On the Lithium Distribution in Halide Superionic Argyrodites by Halide Incorporation in Li_{7-x}PS_{6-x}Cl_x

Ajay Gautam^a, Michael Ghidiu^a, Emmanuelle Suard^b, Marvin A. Kraft^c,

Wolfgang G. Zeier*^{c,d}

^aInstitute of Physical Chemistry, Justus-Liebig-University Giessen, Heinrich-Buff-Ring 17, D-35392 Giessen, Germany.

^bInstitut Laue-Langevin, 71 avenue des Martyrs, 38042 Grenoble, Cedex 9 - France

^cInstitute of Inorganic and Analytical Chemistry, University of Münster, Corrensstrasse 30, 48149 Münster, Germany

^dHelmholtz Institute Münster (IEK-12), Forschungszentrum Jülich GmbH, Corrensstrasse 46, 48149 Münster, Germany

E-mail:

*wzeier@uni-muenster.de

Abstract

Superionic lithium argyrodites are attractive as solid electrolytes for all-solid-state-batteries. These materials of composition Li_6PS_5X (X = Cl, Br, and I) exhibit structural disorder between the X⁻/S²⁻ positions, with higher disorder realizing better Li⁺ transport. Further replacement of the sulfide by chloride anions (for the series $Li_{7-x}PS_{6-x}Cl_x$) has been shown to increase the ionic conductivity. However, the underlying changes to the lithium substructure are still relatively unknown. Here we explore a larger range of nominal halide compositions in this material from x = 0.25 to x = 1.5 and explore the changes with neutron diffraction and impedance spectroscopy. The replacement of S²⁻ by Cl⁻ causes a lowered average charge in the center of the prevalent Li⁺ "cages", which in turn causes weaker interactions with Li⁺ ions. Analysis of neutron diffraction data reveals that the increased Cl⁻ content causes these clustered Li⁺ "cages" to become more interconnected, thereby increasing Li⁺ conductivity through the structure. This study explores the understanding of the fundamental structure–transport correlations in the argyrodites, specifically structural changes within the Li⁺ ion substructure upon changing anionic charge distribution.

1. Introduction

All-solid-state-batteries are attracting significant interest for next-generation lithium-ion batteries due to their device safety and improved energy density, as well as wider operating temperatures compared to conventional lithium-ion batteries.^{1–5} Solid electrolytes typically used for solid state batteries are oxides^{6,7} phosphates,^{8,9} lithium-based halides^{10–13} and lithium thiophosphates^{14–16}; among all of them, thiophosphates are promising candidates for solid-state batteries because of their high ionic conductivity and mechanical softness.^{2,17,18} The lithium argyrodites Li₆PS₅X (X = Cl, Br, I) especially are gaining more attention.^{19,20} This is to a large extent due to the fact that elemental substitution and changing synthetic approaches have shown to improve the ionic conductivity of these materials,^{21–25} which makes the fabrication of solid state batteries with thick electrodes improving the energy densities of all-solid-state-batteries.

Considering the halide-containing argyrodites $Li_{6-x}PS_{5-x}(Cl, Br, I)_{1+x}$, ^{19,26} these can be thought of as derivatives from the simple Li₇PS₆; this otherwise orthorhombic material can only be found in its superionic cubic polymorph above 210 °C. 27,28 The structure of $\rm Li_7PS_6$ consists of $\rm PS_4{}^{3-}$ tetrahedra and unbound S^{2-} anions; in the halogen-substituted phases, the halide X takes the place of some of these 'free' sulfide anions and stabilizes the cubic polymorph at room temperature. The cubic Li₆PS₅Cl is shown in Figure 1a, with the halide anion X^{-} located in a cubic close-packed lattice (Wyckoff 4*a*). P is bonded to four ions forming PS_4^{3-} polyhedra in which the P center occupies the octahedral sites (P and S occupy the 4b and 16e Wyckoff position). 'Free' S^{2-} ions occupy half of the tetrahedral sites (Wyckoff 4d). Analogous to the other halide argyrodite structures, the lithium ions occupy a number of different sites shown in (Figure 1b), so-called T5 (Wyckoff 48h), and T2 (Wyckoff 48h) to form a cage-like arrangement around the 4d sites; recent publications have identified the possible diffusion pathway through the argyrodite structure mostly involving the T2 - T2 distance as the main connecting step from the cage to cage, shown in Figure 1c.^{29,30} Differences in the Li⁺ positions, the Li substructure, has recently been well summarized by calculating the radial distribution of Li^+ around the 4d site, and determining the radius where the mean amount of Li sits, designated as R_{mean} .^{26,29} Importantly, depending on composition and treatment history, 26,29,31 the halide nominally on 4a and the sulfide on 4d can exchange positions, leading to *site-disorder*. This has been invoked as a key factor in changing unit cell volume as well as transport properties.^{24–26,29,32,33} Site-disorder has been shown to affect the lithium substructure

