Design of semiconducting electrides via electron-metal hybridization: the case of Sc₂C

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Abstract:

Electrides are exotic materials that have electrons present in well-defined lattice sites. The existence of Y_2C and Gd_2C as 2D electrides inspired us to examine other trivalent metal carbides, including Sc_2C and Al_2C . It has been proposed that design rules for electride materials include the need for an electropositive cation adjacent to the electride site, but the effect of cation electronegativity on electronic structure in electride materials is not yet known. Here, we examine trivalent metal carbides with varying degrees of electronegativity and experimentally synthesize a 2D electride, Sc_2C , containing the most electronegativity of the cation drives greater hybridization between metal and electride orbitals. Our calculations predict that Sc_2C is a small band gap semiconductor with a band gap of 0.305 eV, with an experimental conductivity of 1.62 S/cm at room temperature. This is the first 2D electride material to exhibit semiconducting behavior, and we propose that electronegativity of the cation drives the change in band structure.

Introduction:

Challenges in energy storage, electronics, and catalysis motivate the search for exotic materials with extreme properties, and electrides—crystals with bare electrons trapped at stoichiometric concentrations¹⁻³—offer some of the most exceptional. These electrons have been ejected from atomic orbitals to reside in vacant lattice sites and, because they are so weakly bound, are better electron donors than alkali metals⁴⁻⁶, can offer electrical conductivity that rivals silver, and can catalyze challenging reactions. These properties have led to the exploration of electrides in applications where electron-rich materials are needed: N₂ and CO₂ reduction^{7,8}, battery electrodes^{9,10}, and electron emitters¹¹⁻¹³. Despite this progress, rules that might predict an electride's properties based on its structure or composition are underdeveloped. For example, unlike in conventional materials, it is unknown how to tune the band gap of semiconducting

electrides. And, even more critically, simple rules that could predict whether a material is an electride remain imprecise^{3,14}.

Presently, just two rules for predicting electrides are widely accepted: (1) the metal atoms that form the coordination sphere around the electride electron must be electropositive¹⁵, and (2) the compound must be electron rich¹⁶. For example, in $[Ca_2N]^+(e^{-})^{(17)}$, calcium is electropositive, and the compound is electron rich because the preferred oxidation states (Ca²⁺, N³⁻) and stoichiometry provide an extra electron, forming an electride. These rules are far too broad, however, because many compounds fulfill these rules but are not electrides.

To improve these rules, experiments have begun to identify the metals that can be adjacent to an electride electron. So far, only Ca⁽⁵⁾, Sr^(18,19), Ba^(17,20), Y^(21,22), Hf⁽²³⁾, and Gd^(24,25) have been incorporated next to the electride site (Figure 1), while Sc and several of the lanthanide series have been identified as candidates computationally^{26,27}. This delicate balance between electron counting and electronegativity was recently explored in the case of Ca₅Pb₃ (not an electride) and Ca₃Hf₂Pb₃ (an electride), as replacing calcium for hafnium makes the compound electron rich¹⁶. Despite these findings, the limit of cation electronegativity in the formation of electrides is not yet known. Therefore, to develop improved design rules for electrides and for their structure-property relationships, we explore the trivalent metal carbides with a layered structure. This family is informative because Y₂C is a known electride²², Sc₂C has been predicted²⁶, and Al₂C is unknown, and further, these metals vary systematically in electronegativity (Figure 1).

The question of electronegativity is essential because two experimentally observed electrides— Y_2C and Hf_2S —are "hybrid" electrides³, where the electride electron has some overlap with the adjacent metal orbitals, and electronegativity of the metal may be correlated with the degree of hybridization. This covalency plays an important role in defining the electronic structure of the electride because the electride electron is normally found around the Fermi level. Thus, as

Number of valence electrons				Unhybridized
One	Two	Three	Four	
Na 0.93	Mg 1.31	AI 1.61	Si 1.90	Hybridized
K 0.82	Ca 1.00	Sc 1.36	Ti 1.54	(e -))
Rb 0.82	Sr 0.95	Y 1.22	Zr 1.33	 Predictive Features 1. Compound electron count 2. Neighboring cation electronegativity 3. Electron-atom hybridization
Cs 0.79	Ba 0.89	RE	Hf 1.3	

