Breaking isolation to form new networks: pH-triggered changes in connectivity inside lipid nanoparticles

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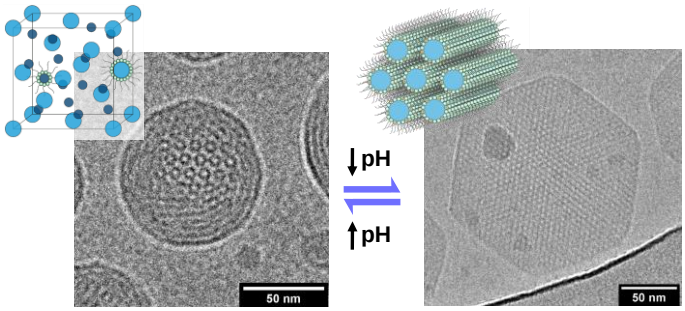
ABSTRACT

There is a growing demand to develop smart nanomaterials that are structure-responsive as they have the potential to offer enhanced dose, temporal and spatial control of compounds and chemical processes. The naturally occurring pH gradients found throughout the body make pH an attractive stimulus for guiding the response of a nanocarrier to specific locations or (sub)cellular compartments in the body. Here we have engineered highly sensitive lyotropic liquid crystalline nanoparticles that reversibly respond to changes in pH by altering the connectivity within their structure at physiological temperatures. At pH 7.4, the nanoparticles have an internal structure consisting of discontinuous inverse micellar 'aqueous pockets' based on space group Fd3m. When the pH is ≤6, the nanoparticles change from a compartmentalized to an accessible porous internal structure based on a 2D inverse hexagonal phase (plane group p6mm). We validate the internal symmetry of the nanoparticles using Small Angle X-ray Scattering and cryogenic Transmission Electron Microscopy. The high resolution electron microscopy images obtained have allowed us for the first time to directly visualize the internal structure of the Fd3m nanoparticles and resolve the two different-sized inverse micelles that make up the structural motif within the Fd3m unit cell, which upon structural analysis reveal excellent agreement with theoretical geometrical models.

TOC



INTRODUCTION

Lyotropic liquid crystalline nanoparticles (LCNPs) have attracted enormous interest for applications such as drug delivery, nanoreactor arrays, biosensors, templating, advanced biomaterials and functional foods. These particles have non-lamellar internal nanostructrures that possess two- or three-dimensional periodicity. LCNPs with internal nanostructures based on inverse bicontinuous cubics (Q_{II}), 2D-hexagonal (H_{II}) or ordered micellar phases are termed cubosomes, hexosomes and micellosomes respectively. They offer potential advantages over current nanoformulations, including favourable payloads due to their high internal surface area, simple preparation protocols, superior ease of conjugation with target biomolecules, biodegradability of the host building blocks and the ability to encapsulate hydrophobic, hydrophilic and amphiphilic substances.[1-6]

Currently there is a large demand to develop smart nanocarriers that can respond to particular environmental stimuli which could be used to release an active at a specific site of action and hence minimize off target toxicity and unwanted side effects. Extrinsic factors such as ultrasound, temperature, magnetic fields and light have been explored to trigger a response however these can only be controlled via external sources.[7] On the contrary, one can harness intrinsic stimuli such as

63 enzymes overexpressed at disease sites or variations in pH throughout the body, such as that found
64 in the stomach, or the more acidic extracellular microenvironment of tumor cells compared to
65 healthy tissue. [8-10] Acidic environments are also present in other organs/organelles such as the skin
66 surface where the pH is around 4-6.5,[11] <4 in the gastric phase,[12] and between 4.5-6.5 in the
67 lysosomal and endosomal lumen,[13] while the pH in the serum is 7.4. Acidic pH is also used *in vivo*
68 to trigger processes such as protein-ligand binding and viral fusion.[14]
69
70 Whilst the majority of pH responsive systems have been focused on polymers,[15-16] LCNPs are
71 mainly made up of biological amphiphiles rendering them more biocompatible and less cytotoxic.[6,
72 17-18] Non-lamellar forming lipids are some of the most common building blocks to generate inverse
73 lyotropic liquid crystalline phases as they can form a variety of structures depending on
74 composition, thermodynamic variables, pH and other parameters.[19-21] The resulting phase behavior
75 is also dependent on the geometric packing of the lipids and can be tuned by increasing the
76 concentration of more strongly curved lipids which will increase the negative spontaneous
77 curvature towards the water and form phases with increasing negative interfacial curvature.
78 (**Figure 1a**). For example, by increasing the fatty acid content in monoolein (MO) based systems,
79 the structure can transform from Q_{II} to H_{II} and on further addition to a discontinuous inverse
80 micellar cubic phase of space group Fd3m.[22-24] In an analogous fashion, the phase behavior can be
81 tuned by introducing ionisable lipids where the lipid headgroup can be protonated or deprotonated
82 depending on the pH and ionic strength. An increase in surface charge density will cause
83 electrostatic repulsion between headgroups, an increase of the effective headgroup area and
84 consequently a decrease in the magnitude of the spontaneous inverse curvature and the adoption
85 of less curved phases (**Figure 1a**). Similarly, an increase in ionic strength will enhance charge

86 screening and hinder the transition to less curved structures. Tuning the phase behavior of these
87 systems is attractive as the release rate of encapsulated actives from LCNPs strongly correlates
88 with the internal nanostructure symmetry as well as, for some LCNPs, the geometry and
89 dimensions of their pores.[10, 25-26]

90 Most studies have focused on pH-induced transitions between lamellar (Lα) and Q_{II}, or between
91 Q_{II} and H_{II} phases. A MO: linoleic acid system was developed to change from Q_{II} at pH 7 to H_{II} at
92 pH 2 to simulate intestinal and gastric phases respectively.[27] Negrini *et al*. further developed a
93 monolinolein: pyridinylmethyl linoleate system that switched from an H_{II} to a Q_{II} phase when the
94 pH was ≤ 5.5 and showed the potential therapeutic role of this targeted system in treating cancer
95 cells by exploiting their more acidic microenvironment.[28] A low pH-induced Lα to Q_{II} phase
96 transition in a MO: dioleoylphosphatidylserine system showed that the Lα phase directly
97 transformed into an H_{II}, and subsequently, the H_{II} slowly converted into a Q_{II} phase.[29]

