Boosting the Efficiency and Stability of Blue TADF Emitters by Deuteration

3 Tianyu Huang¹, Dongdong Zhang¹*, Ge Zhan², Qi Wang¹, Guomeng Li¹, Xiangchen Hong¹,

4 Ziyang Liu¹, Lian Duan^{1,3}*

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6 ¹ Key Laboratory of Organic Optoelectronics, Department of Chemistry, Tsinghua University,

8 ² Beijing National Laboratory for Molecular Sciences (BNLMS), State Key Laboratory of Rare

9 Earth Materials Chemistry and Applications, College of Chemistry and Molecular Engineering,

10 Peking University Beijing, 100871, China

³ Center for Flexible Electronics Technology, Tsinghua University, Beijing 100084, China.

12 E-mail: ddzhang@mail.tsinghua.edu.cn; duanl@mail.tsinghua.edu.cn

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⁷ Beijing 100084, China

17 Abstract:

Deep-blue organic light-emitting diodes (OLEDs) with both high efficiency and stability remain 18 exclusive in scientific literatures after decades of research. A key reason is the chemical bond 19 dissociation through conical intersections between potential energy surfaces (PESs) of 20 energetically-hot excited states and dissociative states. Here, we report a deuteration strategy to 21 stabilize blue thermally activated delayed fluorescence (TADF) emitters. It is unveiled that 22 deuteration would lower high-frequency vibrations to create shallower PESs in both ground and 23 excited states, energetically retarding the cross of the conical intersection point in dissociative 24 process and alleviate vibrational coupling to eliminate nonradiative decay. With a deuterated 25 blue TADF compound, namely 2,3,4,5,6-pentakis(9H-carbazol-9-yl-d8) benzonitrile, as the 26 27 sensitizer, a deep-blue OLED simultaneously achieves a peak external quantum efficiency of 30.3%, a Commission Internationale de L'Eclairage (CIE) coordinate y value of 0.17, and a 28 superb LT80 (time to 80% of initial luminance) of 227 h at 1000 cd/m². And a blue OLED with a 29 $CIE_{\nu} \sim 0.19$ shows an even longer LT80 of 456 h at 1000 cd/m². 30

31 Introduction

Being one of the three primary colors, blue is indispensable for displays and lightings. However, 32 efficient and stable blue organic emitters for organic light-emitting diodes (OLEDs) have 33 puzzled the researchers for decades. As an OLED ages during continuous operation, a small 34 portion of its molecules will chemically degrade over time to form defect sites, which then act as 35 luminance quenchers, nonradioactive recombination centers and deep charge traps, finally 36 arousing the luminance loss.¹ The formation of defects has been convincingly assigned to the 37 annihilation between bimolecular excited states, finally giving birth to an energetically-hot 38 excited states to dissociate molecule bonds. The spin-allowed symmetric singlet excitons (S_1) 39 feature rapid radiatively decay process to compete the annihilations and thus fluorophores that 40 can only harness S₁ exhibit high stability but low efficiency.^{2,3} Conceptually advanced emitters, 41 including noble-metal-based organometallic phosphors^{4,5} and thermally activated delayed 42 fluorophores (TADF).⁶⁻⁸ have been developed to recycle the spin-forbidden unsymmetric triplet 43 excitons. The long-lived triplets, however, are more likely to participate into the annihilation 44 progress.9 Thus, though ultimate high device efficiency, phosphorescent and TADF OLEDs 45 always suffer from worse stability than fluorophores, particularly in the blue region. 46

This invokes a new issue, that is, if the lifetime of blue devices that take fully use of the triplet excitons can be promoted? Principally, the most straightforward way is to reduce the triplet concentration and/or shorten its lifetimes. With this in minds, progresses have been made in sky-blue OLEDs. Forrest and co-workers successfully achieved tenfold increase in the lifetime 51 of sky-blue phosphorescent OLEDs utilizing a graded-dopant-concentration profile to lower triplet concentrations.¹⁰ And later by introducing a manager to quickly consume the hot excited 52 states, the device stability was further improved up to LT80 (LTx means time to x% of initial 53 luminance) of 334 h but only moderate external quantum efficiency (EQE) of ~9% and a 54 Commission Internationale de L'Eclairage coordinate y value (CIE_v) of 0.31 at 1000 cd/m². As 55 for TADF OLEDs, a fast reverse intersystem crossing (RISC) process is believed to be effective 56 in consuming excited states to avoid annihilations. This goal can be fulfilled by introducing 57 hetero-donors to enhance the mixing of excited states with different nature, realizing superb 58 LT97 of 110 h with CIE_v of 0.44 and EOE of 20.8% at 1000 cd/m^{2.11} Cui *et al.* also developed 59 sky-blue TADF emitters with multiple donor units to enlarger spin-orbit couplings and realized a 60 large RISC rate (k_{RISC}) of ~10⁷ s⁻¹, accounting for a long LT90 of ~600 h, an EQE of 28.6% at 61 1000 cd/m^2 and an emission peak at 486 nm, respectively.¹² 62

