Extrapolating DFT towards the complete basis set limit: Lessons from the PBE family of functionals

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Abstract

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Extrapolation of density functional theory results from 2- and 3- ζ calculations is a 2 promising method for extracting higher accuracy data from calculations of systems at the 3 affordability limit. In this work, I present formulas for the determination of extrapolation 4 parameters, that account for the make-up of the density functional approximation. The 5 formulas are fitted to reproduce the complete basis set limit energies of PBE and related 6 density functional approximations, using a set of 30 singlet diatomics. Their performance 7 is extensively evaluated using standard benchmark datasets. The current formulas are 8 shown to be transferable outside the PBE-family of functional approximations, with the 9 resulting extrapolation parameters outperforming the previous, empirical values. A good 10 performance of [2,3]- ζ extrapolations for interaction energies of systems with significant 11 non-covalent character is confirmed, and holds even in systems of ${\sim}100$ atoms in size. 12

13 **1** Introduction

¹⁴ Basis set extrapolation methods allow computational chemists to reliably approximate results of ¹⁵ expensive calculations in larger basis sets by combining results obtained at the same level of ¹⁶ theory from two or more cheaper calculations in smaller basis sets. This is particularly important ¹⁷ in post-Hartree-Fock (HF) wavefunction theory, where the calculation of electronic correlation ¹⁸ dominates the computational cost. It has long been known that the correlation energy (E^{corr}) ¹⁹ approaches its infinite basis set limit (E_{∞}^{corr}) as a function of a cubic power law of the basis set ²⁰ size,

$$E_{\infty}^{\rm corr} = E_X^{\rm corr} - AX^{-\alpha} \tag{1}$$

where the extrapolation parameter $\alpha = 3$, X is the cardinal number of a finite correlationconsistent basis set, and A is a system-dependent parameter fitted to the E_X^{corr} data.¹ Notably, while this cubic power law is correct for all values of Xs, other empirically determined values of α may provide better results for $X \leq 3$.²

Extrapolation methods are not widely used in density functional theory (DFT) calculations. 25 This may be due to several reasons: DFT calculations are computationally inexpensive compared 26 to correlated WFT, reducing the applicability of such methods; there are many different empirical 27 density functional approximations (DFAs) with each requiring its own α , meaning such approaches 28 are seen as less general; reference datasets of complete basis set limit energies from DFT are not 29 as common as their WFT counterparts, hindering the development of such methods; and even 30 the true basis set limit convergence of HF was until recently a matter of debate, putting any DFT 31 extrapolation attempts on an uncertain theoretical basis. 32

In a previous work,³ I have shown that the extrapolation parameters α determined for a set of DFAs using numerical complete basis set calculations are surprisingly transferrable to other DFAs. However, the values of α are strongly dependent on the basis set family. For instance, extrapolated def2-[st]zvpd non-covalent interaction energies generally outperform their
 def2-qzvpd counterparts at 80% of the computational cost, which holds true even for double hybrid density functional approximations (DHDFAs) as long as the correlated perturbation theory
 (PT) component is extrapolated appropriately. The best performing extrapolation function for
 DFT is the exponential–square root function (expsqrt), first proposed by Jensen:⁴

$$E_{\infty}^{\text{fctl}} = E_X^{\text{fctl}} - Ae^{-\alpha\sqrt{X}}$$
⁽²⁾

This function has since been proven to be the correct form for energy convergence in both HF 41 and Kohn-Sham DFT with Gaussian basis sets.⁵ However, several issues higlighted in the previous 42 work remain. The determination of the values of the extrapolation parameters α , as described in 43 Ref. 3, is highly method-specific due to the introduced averaging. Additionally, the values of α 44 clearly scale with the inclusion of exact (HF) exchange in hybrid DFAs, but the introduction of 45 formulas which correct the α by the percentage of exact exchange did not improve the results. 46 Finally, the influence of the reduced correlation component in double hybrid exchange-correlation 47 functionals upon introduction of PT correlation has not been investigated. 48

