

Simulation Meets Experiment: Unraveling the Properties of Water in Metal-Organic Frameworks Through Vibrational Spectroscopy

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Abstract

In nanoporous materials, guest–host interactions affect the properties and function of both adsorbent and adsorbate molecules. Due to their structural and chemical diversity, metal-organic frameworks (MOFs), a common class of nanoporous materials, have been shown to be able to efficiently and, often, selectively adsorb various types of guest molecules. In this study, we characterize the structure and dynamics of water confined in ZIF-90. Through the integration of experimental and computational infrared (IR) spectroscopy, we probe the structure of heavy water (D₂O) adsorbed in the pores, disentangling the fundamental framework–water and water–water interactions. The

experimental IR spectrum of D₂O in ZIF-90 displays a blue-shifted OD-stretch band compared to liquid D₂O. The analysis of the IR spectra simulated at both classical and quantum levels indicates that the D₂O molecules preferentially interact with the carbonyl groups of the framework and highlights the importance of including nuclear quantum effects and taking into account Fermi resonances for a correct interpretation of the OD-stretch band in terms of the underlying hydrogen-bonding motifs. Through a systematic comparison with the experimental spectra, we demonstrate that computational spectroscopy can be used to gain quantitative, molecular-level insights into framework–water interactions that determine the water adsorption capacity of MOFs as well as the spatial arrangements of the water molecules inside the MOF pores which, in turn, are key to the design of MOF-based materials for water harvesting.

Introduction

Hydrogen bonding (H-bonding) between water and host materials has proven to be important for the structure and function of the latter.^{1–3} Strong H-bonds can lead to the formation of extended networks that modulate fundamental processes, including hydration processes,⁴ chemical synthesis and reactions,^{5,6} heat dissipation,^{7,8} and macroscopic structural formations.^{7,9–11} Under confinement, H-bonded networks are disrupted due to physical constraints and host–water interactions.^{7,12} A variety of porous materials, such as carbon nanotubes,^{13–17} zeolites,¹⁸ silica pores and channels,^{12,19} and other materials^{20,21} have been used as model systems to determine the physical properties of water in confinement. These model systems highlight the fact that the properties of the confining environment, whether hydrophobic or hydrophilic, influence H-bonding with the host.

Metal-organic frameworks (MOFs) are an attractive class of porous materials. Comprised of inorganic subunits, also known as secondary building units (SBUs), and organic linkers, MOFs assemble in crystalline three-dimensional structures with large surface areas and high porosity.²² Because of their chemical diversity and tunability, MOFs have found a wide

range of applications including, but not limited to, liquid and gas adsorption,^{23–26} chemical separation,^{27,28} catalysis,²⁹ chemical degradation,³⁰ and proton transport,³¹ as well as in electronic devices³² and magnets.³³ It has recently been shown that some MOFs display remarkable water capture capabilities under ambient conditions.^{34–38} To further capitalize on the potential of MOFs as water harvesting materials, it is necessary to gain fundamental insights on the framework–water and water–water interactions in the confining pores, as these interactions determine both the onset and overall capacity of water adsorption.

An accurate description of the water H-bonding network in MOFs, which can allow for unambiguously disentangling framework–water and water–water interactions, is critical for understanding the physical mechanisms governing water adsorption in MOFs as a function of pore size and shape as well as the physicochemical properties of the framework. In this regard, an integrated approach that combines spectroscopic measurements and molecular simulations can provide such a level of detail. Linear infrared (IR) spectroscopy is sensitive to variations in the strength of the water H-bonding network,³⁹ although all measurements are performed at the ensemble-averaged level, and the decoding of structure from the spectra is often challenging. On the other hand, molecular dynamics (MD) simulations provides a means to directly connect different spectroscopic features to specific H-bonding motifs. However, the calculated spectra, and, therefore, the reliability of the comparisons with the experimental measurements depend sensitively on the ability of the molecular models to correctly describe the underlying molecular interactions.

Diffuse reflection infrared Fourier transform spectroscopy (DRIFTS) can provide functional-group specific information regarding the interaction between water and the framework. A study carried out on MIL-53(Al), which exhibits the so-called “breathing effect”, highlighted the complex mechanisms that modulate water adsorption in this MOF.⁴⁰ It was determined that water interacts tightly with multiple sites of the framework at lower hydration levels, while at higher hydration levels the spectral signatures associated with the water OH stretching vibrations are localized to fewer sites of the framework. A study on the $\text{Co}_2\text{Cl}_2\text{BTDD}$

MOF, a hydrophilic MOF that captures 82% water by weight below 30% relative humidity (RH),³⁶ used DRIFTS to elucidate the adsorption mechanism as a function of RH.⁴¹ Results from Ref. 41 provide a different picture from that derived from the analysis of the IR spectra of water in MIL-53(Al). Water was found to strongly bind to the open Co^{2+} sites of the framework at low RH and to subsequently form disconnected one-dimensional chains of H-bonded molecules bridging between the Co^{2+} sites. Upon further increase in RH, these water chains were found to nucleate pore filling, with water molecules occupying the entire pore volume before the RH reaches 30%. The different results from Refs. 40 and 41 exemplify not only the power of IR spectroscopy in providing functional group-specific information about the water adsorption process but also the variety of hydration mechanisms that can be observed in MOFs, depending on the pore sizes, shapes, and nature of the framework.