by modifying the Li⁺ site locations and occupancies, and in turn changing R_{mean} . A reason for some of these effects is the changing average site charge on 4a and 4d because of the electrostatic difference between S^{2–} and X⁻.^{29,30}



Figure 1: a) Crystal structure of Li_6PS_5Cl shown in anion site-disordered state with 40% Cl⁻ located on the Wyckoff 4a site and the 40% sulfide anion (S^{2-}) on the Wyckoff 4d site. b) Lithium substructure in Li_6PS_5Cl argyrodites, showing T5 (Wyckoff 48h), T2 (Wyckoff 48h) sites to coordinate around an S^{2-} central anion on the Wyckoff 4d site. T2 - T2 offer the shortest intercage jump distance ("red arrows"). The average radial distribution of Li away from the Wyckoff 4d central anion is described by R_{mean} . c) Visual representation of the Li connectivity in which the T2 - T2 jump offers the shortest distance at a larger R_{mean} .

It has been shown that the different type of halide ion can influence ionic transport and whether and to what extent disordering occurs.¹⁹ While iodide versions exhibit virtually no sitedisordering,²⁶ chloride substitutions seem to encourage the highest amount of site-disorder and the highest conductivities.^{19,26} With the multiple 'free' anion sites, from a starting point of Li₇PS₆ and a full chloride occupancy endpoint of Li₅PS₄Cl₂ (full replacement of S²⁻ anions), there is a wide compositional space in Li_{7-x}PS_{6-x}X_x (X = Cl, Br, I) to explore, given the potential for anionic sitedisordering and changing Li carrier density.²⁵ Nazar *et al.* recently showed the substitution series Li_{7-x}PS_{6-x}Cl_x ($x \le 0.5$; functionally equivalent to Li_{7-x}PS_{6-x}Cl_x with an offset in x of unity) with a conductivity of 9 mS cm⁻¹ achieved at x = 0.5 (measured as x = 0.45).³³ In addition to increasing halide, Hu *et al.* reported the Br substitution series Li_{6-x}PS_{5-x}Br_{1+x} ($x \le 0.7$) in which a 10 mS cm⁻¹ conductivity was achieved for sintered Li_{5.3}PS_{4.3}Br_{1.7}.³⁴ The conductivities of these phases are an improvement over the 'standard' composition at x = 0, Li₆PS₅Cl/Br, so clearly they are of interest. However, what is missing is an expansion of the range of x to include compositions closer to Li_{6.75}PS_{5.75}Cl_{0.25} and a full structural study of the changing Li⁺ substructure, and how it relates to changes in transport. This will lead to a more fundamental understanding of how these substructures play a role in optimizing conductivity outside the more oft-studied composition of Li₆PS₅X.

In this work, we explore the entire compositional space of $Li_{7-x}PS_{6-x}Cl_x$, across a range of $(0.25 \le x \le 1.5)$. With this series we are able to perform neutron diffraction to characterize how the lithium substructure changes. This helps gain a deeper understanding of why conductivities are so strongly affected by changes in Cl content. With increases in Cl, disorder on the 4*d* site increases, as well as the expansion of "cages" of Li^+ around the site, all of which seems to be beneficial to the ionic transport.

