

Twistable Dipolar Aryl Rings as Electric Field Actuated Conformational Molecular Switches

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The ability to control the chemical conformation of a system via external stimuli is a promising route for developing molecular switches. For eventual deployment as viable sub-nanoscale components that are compatible with current electronic device technology, conformational switching should be controllable by a local electric field (i.e. E-field gateable) and accompanied by a rapid change in conduction. In organic chemical systems the degree of π -conjugation is linked to the degree of electronic delocalisation, and thus largely determines the conductivity. Here, by means of accurate first principles calculations, we study the prototypical biphenyl based molecular system in which the dihedral angle between the two rings determines the degree of conjugation. In order to make this a gateable system we create a net dipole by asymmetrically functionalising one ring with electronegative substituents (F, Br and CN) with different polarisabilities. In this way, the application of an E-field interacts with the dipolar system to influence the dihedral angle, thus controlling the conjugation. For all three considered substituents we consider a range of E-fields, and in each case extract conformational energy profiles. Using this data we obtain the minimum E-field required to induce a barrierless switching event for each system. We further extract the estimated switching speeds, the conformational probabilities at finite temperatures, and the effect of applied E-field on electronic structure. Consideration of these data allow us to assess which factors are most important in the design of efficient gateable electrical molecular switches.

Introduction

Use of external stimuli (e.g. light, temperature, pH) to controllably induce molecular level structural and/or chemical changes can be exploited for using individual molecules like switches in electronic devices.^{1,2,3} Unlike the top-down fabrication of inorganic silicon-based transistors, molecular switches are synthesised via bottom-up chemical design. As such, there are many different examples of switchable molecular systems based on a wide range of chemistries.⁴ Since electric field (E-field) driven semiconductor-based switches are key components in electrical circuitry, in recent years extensive research efforts have been devoted to fabricating junctions of single-molecules or self-assembled monolayers whose conductance can be reversibly switched by means of external E-fields. Large and reversible jumps of conductance triggered by E-fields can be achieved by means of changes in the redox state of electrochemically active molecules^{5,6,7}, molecular conformational rearrangements^{8,9,10,11,12,13}, and structural changes in the electrode-molecule contacts.^{14,15} Electrically controlled switching is also possible through E-field-induced molecular tautomerizations, isomerizations, conformational changes or rotations in molecular crystals¹⁶, in solution¹⁷, in frozen molecular solids¹⁸, in molecules adsorbed on surfaces^{19,20,21,22,23,24} and in encapsulated molecules.²⁵ It is noted that external E-fields can even catalyze bond rupture processes²⁶ or more complex reactions at the single-molecule level^{27,28,29}.

Although molecular switches are an attractive route towards extreme miniaturization, to eventually compete as viable transistors they must be: i) highly responsive to modest E-fields, and ii) able to control the transport of charge carriers. Herein, we focus in one of the most promising molecular building blocks for single-molecule electronics, namely, biphenyl (see Fig. 1a). In an elegant single-molecule junction study, a set of different substituted biphenyl molecules with variable degrees of steric hindrance between the two aryl rings allowed for the measurement of the transport properties of the biphenyl moiety as a function of the relative twist angle between the two rings.³⁰ It was found that conformational planarity favours through-bond π -conjugation and higher conductance, whereas relatively twisted conformations are highly resistive (see Fig. 1b).^{30,31} More recent experiments have shown that thermal fluctuations of the torsion angle of phenyl groups in junctions containing biphenyl moieties result in stochastic conductance switching³² while an E-field-induced variation of the torsion angle results in a controlled conductance switching.^{11,13} Specific values of the torsion angle together with appropriate electron-donating and withdrawing substituent groups in biphenyl-containing junctions have been predicted to lead to molecular rectifiers.³³ Besides the single-molecule electronics context, the torsion angle between aryl rings in biphenyl moieties is also a key factor to bear in mind for optimal performance of organic functional materials.^{34,35,36,37} In fact,

changes of this torsion angle in response to external stimuli is a powerful tool to develop switchable materials.^{38,39}