On the modeling side, one of the major challenges faced by MD simulations of MOFs is the development of force fields (FFs) that can accurately describe the underlying molecular interactions. Common FFs such as the general Amber force field (GAFF)⁴² and the universal force field (UFF)⁴³ have been used to represent the organic linkers of MOFs. These FFs have the advantage of being applicable to a wide range of systems, although they suffer from a lack of transferability in describing the coordination and geometry of metal centers in MOFs.^{44,45} More specific FFs have been developed to address this deficiency, such as MOF-FF⁴⁴ and QuickFF,⁴⁶ extensions of UFF for MOFs (UFF4MOF),^{45,47} and a zeolitic imidazolate framework force field (ZIF-FF).⁴⁸ These FFs have been extensively used to simulate adsorption isotherms as well as structural, thermodynamic, and dynamical properties of various molecules (e.g., methane and small hydrocarbons, carbon dioxide, and water) adsorbed in the MOF pores. Recently, more sophisticated polarizable models have also been used to model guest–framework interactions.^{49–54} Among different guest molecules, modeling water in MOFs presents particular challenges due to the complex nature of the water–water interactions, which is responsible for the anomalous behavior of water as a function of temperature and pressure.^{55–57}

In this study, we integrate experimental and computational IR spectroscopy with MD simulations to investigate the properties of water in ZIF-90. Compared to most ZIFs that are highly hydrophobic (e.g., ZIF-8),^{58–60} ZIF-90 gives rise to a sharp step in the water adsorption isotherm between 30% and 40% RH (Fig. 1A) due to the presence of the imidazolate-2-carboxyaldehyde linkers, which make the framework hydrophilic.⁶¹ While the overarching goal of this study is to gain general insights into the interplay between framework–water and water–water interactions, direct comparisons between experimental and simulated IR spectra also provide the unique opportunity to assess both merits and shortcomings of current models for MOF simulations and computational IR spectroscopy as well as to identify possible areas of improvement, which are critical for *in silico* screening of MOFs for water harvesting.

Methods

Material Synthesis and Characterization

ZIF-90 was synthesized by fully dissolving imidazole-2-carboxyaldehyde (ICA) in N,N-dimethylformamide (DMF) with heat as described in Ref. 62. The ICA/DMF solution was cooled to room temperature, and pyridine was added. A solution of zinc nitrate hexahydrate in methanol was poured into the ICA/pyridine/DMF solution rapidly and mixed overnight. ZIF-90 crystals were collected through 4 cycles of centrifugation with neat DMF. The collected product was then washed with DMF and dried. To confirm the synthesis of ZIF-90, the product was analyzed using Powder X-Ray Diffraction (PXRD) and Scanning Electron Microscopy (SEM). PXRD data match well with theoretically simulated spectra, and obtained SEM images visually confirm the crystal structure of ZIF-90 (Figs. S1 and S2, respectively).

Spectroscopic Measurements

DRIFTS measurements were carried out on ZIF-90 crystal powder mixed with dry KBr at 5% (w/w) using a Thermo Fischer Nicolet S10 FTIR spectrophotometer fitted with a PIKE

Technologies DiffusIR accessory. The enclosed internal space of the DiffusIR accessory was purged with varying degrees of D₂O humidified air using a humidity generator (Fig. S3). Briefly, dry air streams were mixed with D₂O humidified air until the desired humidity was generated. The enclosed internal space of the accessory was then allowed to equilibrate for 10 minutes before each spectrum and background were recorded. A background spectrum of solid potassium bromide (KBr) was taken for each RH value to correct for signal scatter using a Kubelka-Munk transform.

Molecular Models and Simulations

ZIF-90 was modeled using a flexible force field. All bonded parameters involving the zinc atom were taken from ZIF-FF,⁴⁸ while all bonded parameters involving the atoms of the ligand were taken from the General Amber Force Field (GAFF).⁴² The MOF structure was optimized using periodic density functional theory (DFT) calculations carried out with the Vienna Ab initio Simulation Package (VASP),^{63–66} using the PBE exchange-correlation functional⁶⁷ with Grimme’s D3 dispersion correction⁶⁸ in combination with a projector-augmented wave (PAW) treatment^{69,70} with a 700 eV kinetic energy cutoff. A $2 \times 2 \times 2$ k -point grid was used, and forces were converged to a tolerance of 0.03 eV/Å. Atomic point charges were calculated for the optimized ZIF-90 structure using the density derived electrostatic and chemical (DDEC) charges.⁷¹ Lennard-Jones (LJ) van der Waals (VDW) parameters were taken from ZIF-FF with an initial guess used for the oxygen atom not present in ZIF-FF. Atom types as well as all bonded and non-bonded parameters for the MOF force field are given in Fig. S4 and Tables S1-S4.