2. Experimental Section

Synthesis. All synthesis work and sample treatment for Li_{7-x}PS_{6-x}Cl_x was performed under an argon atmosphere ($O_2 < 1$ ppm and $H_2O < 1$ ppm). Lithium sulfide (Li₂S, Alfa-Aesar, 99.9%), phosphorus pentasulfide (P₄S₁₀, Merck, 99%), and lithium chloride (LiCl, Alfa Aesar, 99.9%) were mixed to obtain 3 g batches of the precursor. The obtained precursor was put into a 45 ml ball milling cup using a 40:1 mass ratio of milling media to precursors (5 mm diameter zirconia milling balls, Fritsch Pulverisette 7 premium line). The precursors were then milled for a total of 60 cycles, each with 10 minutes of milling time followed by 10 minutes of rest. After every 20 cycles, the ball milling cups were opened inside the glovebox to dislodge material adhered to the milling cup inside surfaces. The completely milled material was then manually hand pressed into pellets and filled into ~ 8 cm long quartz ampoules (10 mm inner diameter, first heat treated at 800 °C under a dynamic vacuum for 2 hours to remove all traces of water). The quartz ampoules were sealed under static vacuum. For crystallization, the ampoules were inserted into a tube furnace and heated to specific temperatures: 550 °C for $x \le 1.0$, and 450 °C for $1.25 \le x \le 1.75$ (ramping rate of 100 °C per hour followed by 2 hours dwell time, then natural cooling. This reaction time served to ensure the complete formation of argyrodites. The final product was hand-ground into a powder for neutron diffraction studies and pressed into the cell for electrochemical measurements.

Neutron powder diffraction. The neutron powder diffraction data were collected on the D2B highresolution powder diffractometer at the Institute Laue-Langevin (ILL, Grenoble, France), with incident monochromatic wavelength $\lambda = 1.594$ Å. Starting point scan was repeated several times at various $10^{\circ} < 2\theta \le 160^{\circ}$ values to improve statistics, and each data collection required 4 hours.

Rietveld Analysis. Rietveld refinements were performed using the TOPAS-Academic V6 software package. Initially, the structural information obtained from the neutron refinement of Li₆PS₅Cl from Minafra *et al.* was used.²⁶ These steps were used during the refinements: (1) scale factor, (2) 10 coefficients for a Chebyshev function background, (3) peak shape parameters, (4) lattice parameters, (5) zero error, (6) fractional atomic coordinates, (7) atomic occupancies and finally (8) isotropic atomic displacement parameters. Then multiple correlated parameters were refined simultaneously and constraints table can be found separately in Table S1.

Bond Valence Sum Analysis. The bond valence sum analysis was carried out using the softBV software tool combine with grid size of 0.1 Å. Initially the crystallographic information file extracted from the neutron diffraction data was used to obtain the energy landscape.³⁵ The BVS models were used for finding the transport pathway. More details of the BVS model can be found in ref 35.

Electrochemical Impedance spectroscopy. The ionic conductivities were measured using AC impedance spectroscopy. A steel press-cell was used for all impedance measurements following previous work.³⁶ All compositions were pressed uniaxially with a pressure of 382 MPa for 3 minutes. The resulting densities of the pellets were ~81%. All Electrochemical impedance spectroscopy (EIS) measurements were performed using a VMP300 impedance analyzer (Biologic) with a frequency range of 7 MHz to 100 mHz with an amplitude of 10 mV in a temperature range of $-40 \,^{\circ}$ C to $60 \,^{\circ}$ C. The analyses of obtained spectra were performed using the RelaxIS 3 Impedance Spectrum Analysis software (RHD instruments, Darmstadt, Germany).

2. Results and Discussion

Structural changes. Mechanochemical synthesis was used to synthesize the series $Li_{7-x}PS_{6-x}Cl_x$ ($0 \le x \le 1.75$) and to study how changes in transport properties can be understood through changes to the crystal structure and Li^+ substructure. The precursors Li_2S , P_4S_{10} , and LiCl were mechanically milled to achieve uniform mixing. For subsequent crystallization, these compositions needed to be heated at different temperatures, with 550 °C for samples x = 0.0 to 1.00, or at a reduced temperature of 450 °C (for samples x = 1.25, 1.50 and 1.75). For the structural characterization, neutron diffraction was carried out, and the obtained patterns were analyzed using Rietveld refinements. Figure 2a shows a refinement of $Li_{5.5}PS_{4.5}Cl_{1.5}$ as an example; all refinements and the tabulated structural information can be found in the Supporting Information, Figures S1-S6, and Tables S2-S7.