98 On the contrary, pH-triggered transitions from ordered inverse micellar phases is still in its
99 infancy. These have an internal structure consisting of discontinuous micellar 'pockets' and, unlike
100 the porous Q_{II} and H_{II} phases, these are based on a discontinuous packing of inverse micelles,
101 containing an array of individual nanoscale aqueous compartments separated by fluid hydrophobic
102 regions, which could potentially offer containment control.
103 Salentinig *et al.* showed that by increasing the pH from 6.8 to 9 at 25°C in MO: oleic acid (OA)
104 dispersions, an inverse micellar solution transformed to Fd3m micellosomes, to hexosomes, to
105 cubosomes and finally to liposomes.[30] Similar phase transitions from LCNPs of OA doped with
106 the antimicrobial human cathelicidin LL-37 peptide were observed upon increasing the pH from
107 6 to 8 at 25°C.[31] Recently Fong *et al.* showed that at a certain MO: fatty acid composition Fd3m
108 micellosomes were formed in water with a pH of 4.9 and very low ionic strength. When water was

replaced by PBS (pH 7.4) the micellosomes transformed to hexosomes at 30°C. This transition was however triggered by the combination of pH and ionic concentration changes, and thus cannot be considered a pure pH response.[32]

In this work we aim to develop LCNPs that are switchable between the confined Fd3m and porous H_{II} phases at physiologically relevant temperature and pHs, confirmed with detailed structural characterization by SAXS and cryo-TEM. We designed pH-responsive LCNPs based on MO and oleyl alcohol (Olalc) doped with a small amount (3 mol%) of a pH-sensitive lipid. We show that these LCNPs transform from Fd3m micellosomes at pH 7.4 to hexosomes when the pH is ≤ 6. To the best of our knowledge this is the first report of a pH triggered LCNP transition from Fd3m micellosomes to hexosomes with decreasing pH and compared to previous research on similar transitions, our system can reversibly transform at both room (25°C) and at a biologically-relevant temperature (37°C). We highlight that high resolution cryo-TEM images have allowed us to directly visualize the internal structure of the Fd3m nanoparticles with unprecedented detail allowing us to compare the internal structures of the nanoparticles with theoretical predictions, using cryo-TEM for this purpose for the first time.

RESULTS AND DISCUSSION

Phase behavior of oleyl alcohol, monoolein and DOBAQ mixtures reveal an H_{II} to Fd3m transition

The phase diagram of MO in water has been widely studied;[33-34] above the excess water point of 35 wt%, MO adopts a Pn3m bicontinuous cubic phase between approximately 20 and 80 °C, which then transforms to an inverse hexagonal (H_{II}) phase at higher temperatures. Generally, for the other

pure component systems, DOBAQ forms liposomes in excess water,[35] whereas oleyl alcohol adopts an inverse micellar solution.[25]

The phase behavior of MO with fatty acids, such as oleic acid (OA), has previously been reported: MO:OA dispersions, stabilized by the Pluronic F127 in milliQ water, typically require an excess of 50 mol% OA to form Fd3m micellosomes,[30, 32] and when MO:OA dispersions are made in PBS at pH 7 and 100mM NaCl (above the pK^{app} of OA and hence the OA is negatively charged), over 80 mol% oleic acid is required to form the Fd3m phase.[36]

Here, we have chosen to study mixtures of monoolein with oleyl alcohol, which is structurally similar to OA but is not ionisable, doped with a small amount of the pH-sensitive lipid N-(4-carboxybenzyl)-N,N-dimethyl-2,3- bis (oleoyloxy) propan-1-aminium (DOBAQ). DOBAQ is neutral at physiological pH and becomes cationic as the pH is lowered, with a pKa of approximately 6 when measured in DOPC:DOBAQ 3:1 liposomes.[35] Preliminary exploration of the phase behavior of MO:Olalc:DOBAQ mixtures found that 3 mol% DOBAQ was optimal to reproducibly switch between the Fd3m and H_{II} phases by lowering the pH. The phase behavior and lattice parameter (a) dependence of MO:Olalc dispersions at a fixed 3mol% DOBAQ at 25 and 37 °C and pH 7.4 are shown in **Figure 1b**. At this fixed DOBAQ ratio the system forms hexosomes between 37 and 39 mol% Olalc. Increasing Olalc to 41 mol% at 37 °C or 42 mol% at 25 ° C transforms the system to Fd3m micellosomes, with a small area of phase coexistence between 40 and 41 mol% Olalc at 25 °C and 40mol% at 37 °C.

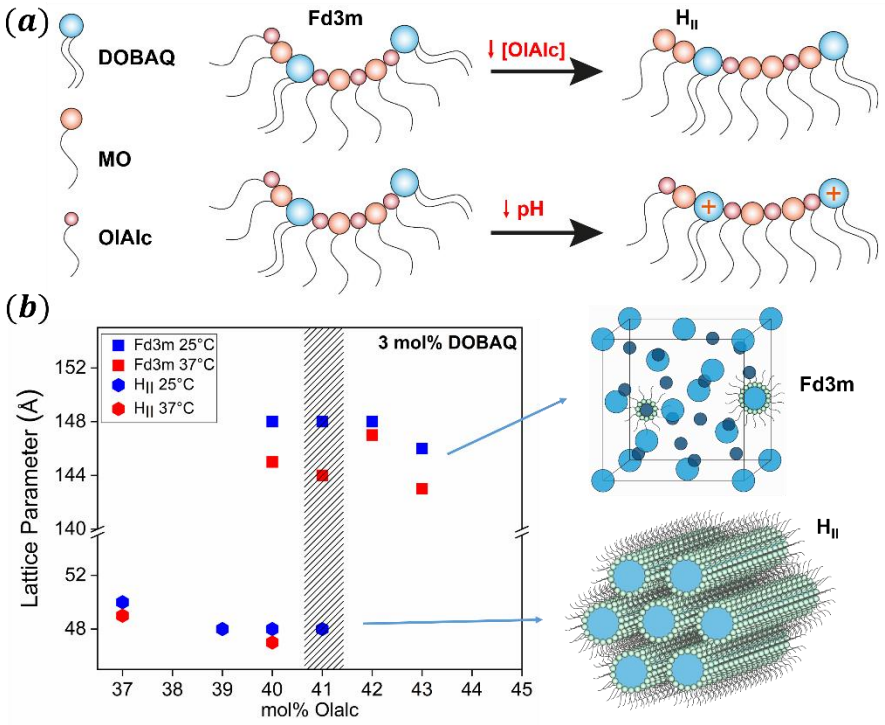


Figure 1. Design parameters to tune the negative interfacial curvature of the nanoparticles. (a) Increasing the amount of the weakly polar ampliphile OlAlc should drive the formation of more curved structures (top). Lowering the pH will cause DOBAQ to become cationic, increase its effective headgroup area and electrostatic repulsion between headgroups which should result in the adoption of less curved phases (bottom). (b) Phase behavior and the effect on the lattice parameter of incorporating Olalc into MO dispersions, keeping the pH sensitive lipid DOBAQ concentration fixed at 3 mol% in all samples at pH 7.4. The dashed region on the graph highlights the composition that was further studied as a function of pH.