Water was modeled using the MB-pol potential energy function (PEF).^{72–74} MB-pol has been shown to correctly predict the properties of water from the gas to the condensed phase,^{75,76} which makes it particularly well-suited for modeling the properties of water adsorbed in MOFs as a function of RH. Furthermore, MB-pol enables accurate simulations of vibrational spectra of water clusters,^{72,77,78} liquid water,^{79–82} the air/water interface,^{83,84}

and ice,^{85–87} and is thus an ideal model to monitor the evolution of the DRIFTS spectra of water adsorbed in MOFs as a function of RH.⁴¹

As in Ref. 41, the framework–water non-bonded interactions were represented in terms of permanent electrostatics and Lennard-Jones (LJ) interactions. Since the extension of our many-body models to generic molecules compatible with MB-pol is under development,^{88,89} in this study the LJ parameters between the MOF atoms and water were obtained by applying the Lorentz-Berthelot mixing rules using the LJ parameters of the TIP4P/2005 water model,⁹⁰ which is the closest point-charge model to MB-pol.⁹¹

All MD simulations were carried out in periodic boundary conditions for systems consisting of $2 \times 2 \times 2$ primitive cells of ZIF-90 and various water loadings, ranging from 25% to 70% RH, using in-house software based on the DL_POLY_2 simulation package.⁹² Constant pressure and constant temperature (NPT) simulations at 1.0 atm and 300 K were performed for 500 ps at each water loading. Since the average box dimensions do not deviate significantly from the experimental values (see Table S5 in the Supporting Information), the latter were used in all simulations to guarantee direct comparisons with the experimental measurements. Each system was equilibrated in the canonical (constant number of atoms, volume, and temperature, NVT) ensemble at 300 K for 1 ns. All dynamical properties, including the IR spectra, were calculated in the microcanonical (constant number of atoms, volume, and energy, NVE) ensemble at 300 K by averaging over 20, 50 ps long independent trajectories. The equations of motion were propagated according to the velocity-Verlet algorithm with a time step of 0.2 fs, and the temperature was maintained at 300 K by a Nosé-Hoover chain of four thermostats. Short-range interactions were truncated at an atom–atom distance of 15.0 Å, and long-range electrostatics were calculated using the Ewald sum.⁹³

To assess the importance of nuclear quantum effects in determining the structural and dynamical properties of water adsorbed in ZIF-90, centroid molecular dynamics (CMD) simulations^{94–98} were also performed. In the CMD simulations, each atom was represented by a Feynman’s ring polymer discretized with 32 beads, and the centroid variables were propa-

gated using the partially adiabatic separation scheme of Refs. 97 and 99, with an adiabaticity parameter of $\gamma=0.25$ and a time step of 0.05 fs. Given the associated computational cost, all CMD simulations were only performed for systems consisting of $1\times 1\times 1$ primitive cells of ZIF-90 at 40% RH.

The theoretical IR spectra were calculated within the time-dependent formalism according to

$$I_{IR} = \left[\frac{2\omega}{3V\hbar c\epsilon_0} \right] \tanh(\beta\hbar\omega) \int_{-\infty}^{\infty} e^{i\omega t} \langle \mu(0)\mu(t) \rangle dt \quad (1)$$

where V is the system volume, c is the speed of light in vacuum, ϵ_0 is the permittivity of free space, and $\beta = (k_B T)^{-1}$, with k_B being Boltzmann’s constant. In Eq. 1, $\langle \mu(0)\mu(t) \rangle$ is the ensemble-averaged dipole-dipole time correlation function that was calculated by averaging over 20 NVE trajectories at each water loading, with μ being represented by the MB- μ many-body dipole moment function.⁷⁹

Due to the neglect of nuclear quantum effects, classical MD simulations predict vibrational spectra that are systematically blue-shifted relative to the corresponding experimental spectra.^{79,83,100,101} Therefore, in all the analyses presented in the following sections, the theoretical spectra calculated from classical MD simulations are red-shifted by 137 cm^{-1} in the stretching region, as shown in Fig. S5 in the Supporting Information, to facilitate comparisons with the experimental DRIFTS spectra. Furthermore, MB-pol has been shown to slightly underestimate the strength of H-bonds in liquid water, which results in a blue-shift of 57 cm^{-1} in the OH stretching region of the quantum CMD spectra.⁷⁹ In the following analyses, the quantum CMD spectra are thus red-shifted by 57 cm^{-1} (Fig. S5) to facilitate comparisons with the experimental DRIFTS spectra.