63 Albeit these progress, deep blue devices with $CIE_{\nu} < 0.2$, however, still lag behind the sky-blue ones. The difficulty arises from the fact that the energy of hot excited-states can be extremely 64 high in deep-blue emitting systems and thus being more conspicuous to deteriorate the device 65 stability. To circumvent this issue, a sensitized fluorescence emission mechanism has been 66 proposed with blue phosphors or TADF materials as sensitizers to assist triplet harvest for stable 67 fluorescent dopants through Förster energy transfer (FET) process. The long-range FET 68 interaction with extremely high rates of $10^8 - 10^{11}$ s⁻¹ enables guick consumption of triplet 69 excitons to suppress triplet-related annihilations and improve device stability.¹³ Recently, 70

multiple resonance (MR) fluorescent emitters with narrow bandwidth emission have been
adopted as the final emitters, offering a purer emission color with relatively lower excited state
energy. Deep blue Ir(III) and Pt(II) phosphors have been adopted as sensitizer for a well-known
MR emitter,

N7,N7,N13,N13,5,9,11,15-octaphenyl-5,9,11,15-tetrahydro-5,9,11,15-tetraaza-19b,20bdiboradin 75 aphtho[3,2,1-de:1',2',3'-jk]pent-acene-7,13-diamine (v-DABNA), realizing extended LT50s of 76 121 h (CIE_v of 0.162, EQE of 23.3%) and 156.1 h (CIE_v of 0.141, EQE of 25.4%) at an initial 77 luminance of 1000 cd/m², respectively.^{14,15} By using v-DABNA as the final emitter, 78 TADF-sensitized-fluorescence (TSF) OLEDs achieved LT50 of 151 h (CIE_v of 0.17, EQE of 79 25.2%) and LT95 of 11 h (CIE_v of 0.20, EQE of 20%) respectively also at an initial luminance of 80 1000 cd/m^2 , comparable or even better than the phosphor sensitized ones.^{16,17} Those preliminary 81 works testify the viability of TSF in breaking the efficiency-stability limitation of deep-blue 82 OLEDs with metal-free all purely organic molecules, which, however, still requires the 83 development more stable blue TADF sensitizers. 84

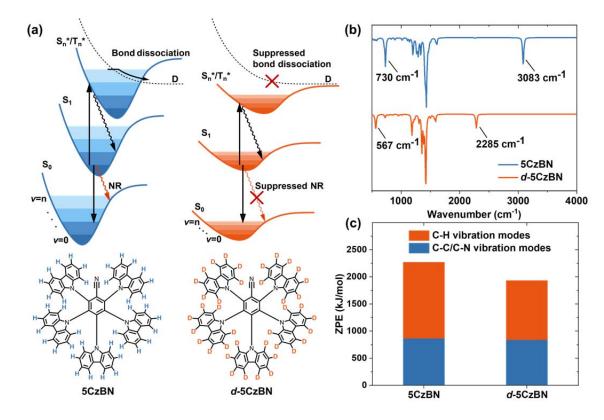
Deuteration had been reported to be an effective way to improve the stability of conventional fluorescent, phosphorescent emitters and host materials as well.¹⁸⁻²¹ However, deuteration effect on TADF emitters remain unexplored and the underlying mechanism is not clear. Herein, we studied deuterated blue TADF emitters with the aim to improve the efficiency and stability. It is revealed that deuteration can lower the energy of vibrational levels establish shallow potential energy surface (PES) by in both ground and excited states. This shallow PES hinders the cross of

91 the conical intersections between the PESs of the hot excited states and the dissociative states and thus suppress molecular dissociation. Also, vibrational coupling via high-energy vibration 92 levels can be alleviated, thus eliminating nonradiative decay. The proof-of-the-concept 93 deuterated blue TADF emitter exhibited both improved electroluminescence stability and 94 efficiency compared with the undeuterated one. And with this deuterated blue TADF compounds 95 as sensitizers, deep-blue OLEDs simultaneously displayed a maximum EQE of 30.3%, a CIE_v 96 value of 0.17, a full-width at half-maximum (FWHM) of 27 nm and a superb LT80 of 227 h at an 97 initial luminance of 1,000 cd/m². Using a redshifted final emitter with a FWHM of only 19 nm, 98 blue device with CIE_{ν} of 0.19 can even realize a LT80 of over 450 h at an initial luminance of 99 1.000 cd/m^2 . This work may shed new light on the development of highly efficient and stable 100 deep-blue emitters. 101