In the current work, I address the three above issues by systematically investigating of PBE-like 49 DFAs. Compared to the set of 9 diatomics used in previous work,³ a larger dataset of complete 50 basis set energies from finite element calculations is used, comprising 30 singlet diatomic species, 51 and including cations and anions. This dataset is evaluated for PBE and related functionals with 52 varying fractions of HF exchange and PT correlation. I investigate trends over a large number of 53 basis set families, including most of the modern basis set families used with DFT and WFT. This 54 leads to a formula which is specific for each basis set family, but universal for all DFAs, taking 55 into account all components of the DFA recipe. The extrapolation parameters α obtained from 56 these formulas are then thoroughly examined using several benchmark databases. 57

2 Computational methods

Analogously to previous work, I use the following nomenclature for the total energy of a DFA (E^{DFA}) and its components:

$$E^{\text{DFA}} = E^{\text{fctl}} + \Delta E^{\text{dh}} + \Delta E^{\text{disp}} + \Delta E^{\text{nl}}$$
(3)

The first component is the self-consistent-field (SCF) energy of the exchange-correlation functional (fctl), the second component is the double-hybrid (dh) correction, the third term corresponds to dispersion corrections, such as Grimme's -D3,⁶ and the fourth term is a non-local correlation term, such as in the VV10 DFA,⁷ or generalised as -NL correction.⁸

⁶⁵ The first two terms can be further split into the following components:

$$E^{\text{fctl}} = (1 - a_x)E^x + a_x E^{\text{HFx}} + (1 - a_c)E^c$$
(4)

$$\Delta E^{\rm dh} = a_{\rm c} E^{\rm PTc} \tag{5}$$

Note that the coefficients a_x and a_c correspond to the fraction of HF exchange (HFx) and PT 66 correlation (PTc) in the DFA recipe. While a_x is constrained to ≤ 1 ,⁹ no such rigorous constraint 67 on ac exists.¹⁰ In practice, most common DHDFAs obey both constraints, with the notable 68 exception of the spin-component-scaled double hybrids (DSD family).¹¹ For simplicity, in the 69 following discussion it is assumed both constraints hold. The fit is always performed using the 70 exact values of the coefficients in front of E^{x} and E^{c} in Eq. (4), even if the latter is not equal 71 to $(1 - a_c)$. Finally, while in principle the ΔE^{nl} term contributes to the SCF cycle, it has been 72 conclusively shown to produce nearly identical results when applied as a post-SCF correction.¹² 73 The post-SCF treatment is applied throughout this work. 74

The $\{a_x, a_c\}$ parameter space explored in this work consists of all 20 combinations of $a_x \in$

⁷⁶ {0.1, 0.3, 0.5, 0.7, 0.9} with $a_c \in \{0.0, 0.2, 0.4, 0.6\}$ as well as the pure variant (i.e. $a_x = a_c =$ ⁷⁷ 0.0). The parameter sets are chosen to cover the range used by most of the common single ⁷⁸ and double hybrids. The exchange and correlation functionals of Perdew, Burke, and Ernzerhof ⁷⁹ (PBE)¹³ are used, as they represent a "non-empirical" DFA, even though it could be argued ⁸⁰ parameter-counting is a questionable metric of empiricity.¹⁴ The generalization of the results ⁸¹ obtained from this family of PBE-like functionals is discussed below.

The complete basis set energies of the 30 diatomic molecules listed in Table 1 are calculated 82 using the diatomic component of the finite element code HelFEM, ¹⁵ using parameters deter-83 mined with the diatomic_cbasis tool. Three anionic diatomics were difficult to converge using 84 the following combinations of a_x and a_c : SF⁻ did not converge with $a_x = 0.1$ and $a_c = 0.6$, and 85 OF⁻ as well as OH⁻ failed to converge with the pure functional, as well as for all four values 86 of $a_{\rm c}$ with $a_{\rm x} = 0.1$. This is attributed to the diffuse nature of the anion exacerbated by the 87 self-interaction error in PBE. Neither increasing the numerical integration radius from 40 Å to 88 120 Å, nor increasing the grid parameters led to convergence. 89