Figure 2: a) Representative Rietveld refinement against neutron diffraction data of $Li_{5.5}PS_{4.5}Cl_{1.5}$. A small fraction of impurity phases corresponding to ~0.7 wt% Li_3PO_4 is present, along with reflections from the vanadium sample container. b) Stacked neutron diffraction patterns of the synthesized materials with different chloride content x of for $Li_{7-x}PS_{6-x}Cl_x$.

All samples (compiled in Figure 2b) exhibited minor impurity phases of Li₃PO₄ (no more than 2.0 wt%). The composition of x = 0, *i.e.* Li₇PS₆ was analyzed using Rietveld refinements against X-ray diffraction data only. The known orthorhombic phase (shown in Figure S7) can be observed. By increasing the halide content to x = 0.25, the orthorhombic structure changes to the cubic polymorph. By further increasing the halide content, the amount of LiCl impurity increases, and

in x = 1.75 compositions around 10 wt% LiCl is present (together with ~ 3 wt% Li₃PS₄), indicating that a solubility limit has been reached. This is in agreement with the findings of Nazar and coworkers.³³ For the following analyses only the cubic polymorphs that have not reached the solubility limit yet, i.e. Li_{7-x}PS_{6-x}Cl_x (0.25 $\leq x \leq 1.50$), were are investigated using neutron diffraction.



Figure 3: a) A comparison of nominal contents of chloride with refined chloride contents in $Li_{7-x}PS_{6-x}Cl_x$. Idealized ratio given as dashed line. b) The lattice parameter decreases as the chloride content increases. c) Top panel shows the percentage of total Cl content distributed across each of the two Wyckoff sites, and bottom panel shows the Cl occupancy of the sites as it increases with increasing Cl content. d) R_{mean} , representing the radial distance between the centers

of the cage (Wyckoff 4d) and the mean Li density of the surrounding cage, increases, and the T2 - T2 distance (inter-cage jump) significantly decreases with increasing Cl content.

The nominal contents (x_N) of chloride used for the synthesis $(Li_{7-x}PS_{6-x}Cl_x)$ and refined contents (x_R) of chloride extracted from the refinement of neutron diffraction data are shown in Figure 3a, indicating that not all of the Cl^- available is taken up at higher x. This is consistent with the observed solubility limit at higher Cl⁻ fractions. The increasing chloride content causes the lattice parameter to decrease (Figure 3b); the kink in the linearity of the lattice parameter decrease is reflected in the deviation in the solubility. The changing lattice volume might not be expected given the similar ionic radii between Cl^{-} (167 pm) and S^{2-} (170 pm). One possibility is that the Li⁺ vacancies caused by the replacement of divalent sulfide with monovalent chloride may play a role in the decrease of the lattice parameter. A second possibility may be the site-disorder of X^{-}/S^{2-} that has also been shown previously to have an effect on lattice parameter, even at fixed composition.^{19,29,33} Nevertheless, chloride added into the structure (increasing x) distributes across the Wyckoff 4a and 4d sites (Figure 3c, bottom panel). A site-disorder in terms of simple occupancy of 4d cannot be established with changing Cl^{-} content, so instead, of the total amount of Cl⁻ from the refined x_R values, the percentage of this amount that sits on each respective Wyckoff site is shown in Figure 3c, top panel. This gives an idea of how the distribution across the two sites changes from a disorder at the nominal composition Li₆PS₅Cl of 62.3(0.7) %. It is clear that it is not constant and that preference seems to change from 4a to 4d, as more Cl⁻ is added. The lithium occupancy in different lithium positions can be found in Figure S8. Due to the charge difference of Cl⁻ and S²⁻, the average charge on the Wyckoff 4*a* and 4*d* sites will change based on its occupancy. As a direct result of the changing local charge environment, the average surrounding Li⁺ substructure is expected to change. Li⁺ positions can be directly extracted from the neutron diffraction results. In order to better conceptualize how the substructure changes with chloride content, we have applied the same analysis from Minafra et al. reported on the lithium argyrodite sublattice of Li₆PS₅X (X = Cl, Br, I)²⁶ and recent work on Li₆PS₅Br with changing S^{2-}/Br^{-} site disorder.²⁹ From the refined crystallographic data, the Li⁺ distribution is determined radially from the 4d site (the center of the Li^+ "cage").²⁶ The radius at which the mean Li^+ density is located is denoted as R_{mean} . A higher value of R_{mean} describes an expansion of the Li⁺ "cage" away from its center, and to a point, a higher level of interconnectivity of the cages.²⁹ The effect of increasing the chloride content x is to increase R_{mean} , from a value of 2.475 Å in $x_{\text{R}} = 0.235$ to 2.525 Å in $x_R = 1.40$. The T2 – T2 distance represents the shortest path between Li⁺ positions of separate lithium cages; this has been invoked in the past as a critical point for improving transport, as the cages must be interconnected for three-dimensional Li⁺ movement in the bulk.^{26,29} The change in T2 – T2 distance is shown in Figure 3d; with increasing chloride content, the distance decreases which is expected to lead to an higher bulk Li⁺ mobility. However, as discussed in our previous work,^{26,29} we believe that R_{mean} is a better general descriptor of the phenomenon. For the following sections, transport is described in terms of this parameter.