Here, we asymmetrically functionalise one of the two rings of biphenyl (i.e. X in Fig. 1) to make a dipolar system that intrinsically is conformationally responsive to applied E-fields. We note that twistable dipolar aryl rings have been proposed as a possible basis for E-field actuated molecular switching in hypothetical 2D graphene–graphyne hybrids.³⁹ Other modelling studies have shown that mono-substituted benzene rings also hold potential for E-field actuated molecular rectifiers⁴⁰ or switches⁴¹ at the single-molecule level. Herein, we use accurate computational modelling to study how to optimise twistable dipolar aryl rings for deployment as molecular switches by comparing three different dipolar functionalisations (X = F, Br, CN). Specifically, we consider the energetic contributions due to molecular conformational and permanent/induced dipoles for a range of twist angles (θ) and applied E-field strengths (F). We also consider switching speed and the effectiveness of the switching at a range of temperatures. Thereby, we provide an overview of which factors are most important for designing low field E-field actuated molecular conformational switches at finite temperatures.

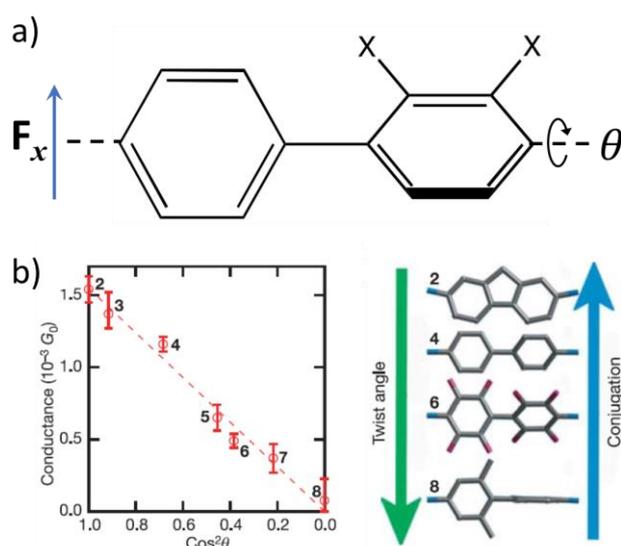


Figure 1. a) General chemical structure of the model biphenyl systems used herein, where F_x is the applied E-field vector in the x -direction, θ is the conformational twist angle, and the X are the functional dipole-inducing substituents (X = F, Br, CN). The left ring is orientationally fixed in the x - y plane and the right ring is geometrically unconstrained. b) Relationship between conformational twist angle (θ) and conductance in biphenyl ring systems (figure adapted with permission from ref 30),

Methodology

The energies and dipoles of each functionalized bi-phenyl system (i.e. with X = F, Br, CN) were computed using density functional theory (DFT) based calculations employing the hybrid PBE0 functional and a 6-31+G(d) basis set.^{42,43} using the Gaussian 09 code.⁴⁴ In each case we consider

conformations resulting from varying θ from 0° to 180° in steps of 5° , where the angle of the non-functionalised aryl ring is kept fixed at 0° . Constrained optimisations for each conformation were performed with zero applied field, and an E-field applied in the x direction (\mathbf{F}_x) which is taken to be parallel to the fixed aryl ring and perpendicular to the inter-ring bond (see Fig. 1b). We consider \mathbf{F}_x values from 0.001 - 0.015 a.u. (i.e. $5.1 \times 10^6 - 7.7 \times 10^7$ V/cm) in steps of 0.002 a.u, and a relatively strong field of 0.03 a.u (1.5×10^8 V/cm). In the solid state, high powered silicon-based metal oxide semiconductor field effect transistors (MOSFETs), also typically operate using E-field strengths towards the lower bound of this range ($\sim 1 \times 10^6$ V/cm⁴⁵). However, in (bio)molecular systems intrinsic local E-fields arising from ionic charge distributions (e.g. solvents and enzyme active sites) have typical strengths in the range $1 \times 10^6 - 1 \times 10^8$ V/cm.^{46,47,48} Such fields have been used in specifically tailored experiments to induce conformational changes of molecules in solution or deposited on surfaces requiring specific set-ups to generate the E-field (e.g. interfacially trapped ions¹⁸ or proximity of a STM tip²⁰).