All other calculations were performed with Psi4 version 1.4a2.dev923,¹⁶ or a development 90 version of the program including basis set extrapolation routines for DFT implemented as part of 91 previous³ and current work. Calculations of the 30 diatomic molecules were carried out using the 92 PK-supermatrix SCF algorithm to avoid issues with missing auxilliary basis sets, with a tightened 93 energy convergence (10^{-10} Eh) , and a large DFT quadrature with 150 radial and 974 angular 94 points. The basis set families investigated in this work are listed in Table 2, significantly expanding 95 on previous work where only four basis set families were compared. All basis sets were used as-96 included with Psi4, or downloaded from the Basis Set Exchange.^{17,18} It is notable that the energy 97 obtained with 4- ζ variants of cc-pwcvXz and aug-cc-pwcvXz is in many cases lower than that 98 with the 5- ζ variants; this non-convergence does not happen with other basis sets. 99

Results of the calculations of the diet100 variant³⁷ of the GMTKN55 database³⁸ are reused from previous work, comprising 4 basis set families and 15 DFAs (see Table S1 for a detailed list); the WTMAD-2 criterion is used as figure of merit.³ Additional calculations of

Table 1: Set of 30 diatomic molecules, including their bond lengths (R), the ℓ_{max} of the basis sets, and the number of radial elements (n_{elem}) used in the finite element calculations. The numerical integration radius was 40 Å in all cases, the number of nodes in each element was 15.

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Diatomic	R (Å)	$\ell_{\sf max}$		<i>n</i> elem
CH ⁺	1.13085	15	11	3
OH^+	1.02890	17	13	3
OH-	0.94246	15	13	3
FH	0.91696	17	13	5
C ₂	1.24780	17	11	3
CN^+	1.17290	17	13	3
CN^{-}	1.17160	17	13	3
N_2	1.09434	17	11	3
NO^+	1.06206	17	13	3
CO	1.12821	17	13	5
CF^+	1.22875	19	15	5
OF ⁻	1.49228	21	17	5
F_2	1.41184	21	15	5
SiH^+	1.50410	25	19	5
SH^-	1.34993	27	21	5
HCI	1.29119	27	21	5
CP^{-}	1.58753	27	21	5
CS	1.53442	27	21	7
SiN ⁻	1.55578	27	21	5
NP	1.49085	27	21	5
SN^+	1.44000	27	21	5
SiO	1.50974	25	19	5
PO^+	1.41900	25	19	5
SF^-	1.70395	29	23	7
CIF	1.66162	29	23	7
SiS	1.93000	31	23	7
P_2	1.89340	31	23	7
PS^+	1.87200	31	23	7
SCI ⁻	2.14846	33	25	7
Cl_2	2.04262	33	25	7

Family	ζ range	Core?	Diffuse?	Reference
cc-pvXz [†]	2 – 6	×	×	Dunning ^{19–22}
cc-pwcv X z †	2 – 5	\checkmark	×	Dunning ^{19–21,23}
aug-cc-pv X z †	2 – 6	×	✓	Dunning ^{19–22}
aug-cc-pwcv X z †	2 – 5	\checkmark	✓	Dunning ^{19–21,23}
def2- X zvp †	2 – 4	×	×	Ahlrichs ²⁴
def2-Xzvpp	2 – 4	×	×	Ahlrichs ²⁴
def2- X zvpd †	2 – 4	×	✓	Ahlrichs ^{24,25}
def2-Xzvppd	2 – 4	×	✓	Ahlrichs ^{24,25}
pc-N	1 – 5*	×	×	Jensen ²⁶
pcseg- N^{\dagger}	1 – 5*	×	×	Jensen ²⁷
aug-pc-N	1 – 5*	×	✓	Jensen ²⁸
aug-pcseg- N^{\dagger}	1 – 5*	×	✓	Jensen ²⁷
Xzapa-nr	2 – 6	×	✓	Ranasinghe ²⁹
Xzapa-nr-cv	2 – 6	\checkmark	✓	Ranasinghe ^{29,30}
jorge-Xzp	2 - 6	×	×	Jorge ³¹⁻³⁴
jorge-aXzp	2 – 5	×		Jorge ^{31,35,36}

Table 2: Basis set families, the ζ range used in extrapolation, and whether the basis set includes tighter core *d*-functions (Core) and/or diffuse augmentation (Diffuse).