102 **Results**

103 Molecular design and theoretical analysis

Fig. 1a shows the Jablonski diagram of the possible relaxation pathways for excitons. The hot state can be either an exciton or polaron state resulting from either exciton-exciton annihilation or exciton-polaron annihilation, respectively. While most hot states rapidly relax to the lowest excited states, those that have sufficient higher energy than the dissociative states can lead to the chemical bond dissociation. Technically, this dissociation process happens at the conical intersection point between PESs of hot singlet/triplet excited states (S_n*/T_n*) and dissociative states. Therefore, the energy of S_n*/T_n* should reach the conical intersection point to induce 111 bond dissociation. Besides the electronic energy level, vibrational energy levels also matter in this process, though having been neglected in previous reports. Emitters with high vibrational 112 energy levels, that have relatively deep PESs, are more susceptible to reach the energy point of 113 conical intersection and induce dissociation process. Alternatively, if the S_n*/T_n* states possess 114 shallow PESs, *i.e.*, low vibration energy levels, the chances to reach the energy crossing point of 115 conical intersection of dissociation process may, theoretically, be significantly reduced and thus 116 suppressing the possibility of bond dissociation. Besides a good long-term stability, a high 117 efficiency is equally important for TADF emitters. The vibrational coupling between high energy 118 vibration levels of the ground state (S₀) and the low energy vibration level of S_1/T_1 has been 119 120 convincingly recognized as main reason for nonradiative decay. It is reasonable to expect that a shallow PES should also favor to alleviate vibrational coupling to eliminate nonradiative decay. 121



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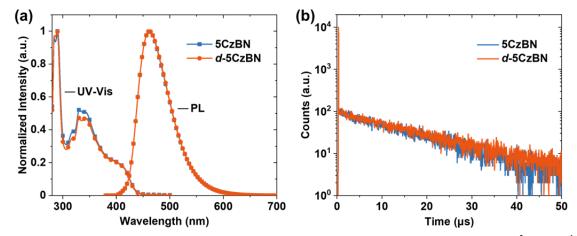
Figure 1 (a) Jablonski diagram of the possible relaxation pathways for excitons and chemical structures of
5CzBN and *d*-5CzBN. (b) Simulated FT-IR spectra of 5CzBN and *d*-5CzBN. (c) Contributions of different
vibration modes to molecular ZPE of 5CzBN and *d*-5CzBN.

126 In the aromatic light-emitting π -systems, high frequency vibration usually originates from C-H 127 stretching and bending. Previous kinetic isotope effect has revealed that replacing hydrogen atom with deuterium (D) can significantly alleviate those high frequency vibrations.²² Therefore, 128 based on the above discussion, deuterated TADF emitters possess the potential to achieve 129 130 superior stability and high efficiency simultaneously. То validate this idea, 2,3,4,5,6-penta(9H-carbazol-9-yl) benzonitrile (5CzBN) was chosen as a template molecule. 131 This multiple-donors compound is one of the widely studied sky-blue TADF emitters.²³ By 132 replacing carbazole (Cz) units with deuterated carbazole, we synthesized the deuterated analog 133 of 5CzBN, namely *d*-5CzBN. To elaborate the isotope effect of *d*-5CzBN, we first carried out 134

135 theoretical calculations. Commonly, quantum chemistry programs are designed based on Born-Oppenheimer approximation, which means the mass of nucleus is not taken into account 136 for the electron structure calculation. In other words, the electronic behaviors including 137 excitation states are not bothered by deuterium effect in theoretical calculation. As a result, the 138 calculated electronic transition properties for both 5CzBN and d-5CzBN are the same and we 139 will verify that the two molecules do exhibit identical S_1 and T_1 energy state in the following 140 content. However, the results of frequency analysis and thermodynamics are affected by isotope 141 setting since the calculations are based on mass-weighted coordinate system. Fig. 1b displays the 142 simulated FT-IR spectrum of both emitters, clearly revealing that the highest energy of observed 143 absorption band for C-H stretching modes is shifted from \sim 3083 to \sim 2285 cm⁻¹ and C-H bending 144 modes shifted from \sim 730 cm⁻¹ to \sim 567 cm⁻¹ after deuteration. The same phenomenon was also 145 recorded from the experimental result shown in Supplementary Fig. 1, evidencing that 146 147 deuteration could effectively suppress the high frequency vibrations of molecules. Also, it is interesting to note that, for other low frequency vibration, deuteration exhibits quite limited 148 influence. Other evidence that proves the alleviated high frequency vibrations should be the 149 150 zero-point energy (ZPE) value of those two compounds. It is found that replacing proton with deuterium can significantly reduce the ZPE of the molecule. And the contribution of different 151 152 normal vibration modes to ZPE of the molecules was analyzed by Shermo program as shown in Fig. 1c.²⁴ The vibrations from aromatic skeleton including carbon-carbon (C-C) and 153 carbon-nitrogen (C-N) bonds contribute similar ZPE value before and after deuteration and the 154