Finally, Fermi resonances between the overtones of the bending vibrations and stretching vibrations of the water molecules in the liquid phase have been shown to be important for a quantitative representation of the OH-stretch lineshape in both water clusters^{77,102} and

liquid water.⁸² To calculate the Fermi resonance contributions to the OD-stretch lineshape, local mode/local monomer (LM) calculations^{103–105} were performed on clusters comprised of a central D₂O molecule surrounded by 16 H₂O molecules, which were used to model the intramolecular (bending and stretching) modes of D₂O in the liquid phase. These clusters were the same clusters used in Ref. 82, which were extracted from CMD simulations of a single H₂O in a D₂O solution. The two hydrogen atoms on the H₂O molecule were optimized, while every other atom was held fixed, using the ORCA package¹⁰⁶ interfaced with our in-house MB-pol software. After this optimization, the central H₂O molecule was converted to D₂O while the 16 surrounding D₂O molecules in the cluster were converted to H₂O for the LM calculations.

Results and Discussion

ZIF-90 exhibits a type V adsorption isotherm (Fig. 1A),⁶⁰ adsorbing water in one pore-filling step that begins at $\sim 30\%$ RH and ends at $\sim 40\%$ RH. To facilitate the analysis of the IR spectra, all measurements and simulations were carried out with D₂O since the OD stretching vibrations do not overlap with any of the framework vibrations. The experimental IR spectra of D₂O in ZIF-90 in the OD stretching region, obtained after subtracting the

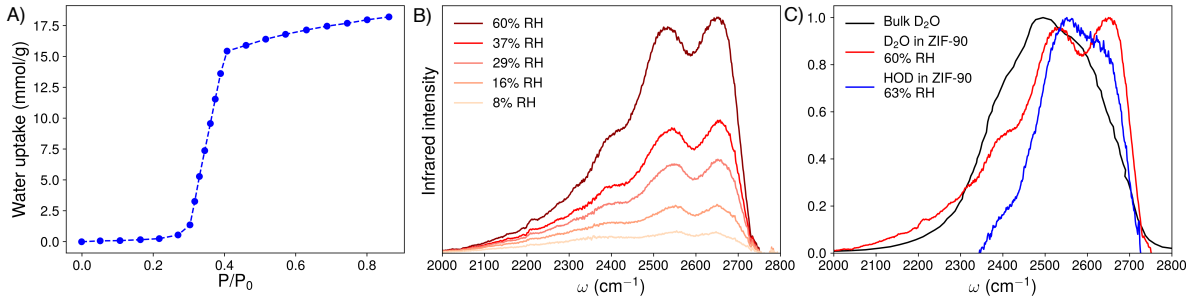


Figure 1: A) Water adsorption isotherm for ZIF-90 from Ref. 60. B) Experimental IR spectra of D₂O in ZIF-90 at 8% RH, 16% RH, 29% RH, 37% RH, and 60% RH. C) IR spectra of experimental bulk D₂O (black) from Ref. 107, D₂O in ZIF-90 at 60% RH (red), and HOD in ZIF-90 at 63% RH (blue). The spectra are normalized to a value of one.

spectrum measured at 0% RH from the spectra measured at different RH values, are shown in Fig. 1B. As the humidity increases to $\sim 10\%$, two peaks centered at ~ 2550 and ~ 2650 cm^{-1} begin to develop, with their intensity monotonically increasing as the RH increases. Importantly, Fig. 1B also shows that the spectral weight shifts from the lower to the higher frequency portion of the lineshape as the RH increases. The comparison with the IR spectrum of bulk D_2O (Fig. 1C) indicates that ZIF-90 induces a blue shift of ~ 40 cm^{-1} to the OD lineshape of the adsorbed D_2O molecules. The OD lineshape in ZIF-90 displays appreciably higher intensity on the blue side of the maximum at 2650 cm^{-1} . Importantly, DRIFTS measurements carried out for HOD in ZIF-90 (Fig. 1C) provide a significantly narrower OD-stretch lineshape that lacks the two distinct peaks that characterize the analogous OD-stretch of D_2O in ZIF-90. Furthermore, the absence of intramolecular coupling in the HOD molecules results in a decrease of the IR intensity on the red side of the spectrum around ~ 2400 cm^{-1} . The comparison of the IR spectra of D_2O and HOD in ZIF-90 thus allows us to unambiguously assign the shoulder at ~ 2400 cm^{-1} in the D_2O spectrum to the Fermi resonances and not to particular ice-like structures of D_2O in the ZIF-90 pores.

To gain insights into the adsorption process, IR spectra were calculated from both MD and CMD simulations carried out at different D_2O loadings. Fig. 2 shows a comparison between the experimental, MD, and CMD OD-stretch lineshapes at 300 K. As discussed in Ref. 82, due to numerical limitations associated with using Cartesian coordinates to propagate the dynamical trajectories, the lineshape calculated using quantum CMD simulations is unable to quantitatively describe Fermi resonances.^{108,109} On the other hand, the Fermi resonances do not appear in the lineshape calculated using classical MD simulations because, in this case, the overtones of the bending vibrations do not overlap with the OD stretching vibrations. Fig. 2A shows the experimental and simulated CMD IR spectrum of bulk D_2O . The CMD spectrum is able to reproduce the experimental D_2O spectrum fairly well, only lacking intensity around ~ 2400 cm^{-1} , which is due to missing intensity from the Fermi resonances.⁸² This inability to reproduce Fermi resonances leads to two distinct peaks in the