* Note that for (aug-)pc-N and (aug-)pcseg-N families, $N \approx X + 1$. † Basis set families used in the ASCDB benchmark.

the ASCDB database³⁹ were carried out using the standard density-fitted SCF algorithm, and 103 the same (150, 974) grid, comprising 8 basis set families (marked with † in Table 2) and 3 PBE-104 like DFAs (PBE-D3(BJ),^{13,40} PBE0-D3(BJ),^{40,41} and PBE0DH-D3(BJ)^{42,43}). Cases including 105 elements for which basis sets are not defined within a given basis set family are excluded from 106 the dataset of the corresponding basis set family. The ASCDB database allows for decomposition 107 of the overall error metric into components of non-covalent interactions (NCI), thermochemistry 108 (THCH), non-local effects (NLE), and unbiased calculations (UNBC).³⁹ allowing a more detailed 109 insight into the performance of an extrapolation method than with the diet GMTKN55 subsets.³⁷ 110 The performance of basis set extrapolations for optimizing geometries of van der Waals 111 complexes was evaluated using the NCDT database.⁴⁴ in its revised form.⁴⁵ using a develop-112 ment version of Psi4 with a tightened set of convergence criteria ($\Delta E < 1 \mu Eh$, max F < 1113 15 $\mu \text{Eh}/a_0$, RMS $F < 10 \ \mu \text{Eh}/a_0$, max $d < 0.6 \ ma_0$, RMS $d < 0.4 \ ma_0$, i.e. interfrag_tight) 114 and an unpruned (150, 974) quadrature grid. The geometries were optimized using the def2-115

¹¹⁶ Xzvpd basis set family, used with revPBE-D3(BJ), M062X-D3, B97M-D3(BJ), B3LYP-D3(BJ), ¹¹⁷ PBE0-D3, and ω B97X-D3(BJ), selected based on their performance in GMTKN55, ⁴⁶ and previ-¹¹⁸ ous NCDT results.⁴⁷

To evaluate the performance of the basis set extrapolations in their most likely use case, that is 119 in large systems with a significant non-covalent character, the interaction energies of the geome-120 tries from the L7 dataset⁴⁸ were evaluated using M062X-D3/def2-Xzvpd and B97M-V/pcseg-N. 121 This benchmark was designed to capture the dominant binding motifs in biological chemistry, 48 122 comprising 7 systems ranging in size between 48 and 112 atoms. The revised reference energies 123 using correlation corrections obtained from DLPNO-CCSD (T_0) /cc-pv[dt]z calculations⁴⁹ are used 124 instead of the original values.⁴⁸ This revised set is in a good agreement with other recent reference 125 data,⁵⁰ however even with significantly different reference values⁵¹ the trends reported here are 126 not affected. The calculations for the L7 dataset are performed with Psi4 version 1.4a2.dev923, 127 using the default unpruned (75, 302) grid and the frozen core approximation for all double hybrid 128 functionals. 129

The values of A and α in Eq. (2) for each of the 30 diatomics are obtained using a non-linear last squares fit of Eq. (6) using the curve_fit function of the Python library scipy.optimize.⁵²

$$\ln(E_X^{\text{fctl}} - E_\infty^{\text{fctl}}) = \ln(A) - \alpha \sqrt{X}$$
(6)

The E_X^{fctl} values correspond to the energies from individual basis sets in each family listed in Table 2, with $X \approx \zeta$. The complete basis set limit value E_{∞}^{fctl} is the value obtained from HeIFEM. The extrapolation parameters α are then averaged based on the system charge, obtaining the means ($\overline{\alpha}$) and medians ($\tilde{\alpha}$) for the cationic (α^+), neutral (α^0), anionic (α^-), and all (α) diatomic species. Finally, a linear fit to the mean values $\overline{\alpha}$ of all combinations of { a_x, a_c } is performed, using Eq. (7), obtaining the intercept α_0 , and two linear coefficients α_x and α_c .

$$\overline{\alpha} = \alpha_0 + \alpha_x a_x + \alpha_c a_c \tag{7}$$

3 Results

3.1 Extrapolation formulas

The results for each basis set family at all a_x with $a_c = 0$ are presented in Fig. 1. The cationic 140 outlier present in all basis sets is OH^+ ; the anionic outlier most visible in def2-Xzvp is OH^- . With 141 the exception of the Jensen basis set families, all $\overline{\alpha}$ s show a positive correlation with a_x . The 142 inclusion of HF exchange into the PBE recipe actually increases the overall self-interaction error 143 (SIE) for hydrogenic atoms, as for the pure functional the exchange and correlation contributions 144 to the SIE cancel each other out, and at non-zero a_x this error cancellation is lost.⁵³ However, 145 for diatomics, the inclusion of HF exchange considerably improves the performance of the DFA 146 in the SIE4x4 benchmark.⁵³ In any case, with a fraction of HF exchange, the electronic density 147 is less delocalised. As the exp $(-\alpha\sqrt{X})$ term accounts for the error in the convergence of energy 148 with respect to basis set size X,⁵⁴ it is not surprising that for more compact densities the errors 149 vanish quicker. This may also explain why no such correlation holds for a_c . 150