Olalc is non-ionisable and has a slightly smaller headgroup compared to OA and hence lower concentrations are required to promote the Fd3m phase. This can be rationalized by increasing the magnitude of the spontaneous inverse monolayer curvature (H_0) towards water,[37] which is related to the structural parameters of the lipids as dictated by the critical packing parameter (CPP) in equation (1).

$$CPP = \frac{v}{a_o l} \quad (1)$$

where v is the molecular volume of the fluid hydrocarbon chain(s), l is the length of a fully extended hydrocarbon chain and a_o is the interfacial area per molecule at the polar-non-polar interface.

Increasing Olalc concentration will decrease the headgroup area driving the formation of structures with increased negative curvature. This is evident in **Figure 1b** where an increase in Olalc causes a transition from hexosomes to the more inversely curved Fd3m micellosomes. Within the pure phase regions a reduction in the lattice parameters of both phases is observed, which is explained by the reduction in water concentration with increasing mean inverse interfacial curvature.

These nanoparticles were engineered and envisioned for future *in vitro/vivo* applications. Consequently, a composition that formed Fd3m micellosomes at physiological temperature was chosen to further study the nanoparticle phase behavior as a function of pH. MO:Olalc:DOBAQ (56:41:3 mol%) was selected as it lies very close to the phase boundary, making it easier to trigger structural transitions between the two phases by changing the pH.

Fd3m to H_{II} transitions can be triggered by decreasing pH at fixed composition

The pH dependent phase behavior as determined by SAXS shows that F127 stabilized MO:Olalc:DOBAQ (56:41:3 mol%) dispersions clearly undergo phase changes in response to pH and demonstrates that the designed system can be tuned to switch between the Fd3m and H_{II} phases (**Figure 2a**). In these experiments, MO:Olalc:DOBAQ dispersions were originally prepared at pH 7.4, aliquots taken and their pH subsequently adjusted by HCl.

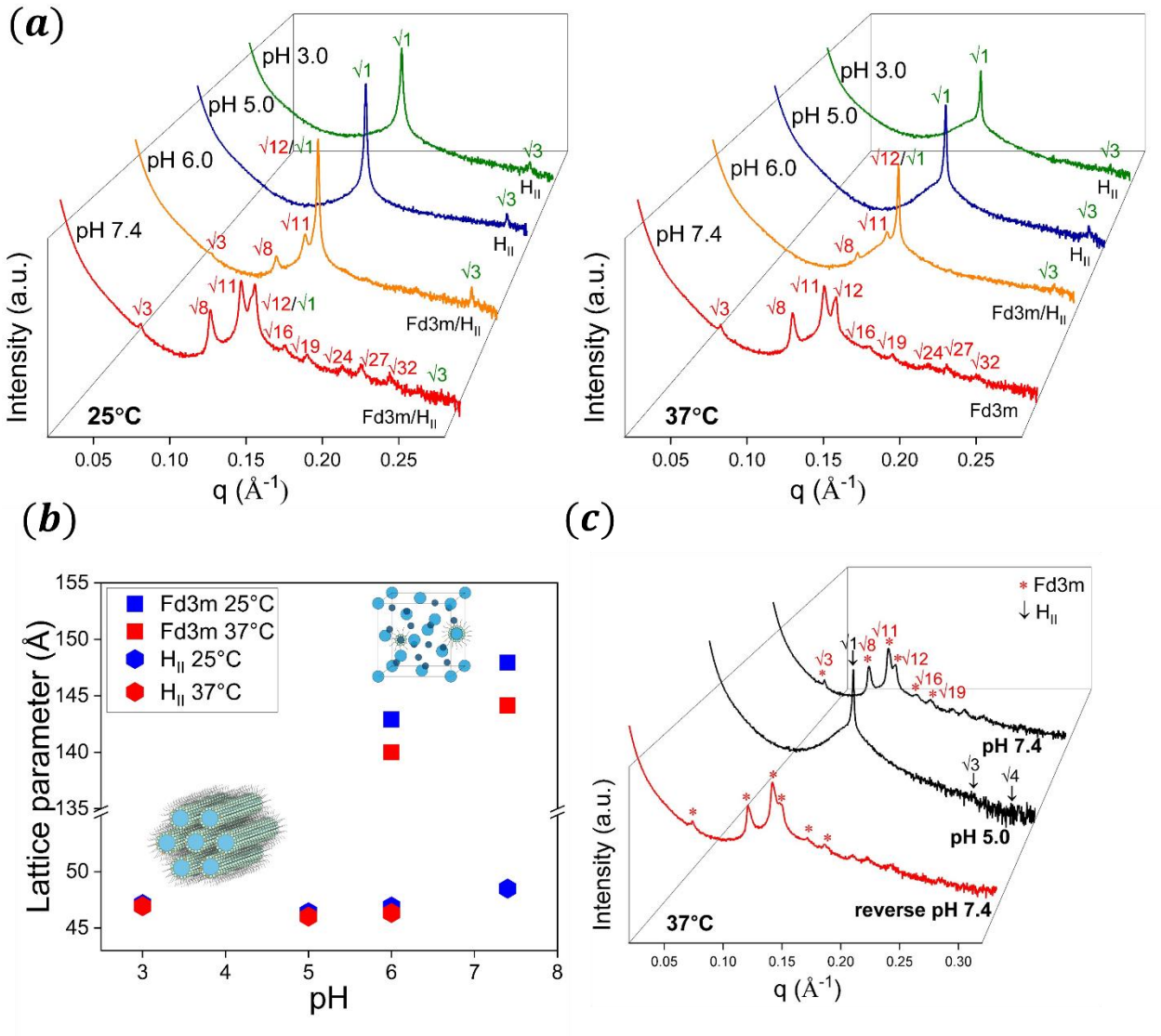


Figure 2. Fd3m to H_{II} transitions can be triggered by decreasing pH and are fully reversible. (a) SAXS patterns of the 5 wt% F127 stabilized 56:41:3 mol % MO:Olalc:DOBAQ nanoparticles in PBS as a function of pH at 25°C (left) and 37°C (right). At pH 7.4, nine Bragg peaks are observed in the ratios of √3:√8: √11:√12:√16:√19:√24:√27:√32 (corresponding Miller indices (hkl) are 111, 220, 131, 222, 400, 331, 422, 333/511 and 440), which index as cubic space group *Fd3m* (red label). At 25°C, the Fd3m phase coexists with a small amount of a 2D inverse hexagonal phase (H_{II}). At pH 6.0, the system adopts an H_{II} phase (green label), with a Bragg peak spacing ratio of √1: √3 (Miller indices (hk) are 10 and 11) indexing to plane group p6mm, coexisting with a small amount of Fd3m phase. When the pH is ≤ 5.0, only H_{II} peaks were observed so the micellar cubosomes have completely transformed to hexosomes. (b) Phase behavior and effect on the lattice parameter of the nanoparticles as a function of pH. The lattice parameters of Fd3m phase decrease with the lowering of pH from 7.4 to 6.0 at both temperatures, then the Fd3m completely transforms to an H_{II} phase when the pH is 5.0. (c) Reversible structural transitions triggered by pH at 37 °C: the pH of the dispersion was gradually decreased from 7.4 to 5.0 (black curve) using HCl and then increased back to 7.4 (red curve) using NaOH. Indexing of the Fd3m (only the first 6 reflections

are indexed for clarity) and the H_{II} Bragg peaks are marked with red stars and black arrows respectively. There is no significant change in the lattice parameter of the Fd3m phase and no additional phases were observed before and after the pH switch, highlighting that the transition is fully reversible.