decreased ZPE value from 5CzBN to *d*-5CzBN mainly originate from the lower frequency vibrations of C-D than C-H. Based on those evidence, it can be concluded that compared with 5CzBN, *d*-5CzBN should possess relatively shallower PES since its high frequency vibrations are suppressed, which is exactly the aim of this molecular design strategy.

159 Material characterization



161 Figure 2 (a) Absorption and PL spectra of 5CzBN and *d*-5CzBN in toluene solution $(10^{-5} \text{ mol } \text{L}^{-1})$. (b) **162** Transient PL spectra of 5CzBN and *d*-5CzBN in toluene solution $(10^{-5} \text{ mol } \text{L}^{-1})$.

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Table 1 Photophysical properties of *d*-5CzBN and 5CzBN.

			HOMO ^{b)} (eV)	LUMO ^{c)} (eV)	$ au_{p}^{d)}$	$ au_d^{(d)}$	$\Phi_{PL}{}^{d)}$	$k_{\rm r}$ (×10 ⁷ s ⁻¹)	$\frac{k_{\rm ISC}}{(\times 10^8 \text{ s}^{-1})}$	k_{RISC} (×10 ⁵ s ⁻¹)	k_{nr}^{T} (×10 ⁴ s ⁻¹)
d-5CzBN	2.91	2.66	5.49	2.66	5.3	11.4	0.95	2.26	1.66	6.89	0.50
5CzBN	2.91	2.66	5.48	2.65	4.9	9.9	0.78	2.24	1.82	6.81	2.61

 ^{a)} Measured in the toluene solution with a concentration of 10⁻⁵ mol/L. ^{b)} Calculated by LUMO and energy gap measured by UV-vis. ^{c)} Measured in dry N, N-dimethylformamide with a concentration of 10⁻³ mol/L.
 ^{d)} Measured in the N₂-saturated toluene solution with a concentration of 10⁻⁵ mol/L.

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To further uncover the effect of deuteration, electrochemical and photophysical properties of 5CzBN and *d*-5CzBN are compared in detail. As shown in **Supplementary Fig. 2**, LUMO energy levels of 5CzBN and *d*-5CzBN were estimated by cyclic voltammetry, being -5.49 and -5.48 eV, respectively. Also, identical optical energy gaps can be obtained for both emitters (2.83
eV) from their ultraviolet-visible (UV-vis) absorption spectra measured in toluene (10⁻⁵ mol L⁻¹)
as shown in Fig. 2a. HOMO energy levels of -2.66 eV for 5CzBN and -2.65 eV for *d*-5CzBN
thereof can be obtained. The almost same energy level of frontier molecular orbitals indicate that
deuterium substitution has nearly no influence on electron configuration.

178 Interestingly, identical photoluminescence (PL) fluorescent and phosphorescent spectra were recorded at room temperature and 77K of 5CzBN/d-5CzBN in toluene solution as illustrated in 179 Fig. 2a and Supplementary Fig. 3, respectively. The S_1 and T_1 energies of 5CzBN and 180 *d*-5CzBN were measured to be exactly the same from the onset and the peak wavelengths, being 181 2.91 and 2.66 eV, respectively. Those performances suggest that deuteration does not affect the 182 radiative decay energy levels. This is understandable that the radiative decay energy levels are 183 mainly determined by the relatively low vibration energy levels, which were not influenced by 184 185 the deuteration as proved above. Actually, the maintained PL spectra has been widely observed for deuterated organic compounds.¹⁸ 186