Figure 1. An adapted periodic table highlighting the valency of various cations. The electronegativity values are shown beneath each element symbol, and the gray blocks highlight the cations that have been experimentally incorporated next to an electride site, which is shown pictorially on the right. The predictive features for potential electrides that we highlight in this work are listed. An explanation of electronatom hybridization is provided pictorially.

hybridization occurs, tunable band structures become possible, and indeed, Y_2C is a semi-metal. This is unique among the layered electrides and raises the question of whether greater degrees of hybridization between metal and electride electron could lead to new categories of electride materials²⁸. These questions motivate our exploration of the M₂C system, and especially of layered Sc₂C.

To highlight uncertainties within the Sc₂C system, we briefly outline previous findings and predictions. Sc₂C was originally synthesized in 1967 and assigned to the NaCl structure type, which we refer to below as the cubic phase²⁹. In 1969, different authors reassigned the structure as layered $(R\bar{3}m)$ on the basis of its powder x-ray diffraction (PXRD) pattern^{21,30}, even though the pattern was missing several important reflections (Table S1). In 1989, the structure was again reassigned as the cubic phase³¹ (Figure 2a, updated with our findings). This assignment appears to be correct and it is not known whether a layered Sc₂C electride exists. Currently, a 2D Sc₂C MXene structure has been explored computationally^{32,33} and calculations predict that layered Sc₂C would be an electride²⁶. In addition, 2D Sc₂C has been predicted to store an impressive 8% hydrogen by weight^{34,35} and reversibly store fluoride as a battery anode at capacities that exceed state-of-the-art lithium-ion batteries¹⁰. Thus, if a layered Sc₂C electride does exist, it could allow these exciting predictions to be realized.

In this study, we predict that layered Sc_2C is stable using an evolutionary algorithm and experimentally explore the phase diagram around the desired composition. We then explore the electronic structure of the synthesized material, and compare it to Y_2C and Al_2C to understand periodic trends. We find that as the cation becomes more electronegative, the band gap opens with Y_2C as a semimetal, Sc_2C as a small band gap semiconductor, and Al_2C as a larger gap semiconductor. The synthesized layered Sc_2C is the first 2D electride with semiconducting properties, opening up a fascinating new class of electride materials with several possible

applications in hydrogen storage and batteries.

Results and Discussion:

Phase diagram exploration:

Although layered Sc₂C is unknown, we and Hirayama *et. al.*²⁶ have hypothesized that it exists because Sc is chemically similar to Y and layered Y₂C is a stable phase. To investigate whether layered Sc₂C is thermodynamically stable, we used USPEX to explore the scandiumcarbon phase space. The phase diagram we obtained indicates that the Sc₂C composition is on the hull, and that the layered $R\overline{3}m$ structure is stable by 28 meV/atom (Figure 2b). This modest relative stability could suggest why the layered phase had not previously been synthesized.

The possibility that layered Sc_2C is thermodynamically stable motivated us to make this compound via hightemperature solid-state synthesis. In the Y_2C synthesis, we observed that high oxygen content in the yttrium led to formation of the cubic phase (YC_x). With



Figure 2a. The binary phase diagram of the Sc-C system adapted from Velikanova *et. al.* The layered Sc_2C has been incorporated (red). **b.** The hull diagram for the Sc-C phase calculated using USPEX. The conventional unit cell for three structures are shown corresponding to the colors on each plot.