At physiological pH, the system forms Fd3m micellosomes whereas at acidic pH the nanoparticles transform to hexosomes. This is attributed to the pH-sensitive behavior of DOBAQ, which is neutral at pH 7.4 due to a quaternary amine present in its headgroup and becomes cationic at acidic pH via protonation of its carboxylate group. The protonation of DOBAQ imposes a change in CCP, due to an increase in the surface charge density and enhanced electrostatic repulsion between charged headgroups. This in turn, increases the effective headgroup area and hence decrease the magnitude of the monolayer spontaneous inverse curvature driving the formation of less negatively curved phases.[2, 6, 38]

Figure 2b shows the variation in the structure and lattice parameter of the system with pH and temperature. At 25°C and pH 7.4, the system adopts an Fd3m phase with a lattice parameter of 148 Å. An additional weak reflection occurring around $q \approx 0.15$ Å^{-1} is due to a coexisting H_{II} phase with a lattice parameter of 48 Å. At 37°C and pH 7.4 the system adopts a pure Fd3m phase with a lattice parameter of 145 Å. Upon decreasing the pH from 7.4 to 6.0, the system transforms to hexosomes. This is supported by recording the characteristic √3-reflection of the H_{II} phase (**Figure 2a**) with a small contribution of coexisting Fd3m micellosomes at both 25 and 37°C. We note that diffraction peak intensities scale with the underlying material's volume fraction. Peak intensities can therefore be used to estimate the amount of material adopted by each respective phase. This transition is in agreement with the pKa of DOBAQ measured by Walsh *et al.* in a liposomal formulation (pH ≈ 6.0).[35] We further note, that the Bragg peaks of the Fd3m phase shift to higher q from pH 7.4 to 6, corresponding to a decrease in a_{Fd3m} to 140 Å and 142 Å at 25 °C and 37 °C

respectively. The effect of increasing temperature, which will increase the chain splay and H_0 resulting in a decrease in the lattice parameter of a mesophase, is consistent with the behavior seen here and with most lyotropic liquid crystalline formulations. On the contrary, the observation of a decrease in lattice parameter with an increase in charge is at first instance counter-intuitive, because an increase in charge should reduce H_0 and hence cause the lattice parameter of the Fd3m phase to increase. This trend could be rationalized by considering the size of the nanoparticles which is also highlighted in the subsequent section on cryo-TEM data. As seen in **Figures 3** and **4**, Fd3m micellosomes have a range of sizes and vary in their lattice parameters, with smaller nanoparticles displaying bigger lattice parameters and larger nanoparticles displaying smaller lattice parameters. The larger nanoparticles offer a bigger confinement volume for lyotropic liquid crystals to form, and as a consequence are less prone to structural defects and repulsive membrane undulation forces which will result in smaller lattice spacings and also render them more stable. We hypothesize that at pH 6, the weak reflections from the Fd3m structure are due to a few remaining more stable, relatively bigger Fd3m nanoparticles which have smaller lattice parameters (**Figure 2b**).

At pH 5.0, only the H_{II} phase was observed and hence the system has completely transformed from the initial Fd3m phase. Further decreasing the pH to 3 did not alter the phase behavior and the lattice parameter of the H_{II} phase varied little (maximum ±2Å) across all pH values measured.

pH-triggered Fd3m to H_{II} transitions are reversible

To evaluate the potential to reverse the structural transition in the designed system at 37°C, the pH of the sample was first lowered from pH 7.4 where the system adopted a pure Fd3m phase, to pH 5.0, where only the H_{II} phase was observed. The pH was then subsequently raised back to 7.4 in

the same sample leading to the reformation of the Fd3m phase. SAXS data demonstrate the reversibility of the Fd3m to H_{II} transition by adjusting the pH of the PBS dispersions with NaOH or HCl (**Figure 2c**). The lattice parameter of the sample in the Fd3m phase before (144.9 Å) and after reversing the pH back to pH 7.4 (147.1 Å) is very similar, highlighting the truly reversible responsiveness of the system. The structure reversibility of the system is due to the protonation and deprotonation of the DOBAQ headgroup. When the pH is 7.4, the DOBAQ headgroup deprotonates and becomes neutral whereas at lower pH protonation causes the headgroup to become cationic. Repulsions between the gradually protonating or deprotonating carboxylic group of 3 mol% DOBAQ embedded at the lipid-water interface, modifies the system's spontaneous curvature.

The nanostructure of the designed system can be easily tuned by only a very small amount of pH-sensitive lipid, and shows the excellent pH sensitivity and selectivity of this system compared to previous studies where large amounts of ionisable lipids were required to switch between these two structures and/or significant changes in ionic strength of the buffer. Moreover, whilst an Fd3m to H_{II} transition has been reported in the literature by increasing pH (although none have demonstrated this at 37°C), this is the first example of an Fd3m to H_{II} transition from physiological to acidic pH. This transition could potentially be advantageous in nanomedicine, where a hydrophilic active could be encapsulated into the inverse micelles of the Fd3m phase significantly hindering its release in the serum, but by rearranging to a porous H_{II} phase, e.g. in the endosome, can facilitate release. As an example, the diffusion coefficient of glucose from a system adopting the H_{II} phase was shown to be approximately 15 times that observed in the Fd3m phase in bulk mesophases however how this translates to LCNPs is not known.[26]

Cryogenic Transmission Electron Microscopy (Cryo-TEM) facilitates quantitative structural characterization of individual LCNPs

The cryo-TEM images of the designed MO:Olalc:DOBAQ 56:41:3 mol% system confirmed the pH-triggered nanostructural transitions observed by SAXS. At pH 7.4, the cryo-TEM image indicates the presence of Fd3m micellosomes coexisting with as small amount hexosomes, with the Fd3m micellosomes dominating in the dispersions (**Figure 3a**) at 25 °C. By contrast, only hexosomes were observed at pH 5.0 (**Figure 3b**).