The augmentation of the basis sets by diffuse functions leads to a clear shift of the extrapolation parameters α for anionic species (•) to values in line with the neutral and cationic species. This trend has been discussed in the context of basis set extrapolations previously.^{3,55} Here, I wish to reiterate the need for the use of diffuse functions whenever anionic systems are involved. The data show no statistically significant difference between $\overline{\alpha^+}$ and $\overline{\alpha^0}$.

The addition of tighter core *d*-electron functions pushes the extrapolation parameters α for many of the diatomics upwards towards ~ 6 . This trend is most apparent for aug-cc-pvXz \rightarrow aug-cc-pwcvXz augmentation, however it is also present in cc-pvXz \rightarrow cc-pwcvXz, as well as for Xzapa-nr \rightarrow Xzapa-nr-cv. It should be noted that the value of α that is optimal for predicting



Figure 1: Extrapolation parameters α obtained with various basis set families at $a_c = 0$ and a_x as shown on the vertical axis. Colour coding corresponds to cationic (•), neutral (•), and anionic (•) species, with the box plots showing the overall median $(\tilde{\alpha}, |)$, mean $(\overline{\alpha}, |)$, quartiles, and the 5th and 95th percentile.

- $_{160}$ the error of the energy expectation value approaches $\pi\sqrt{3}pprox$ 5.44 from below as basis set size
- ¹⁶¹ approaches completeness.⁵⁴

Family	α_0	α_{x}	α_{c}
cc-pvXz	3.622	1.511	0.005
cc-pwcv X z	4.157	1.192	-0.048
aug-cc-pv X z	3.676	1.887	0.139
aug-cc-pwcv Xz	4.485	1.445	0.085
рс- <i>N</i>	6.172	-1.623	0.183
pcseg-N	5.883	-1.825	0.227
aug-pc-N	6.390	-1.874	0.260
aug-pcseg-N	6.166	-2.137	0.296
def2-Xzvp	7.406	1.266	-0.046
def2-Xzvpp	7.408	1.267	-0.046
def2-Xzvpd	7.925	1.370	0.101
def2-Xzvppd	7.927	1.371	0.101
jorge-Xzp	3.531	0.338	-0.011
jorge-a X zp	3.386	0.245	0.033
Xzapa-nr	3.306	2.525	-0.040
Xzapa-nr-cv	5.618	0.490	-0.016

Table 3: Exchange and correlation dependence of $\overline{\alpha}$ for each basis set family studied. Coefficients are as per Eq. (7).

¹⁶² 3.2 ASCDB database

The mean unsigned errors (MUE) in the ASCDB database are shown for three widely used 163 PBE-like functionals in Fig. 2. As expected, the pure functional PBE-D3(BJ) is outperformed by 164 the single hybrid PBE0-D3(BJ). However, the double hybrid PBE0DH-D3(BJ) does not provide a 165 systematic improvement over the single hybrid. Diffusely augmented basis set families outperform 166 the corresponding non-augmented families in all cases, which is not surprising given the inclusion 167 of non-covalent interactions in the ASCDB database. The Jensen basis set families (pcseg-N168 and aug-pcseg-N) systematically outperform both the original Dunning (cc-pvXz and aug-cc-169 pvXz) as well as the Ahlrichs (def2-Xzvp and def2-Xzvpd) sets. The weighted core-valence 170 correlation consistent Dunning sets including tight d-functions (cc-pwcvXz and aug-cc-pwcvXz) 171 offer comparable performance to the Jensen sets, albeit at a higher computational cost. 172



Figure 2: The mean unsigned errors of various functional and basis set combinations in the ASCDB database. Calculations with 2-, 3-, and 4- ζ basis sets shown as bars. Results from the [2,3]- ζ extrapolation (•) and [3,4]- ζ extrapolation (×) from current work compared to previous [2,3]- ζ results (•), where available.