this in mind, compared to the synthesis performed in Rassaerts *et. al.*'s 1967 study, we suspected that our use of high purity scandium (with a lower oxygen content) would allow us to experimentally realize our computational prediction. We arc-melted scandium and carbon in ratios from 2:0.9 to 2:1.15, respectively. These arc-melted samples had a high content of the cubic structure (ScC_x) (Figure S1), but arc-melting has a rapid cool-down, so the structure may not be at equilibrium. Therefore, we followed arc-melting with annealing for 7 days at 1100°C to increase the likelihood of observing the thermodynamically favored layered phase. We found that the cubic structure is stable both at compositions with deficient carbon (Figure S2) and excess carbon (Figure S3). However, at a composition of 2:0.99, the layered structure was observed, and this is reflected in the new binary phase diagram we propose for the Sc-C system (Figure 2a)³¹. There is uncertainty in the assignment of the temperature at which the congruent layered-to-cubic phase change occurs, but our experiments indicate that the layered phase is stable to at least 1100°C.



Figure 3. X-ray diffraction patterns of the experimental structure compared to the calculated layered and cubic Sc_2C phases. The layered and cubic crystal structures are shown, along with the amount present calculated from the Rietveld method. The predicted and experimental lattice parameters are compared.

То fully understand the structure of the layered phase, we first performed PXRD experiments of the product at room temperature and refined the pattern using the Rietveld method 36 (Figure 3). The pattern indexed as the $R\overline{3}m$ structure with lattice parameters of a = b = 3.345 (2) Å and c = 16.37 (2) Å. This refinement includes the incorporation of preferred orientation at {104} with a March coefficient of 0.754. The sample contained up to 15% of the cubic ScC_x. which we attribute to a relatively low stability of the layered phase. To quantitatively compare the predicted and experimental structures, we further optimized the structure from USPEX

using the PBEsol functional. The calculated lattice parameters were 0.45% smaller in *a* and 0.06% smaller in *c*, which is less than the typical error between PBEsol and experiment. This further confirms that the synthesized material is layered Sc_2C .

To better understand the effect of annealing in these samples, we performed transmission electron microscopy (TEM) of the product at room temperature under high vacuum. In Figure 4, we compare the Fast Fourier Transform (FFT) of these samples, where the non-annealed sample shows a ring pattern while the annealed sample shows well-defined spots. From the FFT images, we constructed filtered high resolution-TEM (HRTEM) images to create a color-coded map of crystalline domains where we compare the crystallinity of Sc_2C before and after annealing. The filtered HRTEM images and FFTs confirm an increase in domain size upon annealing, likely due to kinetic trapping from rapid cooling after arc-melting.

In addition, we used TEM to confirm that the layered Sc_2C phase is dominant. While from XRD there is a small amount of cubic Sc_2C (14.5%) compared to layered Sc_2C (85.5%), the

refinement has some error because the layered and cubic patterns are similar. In electron diffraction, the most common spatial frequency corresponds to a d-spacing of 2.3 Å, which we assign to either {104} in the layered structure or {101} in the cubic structure. This ambiguity motivated us to find distinct peaks for the layered phase. To identify the phase of these domains, we sought to acquire diffraction patterns containing the {117} reflection of the layered phase, which is expected at 5.501 nm⁻¹, for which there are no reflections in the cubic phase. In Figure 4c, we observe diffraction spots at 5.479 nm⁻¹, which can only be assigned to the layered phase. Simulations of the cubic structure with varied arrangements of carbon also failed to show a diffraction peak at 5.501 nm⁻¹. Importantly, the prevalence of {117} in diffraction patterns increases upon annealing, in agreement with XRD. In addition, XRD and TEM lattice spacings agree within 0.6%. This data confirms the experimental structure from XRD and corroborates the Rietveld refinement result.



Figure 4a. A high-resolution TEM image of Sc2C before annealing, with colored lattice planes corresponding to different crystallite regions determined through FFTs. **b.** A high-resolution TEM image of Sc2C after annealing for 7 days, with the lattice planes of the single crystallite region in green. **c.** A selected area electron diffraction (SAED) pattern of the (14,7,1) zone axis. The corresponding predicted and experimental d-spacings are compared.