Ionic transport properties. Temperature-dependent electrochemical impedance spectroscopy was performed to determine the ionic conductivity of the chloride enriched lithium argyrodites $Li_{7-x}PS_{6-x}Cl_x$ up to x = 1.50. Exemplarily shown for the Nyquist plots of the nominal composition $x_N = 1.00$, 1.25, and 1.50 at -30 °C in Figure 4a, the data can be fit with an equivalent circuit model of a resistor element in parallel with a constant phase element (CPE), in series with another constant phase element corresponding to the materials resistance and electrode blocking behavior, respectively. All remaining temperature - dependent impedance data can be found in the Figure S9.



Figure 4: a) Representative fit of the impedance spectra of $Li_{7-x}PS_{6-x}Cl_x$ as a function of different nominal chloride compositions; the impedance is shown normalized by the pellet thickness for comparison. b) Representative Arrhenius data of the solid solutions.

The impedance spectra at lower temperatures show suppressed semicircles and a capacitive blocking behavior of the electrodes. The obtained capacitances of the high frequency process are $2.2 - 9.0 \cdot 10^{-10}$ F with ideality factors of $\alpha > 0.85$. At higher temperature, the contribution of the

resistor/CPE shifts to frequencies that are too high to resolve within the probed frequency range and only the tail of the blocking electrode behavior was used for impedance fitting as shown in Figure S9. Clearly, the bulk and grain boundary contributions cannot be deconvoluted and all values represent total ionic conductivities. The temperature dependent ionic conductivities for all compositions are shown in Figure 4b, showing linear Arrhenius behavior as expected.



Figure 5: a) Room-temperature ionic conductivity and b) activation energies of $Li_{7-x}PS_{6-x}Cl_x$ as a function of the refined chloride composition. c) Room-temperature ionic conductivity and (d) activation energy as a function of R_{mean} .

The room temperature ionic conductivities and corresponding activation energies are shown in Figures 5 as a function of the refined chloride content. The ionic conductivity increases from 0.55 $mScm^{-1}$ at $x_R = 0.235$ to 7.2 $mScm^{-1}$ at $x_R = 1.40$, along with a concomitant decrease in activation energy from 0.48 to 0.39 eV, respectively. These obtained values are similar to the recent work on