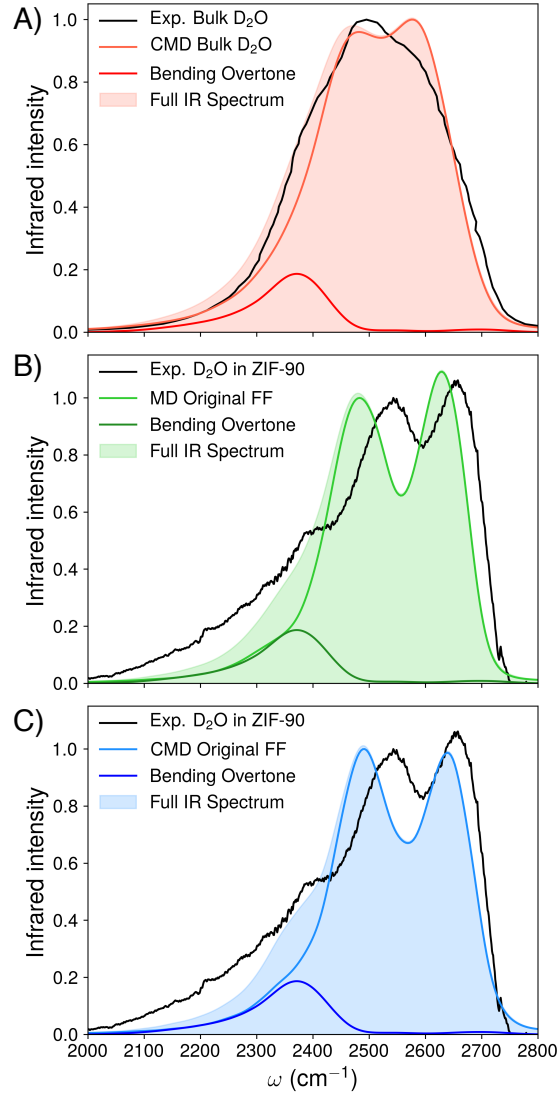


Figure 2: A) Experimental bulk D_2O IR spectrum (black) from Ref. 107, CMD bulk D_2O (pink) spectrum, Fermi resonance contribution (red), and CMD bulk D_2O spectrum with added Fermi resonance contribution (shaded pink). B) Experimental IR spectrum of water in ZIF-90 at 40% RH (black), MD 40% RH (green) spectrum, Fermi resonance contribution (dark green), and MD 40% RH spectrum with added Fermi resonance contribution (shaded green). C) Experimental IR spectrum of water in ZIF-90 at 40% RH (black), CMD 40% RH (light blue) spectrum, Fermi resonance contribution (blue), and CMD 40% RH spectrum with added Fermi resonance contribution (shaded light blue). All MD spectra are red-shifted by 137 cm^{-1} , CMD spectra are red-shifted by 57 cm^{-1} , DOD bending overtones are blue-shifted by 24 cm^{-1} , and all spectra are normalized to a value of one at the lower frequency peak.

CMD spectrum of bulk D₂O instead of the single, broad peak with two shoulders observed in the experimental spectrum (Fig. 2A). After adding the contributions from the Fermi resonances (red trace), the resulting CMD OD-stretch lineshape of bulk D₂O is in remarkably good agreement with the corresponding experimental lineshape.

Figs. 2B and 2C show comparisons between the experimental and theoretical lineshapes for D₂O in ZIF-90 at 40% RH calculated from both (classical) MD and (quantum) CMD simulations, respectively. While the simulated lineshapes correctly reproduce both the blue shift of the OD-stretch lineshape measured experimentally (Fig. S7) and the presence of two distinct peaks, both MD and CMD lineshapes overestimate the separation between the two peaks compared to experiment.

The origin of the blue shift can be explained by investigating the structure of bulk D₂O and D₂O in ZIF-90. As shown in Table 1, D₂O molecules adsorbed in MOF pores form longer, and, consequently, weaker D-bonds compared to bulk. These longer D-bonds are associated with OD stretching vibrations that absorb at relatively higher frequencies, causing the blue shift of the OD-stretch lineshape found for D₂O confined in the ZIF-90 pores.

Since it has been shown that MB-pol correctly describes the properties⁷⁶ and IR spectrum of bulk D₂O (Fig. 2A), the differences between the experimental and simulated IR spectra of D₂O in ZIF-90 are likely due to an inaccurate representation of the interactions between the D₂O molecules and the ZIF-90 framework. The present simulations (Fig. S9) as well as other studies of water in ZIF-90⁵⁹ indicate that the D₂O molecules adsorbed in the ZIF-90 pores can D-bond to the carbonyl groups of the framework. Furthermore, Figs. 2B and 2C show that, compared to experimental results, the low-frequency peak in both MD and CMD OD-stretch lineshapes is red-shifted and misses intensity around $\sim 2550\text{ cm}^{-1}$. This frequency region corresponds to OD stretching vibrations of D₂O molecules that are D-bonded to the carbonyl groups of the framework.¹¹⁰

In order to more accurately describe the strength of D-bonding between the D₂O molecules and the framework, we modified the LJ potential between the D₂O oxygen atom (OW) and

Table 1: Average OD bond, water–water (OW–OW) D-bond lengths, and water–MOF (OW–o) D-bond lengths calculated from MD and CMD simulations for bulk D₂O as well as ZIF-90 at 40% RH using both the original and modified force fields.