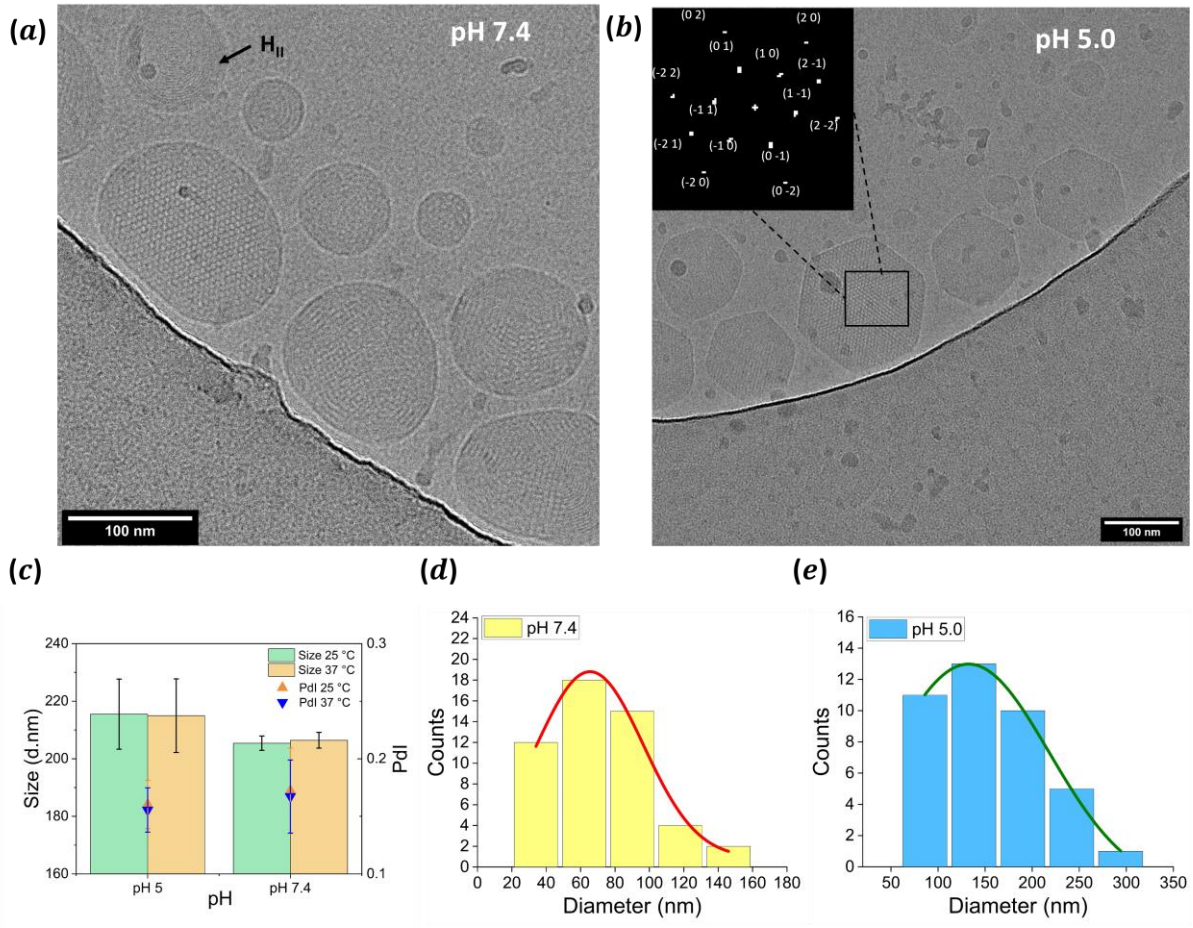


Figure 3. Representative cryo-TEM images of the MO:Olalc:DOBAQ 56:41:3 mol% system prepared at 25 °C and their size distributions. (a) Fd3m micellosomes at pH 7.4 coexisting with a small amount of hexosomes (black arrow), which likely give rise to the weak intensity of the √1 peak in the SAXS data in Figure 2a. Both SAXS and cryo-TEM data are in agreement that the

dominating phase at pH 7.4 at 25 °C is Fd3m. (b) pure hexosomes at pH 5.0. The insert in panel b shows the corresponding intensity of the fast Fourier transform (FFT) applied to the black box region of the individual nanoparticle and the assigned Miller indices. (c) Dynamic light scattering (DLS) data of MO:Olalc:DOBAQ 56:41:3 mol% nanoparticles in PBS at different pH at 25 °C (green) and 37 °C (orange). At 25 °C, Fd3m micellosomes at pH 7.4 have a mean size of ~ 205 nm, whilst hexosomes are ~216 nm at pH 5; At 37 °C Fd3m micellosomes have a mean size of ~206 nm and hexosomes ~215nm. (d) and (e) Cryo-TEM size distribution of nanoparticles hydrated at pH 7.4 and pH 5.0 respectively. Histograms were calculated from 51 nanoparticles at pH 7.4 and 40 for pH 5.0 and fitted to a Gaussian function. From the fit we obtain a mean of 65 and 132 nm for the nanoparticles at pH 7.4 and pH 5 respectively. Both DLS and Cryo-TEM data show nanoparticles at pH 5.0 generally have a slightly larger size compared to that at pH 7.4.

The particle size distribution of the nanoparticles at 25 and 37 °C was investigated using dynamic light scattering (DLS) and the results are shown in **Figure 3c**. The Fd3m micellosomes at pH 7.4 have a mean size of ~ 205and 206 nm, whereas the hexosomes at pH 5 have a mean size of ~216 and 215 nm, at 25 and 37 °C respectively. The DLS results at 25 °C were compared with the nanoparticle size analysis from Cryo-TEM data (51 nanoparticles analyzed at pH 7.4 and 40 at pH 5.0) shown in **Figure 3d** and **3e**. The histograms were fitted to a Gaussian function and we obtain a mean size of 65 and 132 nm for the nanoparticles at pH 7.4 and pH 5, respectively. The mean size of the nanoparticles calculated differs between the two techniques, as DLS weighs the size distribution differently with larger particles emphasized as they scatter light more strongly whereas cryo-TEM often excludes larger particles from the thin ice films, emphasizing smaller particles in these polydisperse populations. Nonetheless, there is a clear trend that hexosomes are larger than the micellar cubosomes. This is because a decrease in pH leads to protonation of the DOBAQ headgroup resulting in a repulsive force at the lipid-water interface which pushes the elongation of the Fd3m micellar 'pockets' to form a tubular structure in the H_{II} phase and an overall increase in the water concentration. However, the overall increase in water volume is only about 6% when hexosomes are formed, which would only account for a 2% increase in particle size (for

estimations please see **Supporting information** Section 1), thus this cannot be the only explanation for the more strongly differing particle sizes. On a mesoscopic level though, there is an unneglectable difference. While the micellosomes are spherical in shape and hence display a maximized volume to surface ratio, hexosomes are by no means perfectly spherical. They display different shapes, like flat hexagonal prisms (**Figure 3b**), and often possess even "spinning top-like" structures.[39] In any case, their shape deviates significantly from that of a sphere, and hence, the apparent mean hydrodynamic size of hexosomes is bigger, whilst their volume has not increased by much. This apparent size increase of hexosomes is even more pronounced in our cryo-TEM measurements, since we measured the diameter of the hexagonal prism section, whilst not considering the perpendicular short extension of the hexosomes.