On the contrary, a photoluminescence quantum yield (PLQY) of 0.95 was obtained for *d*-5CzBN in the degassed solution while only 0.78 was recorded for 5CzBN. This suggests the significantly suppressed nonradiative decay process in *d*-5CzBN. The PL decay curves in the solution was also recorded, revealing a prompt/delayed lifetime of 5.3 ns/11.4 μ s and 4.9 ns/9.9 μ s for *d*-5CzBN and 5CzBN, respectively. The TADF dynamic processes, involving rates of S₁ radiative decay, intersystem crossing, reverse intersystem crossing and T₁ nonradiative decay, have been calculated and listed in **Table 1**. Interestingly, k_r and k_{RISC} are quite similar for both emitters, suggesting that the deuteration would not influence the radiative decay and RISC processes. On the contrary, only a one-fifth k_{nr}^{T} was observed for *d*-5CzBN compared with 5CzBN, evidencing that deuteration favors to suppress the nonradiative decay of T₁.

Unfortunately, the present calculation method cannot distinguish nonradiative decay and ISC 197 process from S_1 .⁹ Therefore, the calculated k_{ISC} utilizing the commonly adopted method should 198 involves both nonradiative decay rate of S₁ (k_{nr}^{S}) and ISC rate (k_{ISC}^{S}). The influence of external 199 factors on ISC and RISC are usually similar, and it is reasonable to assume that $k_{\rm ISC}$ ^S is also not 200 influenced by deuteration. Therefore, the different in $k_{\rm ISC}$ should originates from the different $k_{\rm nr}^{\rm S}$ 201 of both emitters. And it is clear that k_{nr}^{S} in 5CzBN is more significant than that of *d*-5CzBN. 202 Therefore, deuteration should suppress nonradiative decay process of both S₁ and T₁, giving birth 203 to both longer τ_p and τ_d as well as relatively higher PLQY of d-5CzBN. Those results herein 204 205 validate our inspirations that deuteration induced shallow PES favors suppressed nonradiative decay. 206

The photophysical properties of the two emitters were also compared in doped films with a host material named 9-(3-(9H-carbazol-9-yl)phenyl)-9H-3,9'-bicarbazole (mCPBC). As exhibited in **Supplementary Fig.4**, with the increasing doping concentration in mCPBC, both emitters showed red-shifted emission peak. Interestingly, d-5CzBN film exhibited an increased FWHM than that of 5CzBN film. This phenomenon can be attributed to the difference in molecular interaction that leads to different conformations of d-5CzBN and 5CzBN in solid film. As a 213 reference, single crystals of d-5CzBN and 5CzBN were prepared and analyzed. With the structures shown in Supplementary Fig. 6, *d*-5CzBN exhibits smaller dihedral angles between 214 Cz units and central benzene ring compared with 5CzBN. And the difference we believe should 215 arise from the different volume of H and D atom.²⁵ The S₁ energy of *d*-5CzBN and 5CzBN based 216 on their crystal structure have been calculated by TDDFT calculation, exhibiting higher energy 217 values for *d*-5CzBN (Supplementary Table 1). The topologies of the molecular packing are also 218 different from each other (Supplementary Fig. 7). Thus, it is understandable that the different 219 molecular interaction should accounts for the different of PL spectra in doped films. The PL 220 decay curves of the doped films were also measured, displaying relatively longer prompt and 221 222 delayed lifetimes of *d*-5CzBN than 5CzBN (Supplementary Fig. 5). Combining with the higher 223 PLQY of *d*-5CzBN, the different in exciton lifetimes should also be attributed to the significantly suppressed nonradiative decay process in d-5CzBN, which is in agreement with the results 224 225 obtained from the solution.