	Simulation	OD bond (Å)	OW–OW D-bond (Å)	OW–o D-bond (Å)
Bulk D ₂ O	CMD	0.976	1.916	
ZIF-90 40% RH (original FF)	CMD	0.973	1.937	2.146
ZIF-90 40% RH (original FF)	MD	0.969	1.943	2.147
ZIF-90 40% RH (modified FF)	CMD	0.975	1.945	1.713
ZIF-90 40% RH (modified FF)	MD	0.970	1.950	1.715

the oxygen atom (o) of the carbonyl functional groups of the framework by decreasing σ_{OW-o} by 10% and ϵ_{OW-o} by 20%, which leads to a modified FF. By doing this, the OW–o repulsion decreases, which allows the D₂O molecules to approach the framework more closely and establish stronger D-bonds with the carbonyl groups.

The effects of modifying the OW–o LJ potential can be directly seen in the differences between the radial distribution functions (RDFs) calculated from MD simulations carried out with the original and modified FFs (Fig. 3) as well as in the average water–MOF D-bond length in Table 1. In the original FF, the first peak in the OW–OW RDF is located at 2.8 Å (black), while the first peak in the OW–o RDF is at 3.1 Å (gray), suggesting that the D₂O molecules approach each other closer than D₂O molecules can approach the MOF atoms. On the other hand, the first peak in the OW–OW RDF remains at 2.8 Å (royal blue) but the first peak in the OW–o RDF moves to 2.55 Å (light blue) with the modified FF. This indicates that, while the interaction between D₂O molecules is not affected by changes in

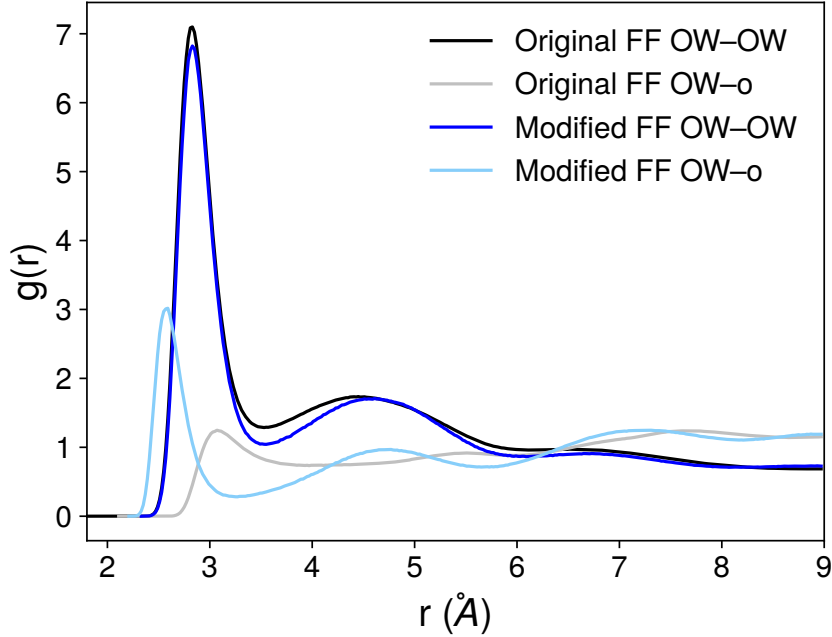


Figure 3: OW–OW (dark colors) and OW–o (light colors) RDFs of water confined in ZIF-90 in the original force field (black colors) and the modified force field (blue colors).

the framework–D₂O interactions, the modified FF leads to D₂O molecules forming shorter and stronger D-bonds with the framework than in the original FF (Table 1). It should be noted that other modifications of the original FF were also investigated and were found to significantly affect the D₂O structure in the ZIF-90 pores, leading to large deviations in both RDFs and IR spectra as shown in the Supporting Information (Figs. S10, S11, and S12).