Traditionally, cryo-TEM cannot be used to unambiguously determine the space group of lyotropic liquid crystalline phases due to the low resolution of the images leading to only a couple of peaks to be observed after applying a fast Fourier transform to the image.[40] However, it is a powerful complementary technique to SAXS. Whilst cubosomes based on bicontinuous cubic phases, and hexosomes are easier to visualize with cryo-TEM and give rise to a number of peaks in their FFTs,[41-42] cryo-TEM images of Fd3m micellosomes reported to date in the literature are of very low resolution, which is insufficient for the visualization of the internal nanostructure of the phase.

Here we have been able to obtain high resolution images and directly visualize the internal structure of the Fd3m nanoparticles along two viewing orientations, which allow us to resolve the large and small micelles of the Fd3m structure and calculate structural information using cryo-TEM for the first time.

Figure 4 shows detailed images of Fd3m nanoparticles in two different orientations, image analysis, indexing and simulated TEM diffraction patterns. **Figure 4a** highlights an Fd3m micellosome along the [110] viewing direction with respect to the electron beam. The simulated TEM diffraction pattern from an Fd3m phase diffraction pattern viewed down the (110) plane normal (**Figure 4b**) is in excellent agreement with the experimentally determined FFT pattern. **Figure 4c** shows the FFT from the image in **Figure 4a** in which the √3, √8, √11 and √12 reflections, corresponding to Miller planes (111), (220), (311) and (222) can be clearly observed. Each theoretical reflection from the simulated diffraction pattern indexed in **Figure 4b** corresponds to the same reflection in **Figure 4c**. The lattice parameter of the nanoparticle in **Figure 4a** was determined to be 149.3 Å (see also **Supporting Information** Section 2 for additional information on determining the lattice parameter).

Figure 4e shows an Fd3m nanoparticle aligned along the [114] direction with respect to the electron beam in which the √8, √11, √27 and √32 reflections, corresponding to Miller planes (220), (311), (511) and (440) were observed. Each theoretical reflection from an Fd3m phase viewed down the (114) plane normal and indexed in **Figure 4f** directly matches the experimentally determined reflections in **Figure 4g**. The lattice parameter of the particle in **Figure 4e** was determined to be 137.0 Å based on the FFT analysis.

The lattice parameters of the individual nanoparticles calculated from the cryo-TEM data are slightly different to those obtained by SAXS, which is expected, as SAXS gives the average lattice parameter of all the nanoparticles in the measured sample whilst we focus on specific nanoparticles for the cryo-TEM analysis. Based on the cryo-TEM images collected, we observe that generally nanoparticles which have a smaller lattice parameter tend to be larger in size. This finding supports our hypothesis that smaller Fd3m nanoparticles undergo a faster structural transition compared to

larger particles, as they are expected to contain relatively more defects (cp. discussion above) and are thus less stable.

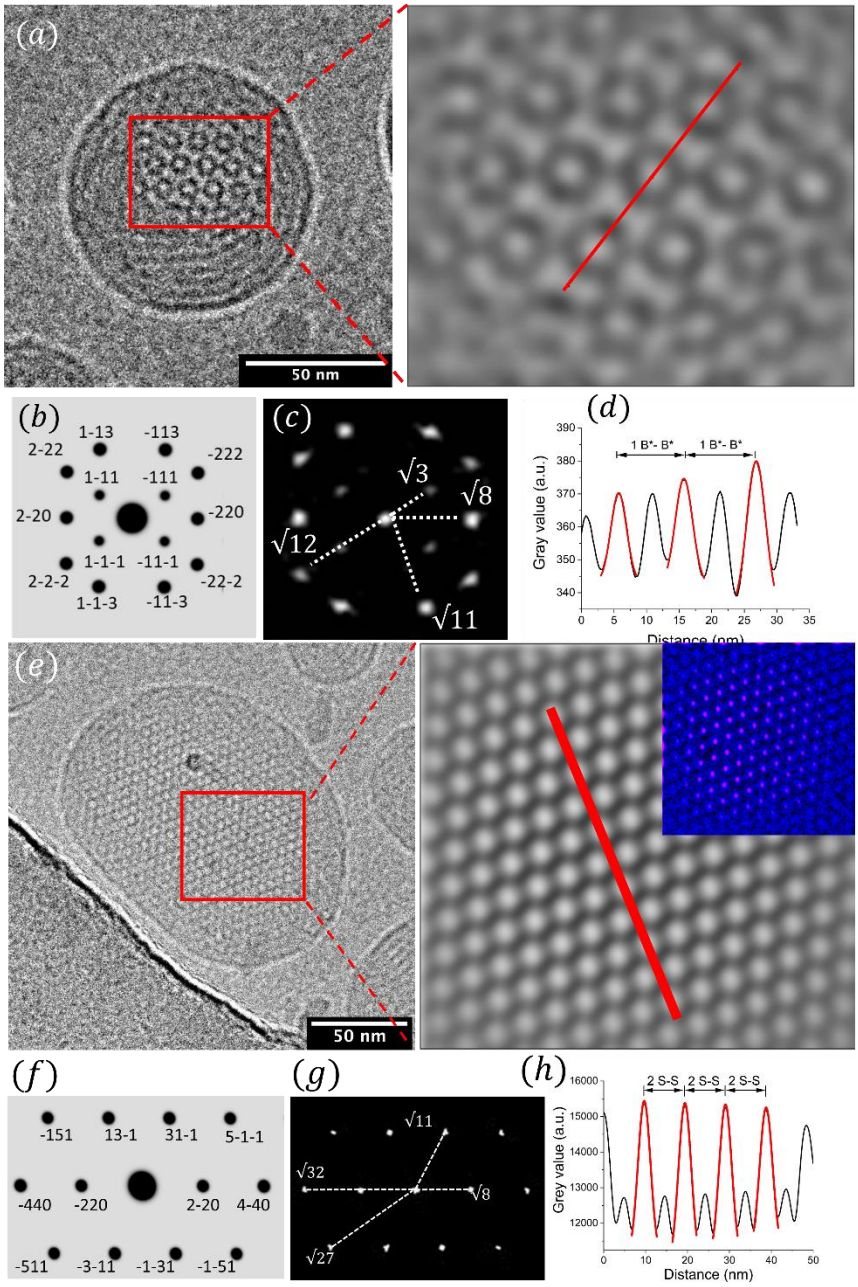


Figure 4. High resolution cryo-TEM images allow direct visualization of the internal structure of the Fd3m nanoparticles. (*a*) Cryo-TEM image on an Fd3m micellosome of MO:Olalc:DOBAQ 56:41:3 mol% prepared at 25°C and pH 7.4.The nanoparticle is aligned along the [110] orientation with respect to the electron beam. The red box indicates the region used to apply the FFT shown in 4c and the zoomed-in region on the right hand side shows the inverse FFT of 4c (after application