Assessment of electride character and electronic structure:

Having confirmed the existence of layered Sc_2C as a thermodynamically stable material, we examined the material's electride character. In the experimental structure determined from our Rietveld refinement, there was a longer scandium-scandium distance (3.53 Å) compared to scandium metal (3.22 Å) and Sc_2O_3 (3.27 Å). This increased distance could indicate that electride electrons are present, so we sought evidence by calculating the electron localization function (ELF) using the hybrid DFT functional HSE-06. The ELF measures the degree of electron spatial localization compared to a reference electron with the same spin, and provides a map of electron probability across a system. The ELF of Sc_2C shows localization of electrons that are not



Figure 5a. The 3D ELF of Sc₂C, where the gradient of isosurface level spans from 0.3-0.6. This scale is the same for each ELF. **b.** The 3D ELF of Y₂C. **c.** The 3D ELF of Al₂C. **d.** The band structure and projected density of states of Sc₂C, where the valence and conduction bands are highlighted in pink and brown, respectively. **e.** The charge density of the Sc₂C valence band at the L point. The gradient of charge density spans from 0.017 e'/Å³ - 0.025 e'/Å³. This scale is the same for each charge density. **f.** The charge density of the Sc₂C conduction band at the L point. **g.** The charge density of the Al₂C valence band at the L point. **h.** The charge density of the Al₂C conduction band at the L point. **h.** The charge density of the Al₂C conduction band at the L point.

associated with either the scandium or carbon atoms, and indicates that there is electride character in the layered structure (Figure 5a).

We calculated the band structure and projected density of states (pDOS) with HSE-06 and found that Sc₂C is a semiconductor with a 0.305 eV indirect band gap and a 0.476 eV direct band gap (Figure 5d). The pDOS was calculated with a dummy atom of no charge to illustrate the contribution of the electride states to the overall density. This shows that the electride electron density mirrors that of the scandium and that the scandium orbitals are hybridizing with the electride electron density in the same way that has been observed in Y_2C and Hf_2S^3 .

We measured the electrical conductivity confirm these to predictions and further corroborate the match experiment between and Using AMSET calculations. the package³⁷. predicted we the conductivity would be 22.2 S/cm at 28.4°C using the band structure of Sc_2C . performed We conductivity measurements at the same temperature. where we used a well-insulated system to accurately measure temperature. We found that the conductivity was 1.62 S/cm, which is an order of magnitude smaller than the value we calculated; however, we expect the calculation to overestimate the conductivity because it does not include impurities in the crystal, which are present in experiment. This conductivity confirms that the synthesized material is indeed a semiconductor with a small band gap.

Emergence of semiconducting character in electrides

To understand how cation electronegativity plays a role in band structure, we compare these results to the electronic structure of Y_2C and Al_2C . The known hybrid electride Y_2C has a similar ELF to Sc₂C, where there is electron character that is not associated with either the yttrium or carbon atoms (Figure 5b). Yttrium has an electronegativity value of 1.22, whereas scandium has a value of $1.36^{(38)}$. Therefore, Sc₂C is the 2D electride with the most electronegative metal adjacent to the electride site yet synthesized. We wanted to further explore a possible electride structure with a similar, yet more electronegative metal, aluminum (1.61). A layered Al₂C has a similar structure to that of both Y_2C and Sc_2C with a space group of $R\bar{3}m$, and there is observed electron localization in this structure between the layers of Al₂C (Figure 5c). We calculated band structures and band gap of Y_2C (Figure S5), Sc₂C, and Al₂C (Figure S6) using the same calculation parameters. Y_2C is a semimetal, while Sc₂C and Al₂C are indirect band gap semiconductors, where the band gap increases from 0.31 to 0.50 eV as the electronegativity of the cation increases.