226 **Device performance**

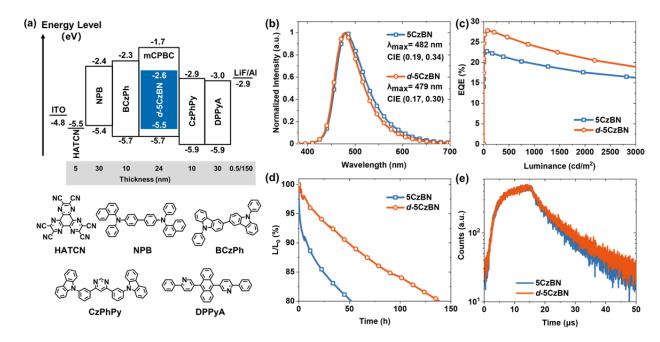




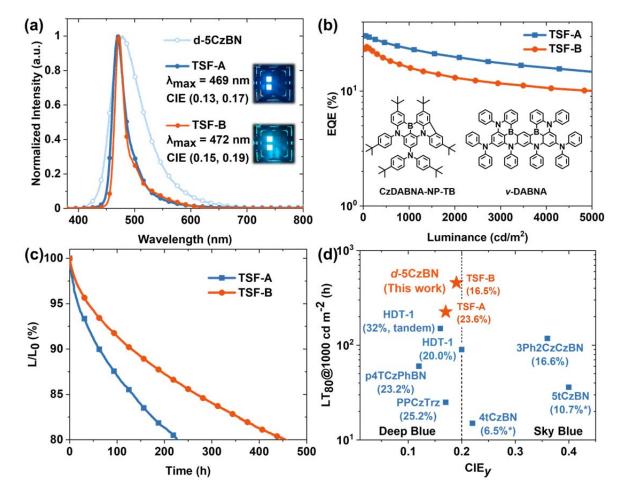
Figure 3 (a) Device configuration and molecular structures of the materials used in devices. (b) EQE vs.
 luminance characters of devices. (c) EL spectra of devices. (d) Lifetime of devices (at an initial luminance of 1000 cd m⁻²). (e) EL transient spectra of devices.

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The electroluminance (EL) performances of 5CzBN and d-5CzBN were compared with device 232 233 structures of ITO/ HATCN(5 nm)/ NPB(30 nm)/ BCzPh(10 nm)/ mCPBC: 20 wt% emitter (24 nm)/ CzPhPy (10 nm)/ DPPyA (30 nm)/ LiF (0.5 nm)/ Al (150 nm). Fig. 3a provides the energy 234 levels and chemical structures of the materials adopted in the device while Fig. 3b exhibits the 235 EL spectra with both emission peaks around 480 nm. Also, similar to the PL spectra of doped 236 237 films, a relatively narrow emission band was obtained for d-5CzBN. Consequently, a relatively bluer CIE coordinate of (0.17, 0.30) was recorded for *d*-5CzBN based device and (0.19, 0.34) for 238 239 device with 5CzBN. The EQE-luminance characters were provided in Fig. 3c, exhibiting a high maximum EQE of 27.8% for d-5CzBN device while a moderate EQE_{max} of 22.7% for 5CzBN 240 device. The relatively higher device efficiency of d-5CzBN naturally should be assigned to its 241

higher PLQY. Also, a relatively low efficiency roll-off was observed for *d*-5CzBN device with EQE remain 24.2% under luminance of 1000 cd/m². A power efficiency as high as 57 lm/W was also obtained for *d*-5CzBN, higher than that of 5CzBN device (53 lm/W).

The operational stabilities of the two devices were also measured at an initial luminance of 1000 245 cd m⁻², of which the results were provided in Fig. 3d. A LT80 of 140 h was recorded for 246 *d*-5CzBN device, which is almost 3-times longer than the 5CzBN-based device (LT80 = 53 h). 247 To reveal the origin of the significantly improved device efficiency and stability, the EL transient 248 decay curves of the two devices were illustrated in Fig. 3e. Similar to PL transient decay curve in 249 mCPBC film, device based on 5CzBN shows a faster decay which means that the device suffers 250 more from non-radiative decay process than *d*-5CzBN device. More importantly, although 251 relatively long exciton lifetime of d-5CzBN device, which should induce more significant 252 exciton annihilations for energetically hot excited states, the 3-fold increasement in LT80 253 254 suggests the good endurance of *d*-5CzBN under excited states, which should arise from shallow PES as aforementioned. Those results validate our inspiration that deuteration induced shallow 255 PES favors high stability under excitation. 256



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Figure 4 (a) EL spectra of the TSF-A, TSF-B and *d*-5CzBN devices; (b) EQE vs. luminance characters of the TSF-A and TSF-B device. (Inset: molecular structure of CzDABNA-NP-TB and v-DABNA). (c) Lifetime of the TSF-A and TSF-B device (at an initial luminance of 1000 cd m⁻²). (d) A comparison of device lifetime with different CIE_y values for stable blue TADF OLEDs. EQE at 1000 cd m⁻² were labeled in the parenthesis. (*EQE at 500 cd m⁻²)

Though the significantly improved device lifetimes, *d*-5CzBN based devices only showed sky-blue emission. Strategy to realize deep blue devices is utilizing *d*-5CzBN as a TADF sensitizer while introducing an emitter with a narrowband emission. Here, a previously reported MR-TADF dopant, CzDABNA-NP-TB, was adopted as the deep blue emitter.²⁶ As illustrated in **Supplementary Fig.10**, this compound shows a narrow band emission peaked at 464 nm with a