of rolling background subtraction and Gaussian blur), and was used to determine the $d_{B^*-B^*}$ distance indicated by the red line. (b) Simulated transmission electron microscopy (TEM) diffraction pattern viewed down the (110) plane normal of an Fd3m phase with assigned Miller indices and (c) the corresponding intensity of the FFT applied to the red box region of 7a which is in excellent agreement with the simulated pattern. Each theoretical reflection from the simulated pattern indexed in 4b corresponds to the same reflection in 4c. (d) Gaussian fitting of the intensity profile along the red line of the inverse FFT, shown in the zoomed in region of 4a, gave an average distance between two interval big micelles ($d_{B^*-B^*}$) of 105.5 Å. The lattice parameter of the nanoparticle in 7a was determined to be 149.3 Å, which gives a theoretical $d_{B^*-B^*}$ value of 105.5 Å. (e) cryo-TEM image of an Fd3m nanoparticle aligned along the [114] orientation with respect to the electron beam. The red box indicates the region used to apply the FFT shown in 4g and the zoomed-in region on the right hand side shows the inverse FFT of 4g (after application of rolling background subtraction and Gaussian blur), and was used to determine the d_{S-S} distance indicated by the red line. The inset shows the overlay of the raw image of 4e on the left shown in blue and the inverse FFT on the right shown in magenta which are in excellent agreement. (f) Simulated transmission electron microscopy (TEM) diffraction pattern viewed down the (114) plane normal of an Fd3m phase with assigned Miller indices and (g) the corresponding intensity of the FFT applied to the red box region of 8a which is in excellent agreement with the simulated pattern. Each theoretical reflection from the simulated pattern indexed in 4f corresponds to the same reflection in 4g. (h) Gaussian fitting of the intensity profile along the red line of the inverse FFT, shown in the zoomed in region of 8a, gave an average plane to plane distance of the small micelles ((d_{S-S}) of 48.5 Å. The lattice parameter of the nanoparticle in 8a was determined to be 137 Å, which gives a theoretical d_{S-S} value of 48.4 Å.

The high resolution cryo-TEM images have allowed us to direct visualize the large and small micelles of the Fd3m structure (eight large and sixteen small micelles per unit cell), determine the micelle positions and compare them to theoretical distances for the Fd3m structure.

From geometry,[43] the big (B) micelle-to-micelle distance is given by equation 2:

$$d_{B-B} = \frac{\sqrt{3}}{4}\, a \qquad (2)$$

and the small (S) micelle-to-micelle distance is given by:

$$d_{S-S} = \frac{\sqrt{2}}{4}\, a \qquad (3)$$

Figure 4a shows an Fd3m micellosome aligned along the [110] orientation. The distribution of the different micelles along the (110) section will depend on where the section cuts through the

plane along the z direction. Here only the big micelles are seen, however this does not represent the d_{B-B} micelle-to-micelle distance defined above which is the smallest distance between two adjacent big micelles, but instead what we define as the d_{B*-B*} distance, which is the distance between two interval big micelles (in a hexagonal arrangement) and is given by equation 4 (see also **Supporting Information** Section 3):

$$d_{B*-B*} = 2d_{S-S} = \sqrt{(\frac{8}{3})}d_{B-B} = \frac{\sqrt{2}}{2}a_{Fd3m} \quad (4)$$

Further discussion on determining the d_{B*-B*} distance as well as comparing the excellent agreement of the inverse FFT (**Figure S2a**) obtained from **Figure 4c** with some 2-D serial sections of the 3-D electron density map from a previously published DOPC:DOG 1:2 mol% system adopting an Fd3m phase (**Figures S2b** and **S2c**)[44] can be found in **Supporting Information** Section 3.

Experimentally, the B*-B* distance was determined by taking the inverse FFT of **Figure 4c**, which has been background subtracted. The resulting, Gaussian blurred, inverse FFT shown on the right of **Figure 4a** was used to determine the B*-B* (d_{B*-B*}) distance by Gaussian fitting (**Figure 4d**) of its intensity profile along the red line in **Figure 4a** (right). The average d_{B*-B*} distance measured was 105.5 Å which is identical to the theoretical distance of 105.5 Å based on the experimentally determined lattice parameter of 149.3 Å (equation 4).

Figure 4e shows an Fd3m micellosome aligned along the [114] orientation and highlights the small micelles of the Fd3m phase. Based on the reflections of the FFT in **Figure 4g**, the lattice parameter was determined to be 137 Å, which gives a theoretical d_{S-S} micelle-to-micelle distance of 48.4 Å (equation 2). Note the √8 refection also corresponds to the S-S micelle distance. The d_{S-S} distance was determined by taking the inverse FFT of **Figure 4g**, which has been background

subtracted. The resulting, Gaussian blurred, inverse FFT is shown in **Figure 4e** (right) was used to determine the d_{S-S} distance by Gaussian fitting (**Figure 4h**) of its intensity profile along the red line in **Figure 4e** (right). The overlay of the raw image of **Figure 4e** (left) shown in blue and inverse FFT (from **Figure 4g**) shown in magenta are in excellent agreement (**Figure 4a,** right, inset). It should be noted that in the [114] direction, d_{S-S} does not correspond to the distance between two adjacent small micelles as shown in Figure **Figure S2b** for the (110) section but is instead the plane to plane distance of the small micelles in this projection (**Figure 4e** (right), see also **Figure S3**). The experimentally determined S-S plane distance (**Figure 4h**) and the (220) reflection from **Figure 4g** were 48.5 and 48.3 Å respectively and match the theoretical distance of 48.4 Å determined above.

CONCLUSION

We have been able to generate highly sensitive pH responsive lyotropic liquid crystalline nanoparticles that require only a very small amount (3 mol%) of pH sensitive lipid to switch between a compartmentalized internal structure based on space group Fd3m at physiological pH to an accessible porous 2D hexagonal structure when the pH is ≤6 at both at 25 and 37°C. The transition is fully reversible by tuning the inverse spontaneous curvature by the protonation or deprotonation of the pH sensitive lipid's headgroup. The internal symmetry of the nanoparticles was confirmed by SAXS and cryo-TEM. Imaging of Fd3m nanoparticles has been challenging with cryo-TEM however we have managed to acquire high resolution images of them and directly visualize the structure in unprecedented detail for the first time. Depending on the direction the particles are aligned with respect to the electron beam we can resolve the large and small inverse micelles of the Fd3m phase which allowed us to calculate structural information such as the

positions and distances of the two different sized micelles and compare these to theory based on geometrical considerations. These promising findings could pave the way to developing superior stimuli-responsive soft nanoparticle formulations that change their connectivity upon encountering acidic pH, which is known for example to be highly important both in human disease as well as plant physiology and crop performance. They can potentially benefit a range of downstream applications in the biomedical, food, agrochemical and environmental remediation industries.