The influence of modifying the OW–o LJ potential is clearly seen in the IR spectra calculated from MD and CMD simulations which are shown in Fig. 4A and Fig. 4B, respectively. Decreasing σ_{OW-o} by 10% and decreasing ϵ_{OW-o} by 20% in the modified FF redistributes the IR intensity between ~ 2500 - 2550 cm⁻¹. Specifically, the shorter water–MOF D-bonds in the modified FF move the IR intensity from the high-frequency range to the mid-frequency range of the OD-stretch lineshape, leaving the low-frequency peak of the strongest D-bonds between water molecules unchanged. With the modified FF, the CMD simulation closely reproduces the intensity of the two main peaks of the experimental spectrum at ~ 2550 and

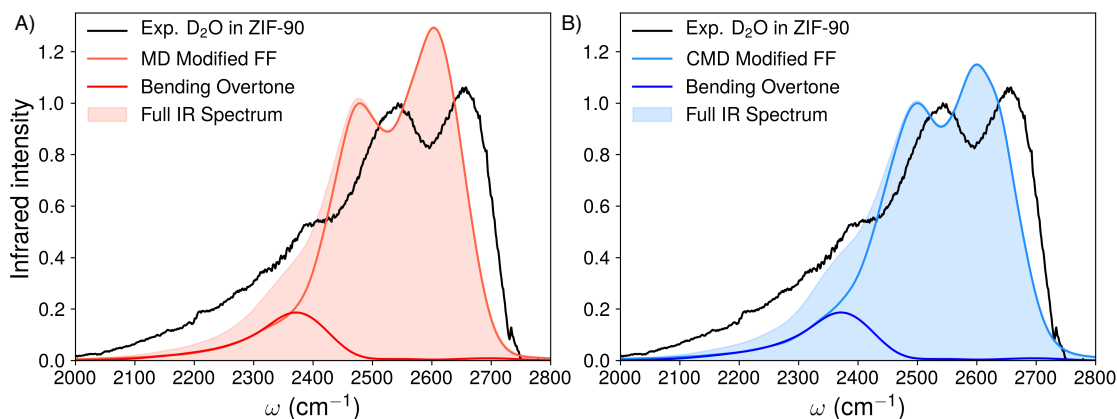


Figure 4: A) IR spectra of water in ZIF-90 at 40% RH from experiment (black), MD at 40% RH for the modified FF (pink), Fermi resonance contribution (red), and MD at 40% RH with added Fermi resonance contribution for the modified FF (shaded pink). B) IR spectra of water in ZIF-90 at 40% RH from experiment (black), CMD at 40% RH for the modified FF (light blue), Fermi resonance contribution (blue), and CMD at 40% RH with added Fermi resonance contribution for the modified FF (shaded light blue). All MD spectra are red-shifted by 137 cm^{-1} , CMD spectra are red-shifted by 57 cm^{-1} , and DOD overtones are blue-shifted by 24 cm^{-1} . All spectra are normalized to a value of one at the lower frequency peak.

$\sim 2650\text{ cm}^{-1}$. Furthermore, adding the Fermi resonance contributions increases the intensity on the red side of both the MD and CMD OD-stretch lineshapes, which makes the CMD lineshape reproduce the corresponding experimental lineshape nearly quantitatively. As seen in Table 1, simulations with the modified FF also lead to slightly longer OW-OW D-bond which results in the loss of some IR intensity from the lower-frequency peak at $\sim 2450\text{ cm}^{-1}$ (Fig. 4).

To provide further insights into the structural arrangements of the D_2O molecules in the ZIF-90 pores in connection with the OD-stretch lineshape, Fig. 5 shows the distribution of different D-bond topologies for both bulk D_2O and D_2O in ZIF-90. The majority of bulk D_2O molecules donate and accept at least one D-bond, with the varying number of donors and acceptors resulting in the broad OD-stretch band. On the other hand, the MD simulations of D_2O in ZIF-90 carried out with the original FF indicate that, among D_2O molecules not involved in D-bonding with the framework (Fig. 5A), a relatively smaller fraction donates

and accepts at least one D-bond compared to bulk D_2O . When the D_2O molecules donate one D-bond to the framework (Fig. 5B), the fraction of D_2O molecules that are involved in D-bonding with themselves is significantly smaller than in bulk D_2O . Therefore, due to the overall smaller number of D-bond donors and acceptors, the vibrations of the D_2O molecules confined in ZIF-90 retain some character of the symmetric and asymmetric stretches of gas-phase D_2O , which results in the two distinct peaks found in the IR spectra.

The comparison between the D-bonding topologies obtained with the original and modified FFs indicates that the latter leads to an overall decrease in the number of D_2O molecules that are D-bonded with themselves when the D_2O molecules are not involved in D-bonding with the framework (Fig. 5A). This is accompanied by an overall increase of D_2O molecules that donate D-bonds to the carbonyl groups of the framework, whose fraction increases from 25% in the MD simulations with the original FF to 40% in the MD simulations with the modified FF, resulting in a more tetrahedral arrangement of the D_2O molecules in the ZIF-90 pores (Fig. S13). Given the differences in strength between the stronger D_2O – D_2O and weaker D_2O –carbonyl D-bonds,¹¹⁰ this analysis demonstrates that the rearrangements of the D-bonding network predicted by the MD simulations with the modified FF are responsible for the redistribution of the IR intensity in the $\sim 2500\text{--}2550\text{ cm}^{-1}$ frequency range of the simulated OD-stretch lineshapes shown in Fig. 4.