269	FWHM of 21 nm and a PLQY of ~100% in toluene solution (10^{-5} mol L ⁻¹) and a narrow sharp
270	absorption peak at 450 nm was also observed. Owing to this small stokes shift, a significant
271	spectra overlap between the wide emission spectrum of sensitizer and the absorption spectrum of
272	dopant can be observed, giving birth to a large radii of FET of 3.62 nm. Efficient energy transfer
273	thereby can be expected from sky-blue TADF sensitizer to deep-blue narrowband emitter. The
274	structure of this TADF-sensitized fluorescent device (TSF-A) was ITO/ HATCN (5 nm)/ NPB
275	(30 nm)/ BCzPh (10 nm)/ mCPBC: 20 wt% d-5CzBN: 1 wt% CzDABNA-NP-TB (24 nm)/
276	CzPhPy (10 nm)/ DPPyA (30 nm)/ LiF (0.5 nm)/ Al (150 nm).

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278 **Table 1** Summary of device performances

Device	λ_{EL} V _{on} (nm) (V)		EQE (%) Max 1000 cd/m ²		Power Efficiency (lm/W) Max 1000 cd/m ²		FWHM (nm)	CIE (x, y)	LT80 ^{a)} (h)
5CzBN	482	2.8	22.7	19.5	53.1	34.2	72	0.19, 0.34	53
d-5CzBN	479	2.8	27.8	24.2	57.3	38.6	69	0.17, 0.30	140
TSF-A	469	2.8	30.3	23.5	50.9	24.1	26	0.13, 0.17	227
TSF-B	472	2.6	24.3	16.5	35.1	16.6	19	0.15,0.19	456

 a^{a} Recorded at an initial luminance of 1000 cd m⁻².

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281 The EL spectra of the TSF devices were shown in **Fig. 4a**, showing strong emission from dopant.

282 Interestingly, owing to the extremely small FWHM of 26 nm, a deep-blue emission with CIE

coordinate of (0.13, 0.17) was observed, even though the emission peak is located at a relatively

long wavelength of 469 nm. This evidence the advantages of narrow band emitter to obtain

285 deep-blue emission. Fig. 4b showed the EQE-luminance characters of the TSF devices,

exhibiting a remarkable EQE_{max} of 30.3% which remains 23.5% at practical brightness of 1000 cd m⁻². The EQE_{max} significantly outperforms that of the device with *d*-5CzBN as emitter. Besides the fully exciton utilization assisted by sensitizer, the high EQE value can be assigned to the enhanced light outcoupling efficiency benefiting from the preferred horizontal orientation of dopant due to its planar structure.²⁷ Also, a high maximum power efficiency of 50.9 lm/W was obtained. Both the EQE and power efficiency represent one of the highest values among all reported results with similar colors.

The device stability was further testified at an initial luminance of 1000 cd m⁻² as illustrated in 293 Fig. 4c. Remarkably, a superb LT80 of 227 h was obtained. Detailed device data are shown in 294 Table 2. Interestingly, compared with device directly utilizing d-5CzBN as emitter, the TSF 295 device showed even longer lifetime with blue-shifted emission. This is reasonable given the fact 296 that the emitter with narrow band though exhibits bluer emission possesses even lower exciton 297 298 energy, defined by the onset energy of fluorescent spectra. Also, the potentially accelerated exciton consumption arising from the rapid radiative decay process of CzDABNA-NP-TB also 299 favors longer lifetimes. 300

Furthermore, we also adopted a well-known MR dopant, *v*-DABNA, as a final emitter to enhance device stability.²⁸ Featuring extremely narrow bandwidth, *v*-DABNA has exhibited good long-term stability in devices based on TSF. Device named TSF-B was fabricated with the same device architecture above. The EL spectrum showed an emission peak at 472 nm with a small FWHM of only 19 nm, corresponding to a CIE_v of 0.19. Though a relatively lower EQE_{max} 306 of 24.3%, a significantly extended long LT80 of 456 h was observed at an initial luminance of $1,000 \text{ cd/m}^2$, which also represents the longest lifetime under this specific color. We summarized 307 recent reported stable bottom emission blue devices with $CIE_v < 0.4$ in Fig. 4d.^{16,17,23,29,30} Clearly, 308 our results here present a significant improvement in device lifetime, together with the 309 state-of-the-art device efficiency, validating our deuteration strategy to stablilize blue 310 compounds. It is anticipated that by optimizing host materials and device structures such as 311 top-emitting devices, both efficiency and stability can be further improved accompined by even 312 blue-shifted emission. 313