The performance of the [2,3]- ζ extrapolation (•) using the formulas from Table 3 is encourag-173 ing. The extrapolations generally outperform the 3- ζ calculations, with the following exceptions: 174 The extrapolation using aug-cc-pvXz basis sets with PBE-D3(BJ) performs worse than the 2- ζ 175 calculations; and the extrapolations using def2-Xzvp, def2-Xzvpd, cc-pvXz, or aug-cc-pvXz basis 176 sets with PBE0DH-D3(BJ) perform worse than 3- ζ results. These exceptions will be analysed 177 further below. Notably, the previously determined [2,3]- ζ extrapolation parameters (•)³ perform 178 worse than the current formulas in all cases. The results are especially encouraging as the revised 179 formulas perform significantly better with the Jensen basis sets. Finally, the [3,4]- ζ extrapolation 180 (\times) rarely outperforms the 4- ζ results, and as such cannot be recommended. 181

Detailed results for the various subsets of the ASCDB database are included in the Supple-182 mentary Information (Figs. S1–S4). In the non-covalent interactions subset, $[2,3]-\zeta$ extrapolation 183 performs significantly better than the 3- ζ calculations using all basis set families and functionals. 184 Generally results comparable to 4- ζ quality can be expected. In the thermochemistry subset, the 185 performance of the [2,3]- ζ extrapolation is more mixed. In the most extreme cases, the MUE 186 increases with basis set size (e.g. PBE-D3(BJ) with aug-cc-pvXz or aug-cc-pwcvXz); the re-187 sulting poorer performance of extrapolated data can be attributed to the functional as opposed 188 to the basis sets. That said, when 3- ζ and 4- ζ results are compared, the MUE in the thermo-189 chemistry data does not improve as significantly as for non-covalent interactions, and the [2,3]- ζ 190 extrapolation is unlikely to increase the errors. In systems with significant non-local character, 191 extrapolation methods offer no benefit; in fact it is this class of problems which causes the poor 192 performance of PBE0DH-D3(BJ) when coupled with a [2,3]- ζ extrapolation using def2-Xzvp, 193 def2-Xzvpd, aug-cc-pvXz, or aug-cc-pvXz basis sets. This is of course a more general issue with 194 DFT as opposed to a basis set completeness error. Notably, the non-local effects subset is the 195 only subset where the double-hybrid PBE0DH-D3(BJ) consistently outperforms the single hybrid 196 PBE0-D3(BJ) when used with 4- ζ basis sets; the [2,3]- ζ extrapolation applied to cc-pwcvXz and 197 aug-cc-pwcvXz yields results comparable to the 4- ζ data. Finally, in the unbiased benchmarks 198 subset of ASCDB, the performance of the $[2,3]-\zeta$ extrapolation is also underwhelming. This 199 might be correlated with the significant increase in performance of the 3- ζ basis sets compared to 200 the 2- ζ results, and the comparably minor improvement upon a further increase in the basis set 201 size to 4- ζ . The 2- ζ results are likely too inaccurate to improve the 3- ζ results by extrapolation. 202

203 3.3 GMTKN55 database

The results of basis set extrapolation with four representative DFAs for the diet100 subset of the GMTKN55 database is shown in Fig. 3. More detailed results, calculated with a wider range of functionals and the cc-pvXz-pp, def2-Xzvp, def2-Xzvpd, and pcseg-N basis set families, are shown in the Supplementary Information (Figs. S5). In the vast majority of cases, the current



Figure 3: The weighted total mean absolute deviation (WTMAD-2) of various functional and basis set combinations in the diet100 subset of the GMTKN55 database. Calculations with 2-, 3-, and 4- ζ basis sets shown as bars. Results from the [2,3]- ζ extrapolation (•) and [3,4]- ζ extrapolation (×) from current work compared to previous [2,3]- ζ results (•), where available.