A small red shift between the experimental and simulated OD-stretch lineshapes still remains after modifying the FF and including the Fermi resonances (Fig. 4). In this context, it should be reminded that the intrinsic shifts applied to the simulated lineshapes are based on the analysis of the IR spectrum of bulk D_2O . As highlighted in Fig. 1C, there is a $\sim 40\text{ cm}^{-1}$ blue shift in the experimental spectra of D_2O confined in ZIF-90 relative to bulk D_2O . This blue shift accounts for the small differences that remain in the higher-frequency portion of the lineshape. Although the level of agreement with the experimental OD-stretch lineshape is nearly quantitative, the simulated lineshapes display slightly lower intensities between 2000 cm^{-1} and 2400 cm^{-1} . As discussed above, this frequency range is primarily associated

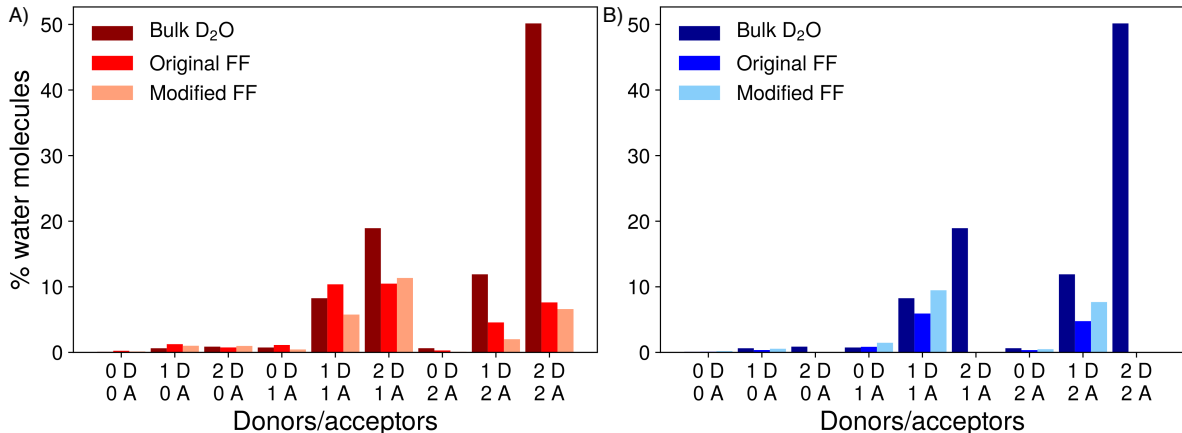


Figure 5: A) Donating (D) and accepting (A) D-bond distribution of water molecules when there are zero D-bonds donated to the MOF. D-bond distributions are shown for bulk D₂O (dark red), the original FF (red), and the modified FF (light red). D-bonds are plotted as a percent of the total number of water molecules in each simulation. B) Donating (D) and accepting (A) D-bond distribution of water molecules when there is one D-bond donated to the MOF. D-bond distributions are shown for bulk D₂O (dark blue), the original FF (blue), and the modified FF (light blue). D-bonds are plotted as a percent of the total number of water molecules in each simulation.

with Fermi resonances. While this is the first study where Fermi resonance contributions are included in simulations of the IR spectra of water in MOFs, state-of-the-art calculations currently only allow for determining these contributions from simulations of relatively small gas-phase water clusters. This approximation is likely responsible for the remaining differences between experimental and simulated lineshapes in the 2000-2400 cm⁻¹ frequency region.

Conclusions

In this study, we have characterized the structure of D₂O adsorbed in ZIF-90 as a function of relative humidity by integrating experimental and computational IR spectroscopy. Due to the confining environment provided by the framework, the D₂O molecules are found to establish longer and relatively weaker D-bonds that manifest in an OD-stretch band that is blue shifted compared to bulk. A systematic analysis of the OD-stretch band, carried out

using both (classical) MD and (quantum) CMD simulations demonstrates the importance of explicitly taking into account both Fermi resonances and nuclear quantum effects. Our simulations indicate that the D₂O molecules preferentially interact with the carbonyl groups of the framework. However, it is found that “off-the-shelf” force fields commonly used in MD simulations of MOFs are unable to correctly represent the strength of the framework–water interactions. Guided by direct comparisons with the experimental OD-stretch lineshape, we have demonstrated that it is possible to use the comparison between experimental and simulated IR spectra to guide the refinement of the force field parameterization representing the framework–water interactions. This “reverse engineering” process has led to a refined set of Lennard-Jones parameters describing the interactions between the oxygen atoms of the water molecules and the carbonyl groups of the framework, which results in nearly quantitative agreement between the experimental and simulated OD-stretch lineshapes and, in turn, allows for an accurate, molecular-level characterization of the structural arrangements of the D₂O molecules inside the ZIF-90 pores. Our results suggest that the integration of experimental and computational vibrational spectroscopy can play an important role in characterizing the molecular properties of water adsorbed in MOFs and identifying key framework–water interactions, which is critical for the design of efficient MOF-based materials for water harvesting.

Associated Content

Experimental PXRD, SEM, and humidity generator; all force field parameters for ZIF-90; additional comparisons of experimental and simulated IR spectra; structural and dynamical properties of water confined in ZIF-90.

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TOC Figure

