF, CASTEP and StoBe. The theoretical results show the individual simulations for the octahedral and tetrahedral Ga sites (dotted blue and dashed red lines, respectively), as well as their average (solid green lines). The site contributions are weighted equally in the average, according to the experimental structure. The CASTEP-computed components are shifted to align with the maxima computed by FEFF.

4. Conclusions

In summary, first principles simulations of the XANES for well-defined Ga(III)containing materials with known structures provide chemical insight into the large variability reported in their spectra. Comparison of the results using three different theoretical methods shows that their accuracy varies with the local environment, thus simulation results must be interpreted with caution. The large differences in edge energies for various Ga(III)-containing materials are reproduced, confirming that the XANES cannot be used as a simple fingerprint of the Ga oxidation state. Several descriptors, including the coordination number and symmetry of the local ligand environment, as well as the metal-ligand bond distances, correlate with spectral features because of ligand effects on the core ionization energy and the valence state energy.

The roles of such descriptors in predicting XANES features will be especially important in guiding the construction of models for active sites and reaction pathways in Ga-containing heterogeneous catalysts, whose structures are generally not well-known. We anticipate that using descriptors to interpret the spectra will be a more reliable approach than the common practice of remarking on the resemblance to the experimental XANES of reference compounds, such as the assignment of Ga(I) features by comparison to mixed-valent Ga(I)Ga(III)Cl₄.¹⁰

This study also highlights a critical need for a similar analysis of Ga(I) XANES, to identify descriptors that can distinguish between Ga(I) and Ga(III) sites that may be present in heterogeneous catalysts. We anticipate that such descriptors will be powerful inputs in data science and machine learning approaches, which show promise as computationally efficient ways to simulate experimental spectra.⁴⁰⁻⁴⁴

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Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/doi/xxx/xxx

Quantitative comparison between experimental and simulated spectra (Table S1, Figure S10), structures used for the computation of the XANES (Figures S1-S9), functional dependence of the computed XANES (Figures S11-S13), dependence of the computational spectra on Ph substitution, comparison between computed and experimental XANES for GaMe₃ (Figure S15), excited density of states for selected Ga complexes, with associated orbital visualizations (Figures S16-S18), crystallographic data for the DFT-based structures.

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