[2,3]- ζ extrapolation (•) outperforms or matches the previous extrapolation results (•), with the 208 only exception being the cc-pvXz-pp results (see e.g. M052X-D3 and B2PLYP-D3(BJ) data). 209 The four DFAs shown in Fig. 3 are all significantly different from the PBE family of functionals 210 (the closest being revPBE, where the Lieb-Oxford bound is relaxed ⁵⁶), therefore this improvement 211 in performance confirms the transferrability of the extrapolation formulas, which are based on 212 PBE-like DFAs. Additionally, in all but two cases (DSD-BLYP-D3(BJ) with def2-Xzvp and pcseg-213 N basis sets), the current [2,3]- ζ extrapolation performs at least as well as 3- ζ calculations. As 214 also shown in previous work,³ the extrapolation using def2-[s,t]zvpd basis sets matches def2-215 gzvpd performance in this GMTKN55 subset. This is encouraging, as diffusely-augmented basis 216 sets are necessary for proper description of anionic systems, and the def2-Xzvpd sets are defined 217 for the whole periodic table. The performance of the [2,3]- ζ extrapolation with pcseg-N basis 218

²¹⁹ sets is generally also improved, especially with DHDFAs. However, as also noted previously, ³ the ²²⁰ pcseg-2 results (3- ζ quality) generally outperform the [2,3]- ζ extrapolation with def2-Xzvp and ²²¹ cc-pvXz-pp and are often comparable to the cc-pvqz-pp results. The choice of an appropriate ²²² basis set family therefore remains crucial and cannot be remedied by basis set extrapolation.

223 **3.4 NCDT database**

Given the good performance of extrapolation schemes for interaction energies of systems with significant non-covalent character, a similarly good performance may be expected for obtaining geometries of non-covalent complexes with such schemes. The results for the NCDT database, presented as the mean absolute deviations of selected bonds in each complex, summed over all 17 complexes, are shown in Fig. 4 for six density functional approximations, three of which have been recommended based on the GMTKN55 benchmarks, ⁴⁶ and two (B3LYP-D3(BJ) and PBE0-D3) have performed well in previous NCDT benchmarks.⁴⁷



Figure 4: The total of the mean absolute deviations in the optimized structures of the 17 complexes in the NCDT database. The performance of the def2-[st]zvpd extrapolation is indicated by the gray bar.

To bring perspective to the extrapolation results (indicated by an the gray bar in Fig. 4), let 231 us start with a few comments on the general performance of the DFAs and trends in the dataset. 232 Firstly, with the exception of M062X-D3, all DFAs struggle with optimization of the interfragment 233 angle of the HCN \cdots H₂CO structure; in the case of B3LYP-D3(BJ) this complex alone accounts 234 for $\sim 25\%$ of the total MAD. Secondly, the three functionals chosen based on their GMTKN55 235 performance perform significantly worse than both B3LYP-D3(BJ) and PBE0-D3, which remain 236 good methods for structure optimization. However, the comparison is not quite fair, as the 237 range-separated hybrids ω B97M-V and ω B97X-V, which were previously found to offer excellent 238 performance with both GMTKN55 and NCDT, are here excluded from the study as analytical 239 gradients for the non-local correction are not available in Psi4. Accordingly, the ω B97X-D3(BJ) 240 variant, for which analytic gradients are available, is also evaluated and with a 4- ζ basis it shows 241 results on par with B3LYP-D3(BJ). Finally, in most cases, the improvement from def2-svpd to 242 def2-tzvpd is significant, while the difference between def2-tzvpd and def2-gzvpd is marginal, 243 if any. With revPBE-D3(BJ), increasing the basis set size decreases the agreement with the 244 reference data, which is likely an issue with the DFA. The M062X-D3 functional was trained 245 using a 3- ζ basis set, therefore a decrease in performance with a 4- ζ basis is not unexpected. 246 The comparably poor performance of the basis set extrapolation can then be attributed to the 247 small potential for improvement in geometries past 3- ζ basis sets, and the fact that more ac-248 curate interaction energies do not necessarily translate to more accurate geometries. Basis set 249 extrapolations therefore cannot be recommended to obtain more accurate structures with DFT. 250

251 **3.5 L7 database**

The most likely use-case of extrapolation methods in DFT would be for systems where a calculation of 3- ζ quality would approach the limit of affordability. The systems in the L7 database are comparably large, between 48 and 112 atoms in size,⁴⁸ yet reference interaction energies of near-CCSD(T) quality are now available.⁴⁹ This dataset is therefore an appropriate proxy for the performance of methods in such applications.