314 **Discussion**

Highly efficient and stable deep-blue OLEDs have continuously been an obstacle for further 315 development of this technology and still face formidable challenges until now. The emergence of 316 TADF emitters have provided a viable alternation for this goal. Finding strategies that can 317 318 improve stability of TADF emitters have been an exigent task and highly desired in the industry and academia. The deuteration strategy proposed here can simultaneously enhance efficiency and 319 operational stability of blue TADF emitters. From the perspective of both theoretical calculation 320 321 and experiments, it is confirmed that deuteration can significantly slow down high-frequency vibrations and create a relatively shallow PES in both ground and excited states, energetically 322 favoring to suppress dissociative process and eliminating nonradiative decay. The 323 324 proof-of-the-concept deuterated blue TADF emitter d-5CzBN realized nearly 3-times longer device lifetime with ~ 1.3 times higher EQE_{max} than the undeuterated one. And the TSF device 325

326 with d-5CzBN as sensitizer while a narrowband dopant as final emitter, a deep-blue device realized a superb long LT80 of 225 h and extremely high EQE_{max} of 30.3% simultaneously at a 327 CIE_{ν} of 0.17. By further optimizing the final emitter, blue device with CIE_{ν} of 0.19 can even 328 realize a LT80 of 456 h at an initial luminance of 1,000 cd/m². Quite recently, Bae et al. revealed 329 that deuterating the vulnerable benzylic C-H bond can greatly increase the operational lifetime of 330 phosphors by hampering the possible C-H bond cleavage.³¹ And Adachi's group also studied the 331 isotope effect of the host material and found a positive effect on the device stability,²¹ which was 332 assigned to the enhanced stable amorphous nature and balanced carrier transport properties. The 333 mechanism disclosed here also works well for previous studies. Our work may pave the way for 334 developing better deep-blue TADF emitters and devices that can rival conventional fluorescent 335 devices. 336

337 Methods

General Information. All commercially available reagents were used as received unless 338 otherwise stated. 5CzBN and d-5CzBN were synthesized according to Supplementary 339 Information. Mass spectra were recorded on a Shimadzu MALDI-TOF mass spectrometer. The 340 341 electrochemical measurements were performed with a CHI600E electrochemical workstation by using Pt as the working electrode, platinum wire as the auxiliary electrode, and an Ag wire as the 342 reference electrode standardized against ferrocene/ferrocenium. Organic films for optical 343 344 measurements were fabricated by thermal evaporation under high vacuum onto clean quartz substrates. UV-vis absorption spectra were recorded by an Agilent 8453 spectrophotometer. 345

Fluorescence and phosphorescence spectra at steady state were recorded by HITACHI F-7000
Fluorescence Spectrometer. Fluorescence lifetime measurement was carried out with Edinburgh
fluorescence spectrometer (FLS1000) using picosecond pulsed diode laser under the excitation at
365 nm. PLQYs were measured by Hamamatsu absolute PL quantum yield spectrometer
(C9920-02G) with an integrating sphere.

Theoretical calculation. All theoretical calculations were performed using a density functional theory (DFT) method as implemented in Gaussian 16 and ORCA in gas state.^{32,33} Ground-state geometries were optimized by B3LYP functional with the 6-31G(d) basis set. Frequency analysis was performed under the same level with corresponding frequency scale factors.³⁴ Single point energy was calculated by double hybrid functional B2PLYP-D3 with the def2-TZVP basis set.

Device fabrication and Characterization. Before device fabrication, the ITO glass substrates 356 were precleared carefully. Then the sample was transferred to the deposition system. The devices 357 were prepared in vacuum at a pressure of 5×10^{-6} Torr. The organic layers were thermally 358 evaporated at a rate of 1.0 Å s⁻¹. After the organic film deposition, 0.5 nm of LiF and 150 nm of 359 aluminum were thermally evaporated onto the organic surface. All the organic materials used 360 361 were purified by a vacuum sublimation approach. The electrical characteristics of the devices were measured with a Keithley 2400 source meter. The electroluminescence spectra and 362 luminance of the devices were obtained on Hamamatsu external quantum efficiency 363 364 measurement system. All the device fabrication and characterization steps were carried out at 365 room temperature under ambient laboratory conditions.

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376 Author contributions

Tianyu Huang performed the experiments and prepared the manuscript. Ge Zhan synthesized
deuterated materials. Dongdong Zhang and Lian Duan conceived and supervised the research.
All authors participated in the discussion on experimental results and revised the manuscript.

380 Competing interests

381 The authors declare no competing interests.

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Reference

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459 Data availability

- 460 The data that support the plots within this paper and other findings of this study are available
- 461 from the corresponding author upon reasonable request.

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