The metal-metal distances in these M_2C systems are all increased with respect to their metallic counterparts, and compared to one another, the Y-Y distance is the largest (3.96 Å), Sc-Sc is in the middle (3.53 Å), and Al-Al is the shortest (3.09 Å). Although one expects a different ionic radius for each cation, the metal-metal distances span such a large range that the differences in metal ionic radii are negligible; we conclude there is an increase in metal-electron hybridization from Y to Sc to Al. From these observations, we propose a model in which increasing electronegativity of metal leads to increased hybridization between unoccupied metal orbitals and the electride electron. This hybridization is because of similar energies and because of spatial proximity. Increased hybridization between electride state and empty orbitals leads to greater splitting between the filled valence band and empty conduction band – and thus a larger band gap. The relatively large amount of hybridization in the Al₂C structure would lead to a small amount of electron localization. Al₂C as a low electride character species supports the idea that neighboring cation electronegativity in electride materials likely lies somewhere just below the electronegativity value of aluminum.

If this model is correct, we should observe electride character in both the valence and conduction bands. Thus, we examined charge density at the top of the valence band and bottom of the conduction band in Sc₂C and Al₂C (Figures 5e-h). This indeed revealed electride character in both states, which could only result from hybridization of the electride electron and empty metal orbitals. This hybridization is analogous to a donor-acceptor interaction in molecules, such as BH₃-NH₃, but with the 2D electron layer playing the role of the donor and the 2D M₂C atoms playing the role of acceptor. Therefore, as overlap increases between donor (electride electrons) and acceptor (empty metal-carbide orbitals), the valence and conduction bands split more, and the band gap increases. Presumably, the greatest hybridization seen in the Al₂C system is facilitated by aluminum being most electronegative, since the empty Al orbitals are closer in energy to the electride electrons. This directly leads us to a powerful design rule for layered semiconducting electrides: that increasing metal electronegativity leads to increased electron-metal hybridization, yielding an increased band gap.

Exfoliation potential:

To assess anisotropy in the layered Sc₂C, we computationally predicted the work function using the thin slab method³⁹. We found that Sc₂C has a work function of 3.78 eV for the (0001) plane (Figure S7) and 2.89 eV for the (11 $\overline{2}0$) plane (Figure S8), which are comparable to the Y₂C values of 3.72 eV and 2.84 eV at (0001) and (11 $\overline{2}0$), respectively²². The change in work function in different directions shows the anisotropic nature of Sc₂C. The similar work functions of Sc₂C and Y₂C support their similarity, but Sc₂C has a slightly larger work function, consistent with greater electronegativity of scandium. Nevertheless, Sc₂C is still a good reducing agent. To explore anisotropy in a different way, we calculated the interlayer and intralayer stiffness tensors of Sc₂C to be 136.7 GPa and 286.8 GPa. Compared to Y₂C, these values indicate a greater degree of anisotropy because the intralayer stiffness tensor of Sc₂C is 110% greater than its interlayer counterpart, whereas the intralayer tensor of Y₂C is only 44% larger⁴⁰ (Table S5).

Additionally, we calculated the energy to separate two monolayers of Sc_2C (i.e. binding energy). We obtained a value of 2.76 J/m², which is larger than that of Y_2C^{41} . Compared to other known layered electrides with similar structures (Ca₂N and Sr₂N), they exhibit a similar trend where the electride with the smaller cation, Ca₂N, has a larger binding energy than Sr₂N. This is likely due to the smaller cation electride having a smaller interlayer distance, giving a higher electrostatic attraction between the layers. The larger binding energy of the carbide systems overall is likely due to additional covalency, which is present in the nitride systems to a lesser degree.

Finally, we calculated the mechanical properties of Sc_2C as compared to Y_2C (Table S6). We found that Sc_2C has a higher bulk modulus than Y_2C , while it's shear modulus and Young's modulus are both lower. This indicates that Sc_2C is more resistant to compression than Y_2C ,395.6 but its strain response to stress is much higher. The higher bulk modulus, lower shear modulus, and higher degree of anisotropy indicate that the material has a higher exfoliation potential than Y_2C . We explored this experimentally, and found that the material wouldn't exfoliate after sonicating in 1,3-dioxolane for 800 minutes. SAED patterns and TEM images show that no nanosheets or aggregates are present (Figure S9). Nevertheless, experimental exfoliation could lead to realization of many exciting applications proposed for 2D Sc_2C , and these calculations predict that it should be explored further.