The results with a variety of combinations of DFAs and basis set families are shown in Fig. 5. 257 The three double hybrids (DSD-BLYP-D3(BJ), DSD-PBEP86-NL, and B2PLYP-D3(BJ)) per-258 form poorly compared to the less expensive methods, with a 3- ζ MAE above 20 kJ/mol and a 259 systematic overbinding regardless of the basis set family. This may be due to the exceptionally 260 slow convergence of the double hybrid correction ΔE^{dh} , especially in the two spin-component-261 scaled double hybrids. Indeed, the reported MAE of MP2 extrapolated towards the complete 262 basis set limit exceeds 30 kJ/mol.⁴⁹ On the contrary, a comparably good performance of double 263 hybrid DFAs with -NL correction for the L7 database has been reported by Calbo et al., who 264 compared B2PLYP-NL and revPBE0-DH-NL to their single hybrid and non-hybrid counterparts, 265 albeit with different reference energies.⁵¹ The use of the reference energies of Calbo et al.⁵¹ or 266 Al-Hamdani et al.,⁵⁰ as opposed to the data of Ballesteros et al.,⁴⁹ has no impact on the trends 267 shown in Fig. 5. A more detailed analysis of the poor performance of double hybrid DFAs is 268 beyond the scope of this work. 269



Figure 5: The mean absolute deviations in the interaction energies of the 7 structures in the L7 database. The performance of the [2,3]- ζ extrapolation methods is indicated by the gray bar.

As also shown in Fig. 5, the [2,3]- ζ extrapolations outperform the 3- ζ results in all cases, with significant improvements over 3- ζ results using B97M-V/pcseg-[12], B3LYP-D3(BJ)/def2[st]zvpd, as well as the three double hybrids. Notably, none of the six methods presented in
 Fig. 5 are members of the PBE-family of functionals, confirming the good transferability of the
 extrapolation formulas developed in the current work.

275 4 Conclusion

A systematic way of extrapolating density functional theory results towards the complete basis set limit can be derived from data computed with the PBE family of density functional approximations. The proposed formulas adjust the extrapolation α for each DFA based on the admixture of Hartree-Fock exchange and second-order perturbation theory correlation into the functional recipe. The extrapolation parameters α are also dependent on the basis set family used in the extrapolation. The use of diffusely-augmented basis sets is strongly recommended for anionic species.

A set of complete basis set energies for 30 singlet diatomic molecules has been calculated 283 with 21 PBE-related DFAs using the finite element code HelFEM. The dataset is complemented 284 by energies of the same 30 diatomic molecules, calculated with 16 basis set families and the 285 same 21 PBE-related DFAs. This comprehensive set was used to fit linear formulas for scaling 286 the extrapolation parameter α based on the fraction of HF exchange a_x and PT correlation a_c in 287 each of the 21 DFAs. The formulas for calculating the parameter α are determined individually 288 for each basis set family, while the extrapolation parameters α derived from such formulas are 289 DFA as well as basis set family specific. 290

The current, systematically derived, extrapolation parameters outperform the previous, empirically averaged values. For single point energies in the ASCDB database, the performance of $[2,3]-\zeta$ extrapolations exceeds 3- ζ and generally approaches 4- ζ results in systems with significant non-covalent character. For thermochemistry calculations the $[2,3]-\zeta$ results at worst match and often outperform 3- ζ results. The improvement over 3- ζ results in interaction energies for the L7 database of large systems is also significant. On the contrary, for geometry optimizations of

²⁹⁷ non-covalent complexes, the [2,3]- ζ extrapolation barely outperforms 3- ζ results and therefore ²⁹⁸ cannot be recommended. The basis set extrapolation methods are implemented in a development ²⁹⁹ version of Psi4, with extrapolated analytic gradient calculations as well as extrapolated numerical ³⁰⁰ Hessians available.

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Supporting Information Available

Supporting information available: List of methods used with the diet100 variant of GMTKN55; Figures of results of the subsets of the ASCDB database; Figures of results of the diet100 subset of the GMTKN55 database. Additional supporting information, including the scripts generating the inputs, parsing the outputs, and creating figures are available on Zenodo under DOI: 10.5281/zenodo.4783007.

³¹⁴ Table of contents graphic

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