the substitution series $Li_{6-x}PS_{5-x}Cl_{1+x}$, which achieved a conductivity of over 9 mScm⁻¹ and activation energy of 0.29 eV at x = 0.5.³³ Small differences in conductivities between these same compositions likely arise from measurement practices across working groups.¹⁶ Figure 5c shows the ionic conductivity as a function of R_{mean} . The conductivity increases with R_{mean} , from a value of 0.55 mS cm⁻¹ (2.475 Å) to 7.2 mScm⁻¹ (2.525 Å). This is expected as higher R_{mean} means a higher level of interconnectivity of the cages, easing Li movement through the structure.^{26,29} Figure 5d show that the activation energy decreases as a function of R_{mean} . Interestingly, the increase in conductivity is not linear, instead exhibiting a drastic increase after x = 1.00. This behavior is not reflected as much in the activation energy. Clearly something significant happens at Cl⁻ content higher than Li₆PS₅Cl (x = 1.00); whether this has to do with having more Cl⁻ than a single 4a/4dsite could hold, or reaching a critical lower concentration of Li⁺ per unit cell, or a different cause, is unclear at this time.



Figure 6: a): The Li-ion migration pathways were analyzed using the bond valence sum method with additional chloride content, showing T2 - T2 jump distance which represent the interconnected cages. b) Migration barrier along the T5 - T2 - T2 - T5 - T5 pathway that

corresponds to transport from one cage to the next. c) The experimental activation energy, calculated T2 - T2 migration barrier that can be observed with additional chloride content.

In order to gain insight into the energetics of lithium diffusion pathways, we performed bond valence sum analyses. The Li⁺ positions and structural framework are determined from the neutron diffraction data. The experimental activation energy is calculated using impedance spectroscopy and obtained T2 – T2 activation energies are extracted from the bond valence sum. Figure 6a shows a visual representation of cages along with Li-ion migration pathways. The Li⁺ pathway is chosen along the T2 and T5 sites: T5 - T2 - T2 - T5 - T5 (Figure 6 b). As the Cl⁻ content increases in Li_{7-x}PS_{6-x}Cl_x, the activation energy for a T2 – T2 jump decreases and for a T2 – T5 jump increases. However, the T2 – T2 jump shows the highest energy barriers, suggesting to be the rate limiting step. These T2 – T2 activation energies are shown together with the experimental activation energies in Figure 6c; the magnitude of the experimental results is higher than predicted, but a similar trend with increasing Cl⁻ content is found. The more interconnected the Li⁺ cages are, corroborated by an expanding R_{mean} and a decreasing T2 – T2 distance, the lower the activation barriers and better ionic transport properties.

3. Conclusion

We have synthesized the chloride enriched lithium argyrodites $Li_{7-x}PS_{6-x}Cl_x$ with compositions from $x_N = 0.0$ to $x_N = 1.75$. Neutron powder diffraction was utilized to study the Li⁺ substructure and its effects on the ionic transport. With increasing Cl⁻ content, the Li⁺ "cages" expand, indicated in both decreases of T2 – T2 site distances and, in a more general sense, increases in the radial distribution of Li⁺ away from the center of the cage, characterized by R_{mean} . The changing average anionic charge distribution by Cl⁻ addition has a strong effect on R_{mean} and ultimately the conductivity as it leads to a much better connected Li⁺ landscape. This work offers a common thread to previous studies showing that, while disorder and composition play a role in increasing conductivity, one of the underlying concepts is the importance of the anionic charge distribution and how much the Li⁺ substructure is interconnected. Fully understanding how the substructure is impacted by chemical changes is a key to tailoring transport properties, and this general methodology should be applicable to many other superionic conductors.

Supporting Information

All refined neutron powder diffraction data of the chloride enriched lithium argyrodites $Li_{7-x}PS_{6-x}Cl_x$ series are reported here, as well as additional information on the constraints applied during the refinements and the refinement results. Further shown are the X-ray diffraction data of Li_7PS_6 All impedance spectra are provided. Crystallographic information files (CIF) are also provided.

Acknowledgments

The research was supported by the Federal Ministry of Education and Research (BMBF) within the project FESTBATT under grant number 03XP0177A. The D2B high-resolution diffractometer and monochromatic neutron source was used for measurement of neutron diffraction, operated by the Institut Laue-Langevin, Grenoble, France.

Conflict of Interest

The authors declare no competing financial interests.

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