MATERILAS AND METHODS

Preparation of MO/Olalc/DOBAQ nanoparticles. Monoolein (MO; was a kind gift from Croda Personal Care, Goole, UK), Oleyl alcohol (Olalc; Sigma Aldrich, Gillingham, UK), N-(4-carboxybenzyl)-N,N-dimethyl-2,3-bis(oleoyloxy)propan-1-aminium (DOBAQ; Avanti polar lipids, AL, USA) were weighted and dissolved in chloroform (Sigma-Aldrich, Gillingham, UK), then mixed at the desired MO:Olalc:DOBAQ molar ratios. MO had a purity of >92% whilst the rest of the lipids had a purity of >99% and were used without further purification. The chloroform solutions were dried using a stream of N_2 gas and the dry thin films were left under vacuum overnight to remove any residual organic solvent. Thin films were then heated to 60°C and hydrated with phosphate-buffered saline buffer (PBS; Sigma-Aldrich, Gillingham, UK) containing 6.923 mg/ml Pluronic F127, 5 wt% F127 of total lipids, (Sigma-Aldrich, Gillingham, U.K.) to achieve a sample concentration of 100 mg/ml. Dispersed samples were prepared by tip sonication (VibraCell 750 with a stepped microtip, Sonics & Materials Inc, Newtown, USA) for 30 minutes in pulse mode (1s pulse, 1s break) at 35% of its maximum power (750 W). The pH of the NP dispersions was readjusted to 7.4 using 0.1M NaOH (Fisher Scientific, Loughborough, UK) or 0.1M HCl (Fisher Scientific, Loughborough, U.K.). Samples were allowed to equilibrate at 37°C

overnight before further experiments. The pH of the dispersed nanoparticles was first gradually decreased from 7.4 to 3.0, then increased back to 7.4 while taking out samples at each selected pH point for SAXS measurements at both 25 °C and 37 °C. For the dynamic light scattering and cryogenic transmission electron microscopy data, lipid nanoparticles were hydrated by PBS at pH 7.4 or pH 5.0 directly.

Synchrotron Small-Angle X-ray Scattering (SAXS). Synchrotron SAXS data were collected on beamline I22 at Diamond Light Source. The synchrotron X-ray beam was set to 18 keV which corresponds to a wavelength of 0.69 Å. The sample to detector distance was set at 6 or 8.7 m and the 2-D powder diffraction patterns were recorded on a Pilatus 2M detector. SAXS data were analyzed using the DAWN software.[45-46] The lattice parameters of the Fd3m and H_{II} phases were calculated using the equations below:

$$a_{Fd3m} = \frac{2\pi}{q_{hkl}}\sqrt{h^2 + k^2 + l^2} \quad (5)$$

$$a_{H_{II}} = \frac{4\pi}{\sqrt{3}\, q_{hk}}\sqrt{h^2 + k^2 + hk} \quad (6)$$

where q_{hkl} or q_{hk} is the scattering vector and hkl or hk are the Miller indices of a given Bragg peak. Samples were equilibrated at 25 °C and 37 °C for 5 min before taking measurements.

Dynamic Light Scattering (DLS). For DLS measurements, particles were made at a concentration of 5mg/mL. Samples were diluted to 25 μg/mL by isotonic PBS at different pH before DLS measurement. The mean hydrodynamic radius and polydispersity index width (PDI) of the nanoparticles were measured at pH 7.4 and pH 5.0 using a Malvern Zetasizer Nano ZSP

instrument (Malvern Panalytical, Malvern, UK) equipped with a 633 nm helium-neon laser. Samples were measured at 25 °C or 37 °C at a fixed 173° backscattering angle. The refractive index and absorption of material is set as 1.46 and 0.001 respectively. The refractive index of dispersion is 1.334 and its viscosity is 0.9110 mPa.s at 25 °C and 0.7130 mPa.s at 37 °C.

Cryogenic-transmission electron microscopy (Cryo-TEM). Samples were diluted once (final concentration of 50 mg/ml) in isotonic PBS at the required pH for the morphological characterization using cryo-TEM. In brief, 3 μL of the diluted sample was applied to freshly glow discharged Cu QUANTIFOIL grids (R2/R2, 300 or 400 mesh) with a hold time of 30 s. The carbon-coated grids were glow discharged at 10 mA for 20 s for samples at pH 7.4 and 30 s for samples at pH 5.0. The grids were blotted for 6 s with a blotting force of 7 at 25 °C under 100% relatively humidity, then plunged into liquid ethane using a Vitrobot™ mark IV (Thermo/FEI). Grids were kept under liquid nitrogen before measurement.

Data acquisition was carried out on a Titan KRIOS microscope (Thermo Fisher Scientific, US) with an accelerating voltage of 300 KV and a defocus value of -1 μm at a nominal magnification of 45k or 75k. The pixel size for these images were 1.76 and 1.065 Å, respectively. Image processing was done using Fiji.[47] The Fast Fourier Transforms (FFTs) of the raw images were background subtracted. The d-spacing of each reflection in the FFT was determined using TrackMate.[48] The micelle distances were calculated from the inverse FFTs after application of rolling background subtraction and Gaussian blur using Fiji.

Simulation of TEM diffraction patterns

The program CrystalMaker® (CrystalMaker Software Ltd, Oxford, England- www.crystalmaker.com) was used to visualize 2-D sections (110 and 114 planes) through the model structure of the Fd3m cubic phase. Fd3m origin 2 (at $\bar{3}m$) was used. These sections were then sent to the SingleCrystal™ program to simulate transmission electron microscopy (TEM) diffraction patterns.

ASSOCIATED CONTENT

Supporting Information

Discussion on the estimation of micellosome and hexosome size ratios, details on the calculation of lattice parameters from cryo-TEM data and determination of the micelle positions in Fd3m nanoparticles.

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Author Contributions

A.I.I.T. conceived and designed the experiments. J.M.S, P.A.B., M.R. and A.I.I.T contributed to study design. Z.X. performed all the experiments. Z.X., M.R. and A.I.I.T. analyzed the data. Z.X., P.A.B., M.R. and A.I.I.T. interpreted the results. Z.X. M.R. and A.I.I.T co-wrote the manuscript. All authors discussed the results and commented on the manuscript.

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Notes

The authors declare no competing financial interest.

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