Stability in air:

Finally, we tested the stability of this novel material in air and further explored the effect of oxidation on the crystal structure of Sc₂C. We performed PXRD experiments, taking data at various timepoints after letting the product sit in air (Figure 6). The data shows that the material oxidizes slowly, with almost no change after sitting in air for one day. However, after four days, there was a shift to higher angles, decreased intensity, and increased width. These attributes indicate decrease а in crystallinity more as oxygen incorporates. This correlates with



Figure 6a. X-ray diffraction patterns after various amounts of oxidation. **b.** The diffraction pattern zoomed in between 54-56° two-theta. **c.** The diffraction pattern zoomed in between 32.4-33.4° two-theta.

formation of an amorphous oxide, which is the dominant phase by the final timepoint, where there is little intensity throughout the diffraction pattern. The particle size is, on average, greater than 1 μ m (Figure S9), indicating that the oxidation process isn't self-limiting.

To better understand the mechanism of oxidation, we performed x-ray photoelectron spectroscopy (XPS). High-resolution scans of scandium and carbon were taken at 0 days, 1 day, and 11 days sitting in air (Figure 7). Scandium shows four species at the initial timepoint. indicating several scandium species were present. Component A corresponds to $(Sc^{3+})^{42}$, likely due to surface oxidation during transfer into the XPS. We observe one species (component B) at binding energy that were previously assigned to $(Sc^{2+})^{43}$. We assume that component C is an intermediate species as the carbide transitions to the fully formed oxide, with peaks that lie between the component A and B doublets. Finally, there is a small peak with a low binding energy (395.6 eV), where it is unclear whether this is scandium or possibly nitrogen as the binding energy of peak is low for either element. There are four species at the initial timepoint for carbon, where component E and G are adventitious carbon and its oxidized counterpart, respectively⁴⁴. Component F is likely the carbide species, as peaks at a similar shift have been seen in other transition metal carbide species⁴⁵. The carbide peak energy in this material is relatively reduced compared to other known transition metal carbide species, which is not surprising from an electride. Component H is likely an intermediate species as the carbide oxidizes. After one day of exposure to air, the Sc^{2+} peak decreases, the Sc^{3+} peak increases, and the carbide carbon peak decreases in intensity, indicating the sample is oxidizing. In the carbon spectra, a new species,



component I, becomes apparent, and likely results from the adventitious carbon oxidizing. At the 11-day timepoint, the peaks associated with Sc₂C in both the scandium and carbon spectra have disappeared, indicating the surface was completely oxidized. The XRD data shows that there are still crystalline peaks present at the 11-day timepoint. From XRD and XPS, we conclude that an amorphous oxide forms on the surface of crystals first, and as the material continues to be exposed to air, oxygen diffuses further into the crystals until the bulk material becomes an amorphous oxide after 28 days. Therefore, handling material in air will this require passivating coatings.

Figure 7. X-ray photoelectron spectroscopy spectra of scandium and carbon as-made, after 1 day in air, and after 11 days in air.

Conclusions

We have synthesized a new 2D inorganic electride, Sc_2C , updating the previously published Sc-C phase diagram. This electride includes the most electronegative cation adjacent to the electride site. We determine that while there needs to be an electropositive cation present in the structure next to the electride site, the upper limit of electronegativity likely falls just above that of aluminum. We find that this material is the first synthesized small band gap semiconductor 2D electride. Along with our exploration of Al₂C, the trivalent metal carbides lead to a new design rule for semiconducting 2D electrides: that more electronegativity of the cation, the band gap in these electride materials can be tuned. The small band gap and good conductivity also allow applications as a battery electrode or in IR photodetectors. Semiconducting electrides under photoexcitation could exhibit both interesting reduction and oxidation capabilities, which could be useful for photocatalysis or catalysis. The experimental realization of layered Sc_2C will allow predictions of its enormous H₂ and F⁻ storage capacities to be explored.

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Supporting Information Available

Additional details regarding the literature search, computational and experimental methods, and experimental observations.

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