Generally, the total energy (E_{tot}) of a molecule in an E-field is given by:

$$E_{tot} = E^0 - \sum_{n=1}^{\infty} \frac{F^n}{n!} \cdot \mathbf{a}_n = E^0 - \boldsymbol{\mu} \cdot \mathbf{F} - \frac{1}{2} \boldsymbol{\alpha} \cdot \mathbf{F}^2 - \frac{1}{6} \boldsymbol{\beta} \cdot \mathbf{F}^3 - \dots \quad (1)$$

where E^0 the energy without E-field, \mathbf{F} is the E-field, and $\boldsymbol{\mu}$, $\boldsymbol{\alpha}$, $\boldsymbol{\beta}$ are the dipole moment, polarizability and the first hyperpolarizability, respectively.^{49,50} We note that in the case of a weakly polarizable dipolar molecule in a weak E-field, E_{tot} can be well approximately by only the first two terms of the expansion. Assuming this dipolar approximation, the torque (τ) on the dipole with respect to an applied E-field is given by:

$$\tau = \frac{dE_{tot}}{d\theta} \quad (2)$$

When the moment of inertia (I) is constant, we also have:

$$\frac{\tau}{I} = \ddot{\theta} \quad (3)$$

where $\ddot{\theta}$ is the angular acceleration. We extract τ from our calculated θ -dependent total energies (eq. 2) and solve the differential equation (eq. 3) to obtain θ as a function of time, i.e. $\theta(t)$. Thereby, we provide order of magnitude estimates of switching speeds. See S1 for more details.

The probability of finding different twist angle conformations, P_θ , at finite temperatures with different applied fields is also estimated using standard Boltzmann statistical mechanics using these DFT-calculated total energies, i.e.

$$P_{\theta} = \frac{e^{\frac{-E_{\theta}}{k_B T}}}{Z}, \text{ with } Z = \sum_{\theta=0}^{360} e^{\frac{-E_{\theta}}{k_B T}}$$

where the E_{θ} are total system energies for each considered value of θ , k_B is the Boltzmann constant, T is the temperature, and Z is the partition function.

Results and Discussion

Considering first the variation of E_{tot} with respect to θ for zero applied field, the trends can largely be understood through the interplay between π -conjugation and steric hindrance. The di-fluoro system shows two similar energy minima for $\theta = 45^{\circ}$, 135° , separated by a barrier of 6.7 kJ/mol at 90° and higher energy maxima at $\theta = 0^{\circ}$, 180° of >10 kJ/mol (see Fig. 2a). For the latter maxima when the two aryl rings are aligned in the same plane, although here π -conjugation is maximized, the steric hindrance between the close F and H substituents is also maximized and this strong repulsive interaction dominates the total energy. Twisting away from these maxima rapidly reduces the steric hindrance and allows for the energy lowering effect of π -conjugation to stabilize the system. However, π -conjugation is most favored when the two rings are most aligned and as the rings become more mutually twisted the π -conjugation decreases. This scenario explains the appearance of the small maximum for a perpendicular ring alignment (i.e. $\theta = 90^{\circ}$) where π -conjugation is zero. It is noted that the conformational energy profile obtained for the difluoro system at zero field is in line with the results reported in Ref. 51 for 2-fluorobiphenyl. This general scenario also leads to a very similar zero field total energy curve for the di-nitrile system (see Fig. 2e), where the barrier at $\theta = 90^{\circ}$ is found to be 7.1 kJ/mol. For the di-bromo system the same competition between π -conjugation and steric hindrance also explains the zero-field curve. However, here the relatively larger Br atoms cause the steric hindrance to dominate for all θ , resulting in a single minimum at $\theta = 90^{\circ}$, where this effect is minimized (see Fig. 2d). In all cases, at zero field an ensemble of any of these three biphenyl systems will exhibit a symmetric distribution of conformational angles with respect to $\theta = 90^{\circ}$.

The effect of applying an E-field in a direction parallel to the plane of the fixed ring (i.e. \mathbf{F}_x) on the conformational energy profiles of these bi-aryl ring systems is shown in detail for the case of the di-fluoro system in Figs. 2a-c. The first clear effect of the applied E-field is to induce an energetic stabilisation of only one minimum. In this way the E-field thus switches the system from a relatively broad ensemble of conformational possibilities to a specific narrower range of conformations. With, with increasing E-field strength the scale of the total energy variation with respect to θ rapidly dominates the variation of the system energy at zero field. The applied E-field clearly energetically favors conformations with lower values of θ where the permanent dipole of the F-functionalised ring

is aligned with the field. However, for higher values of θ where the permanent dipole is anti-aligned with the field, the total energy is less affected. In fact, apart from the highest field strength considered, the total energy of the anti-aligned di-fluoro system remains close to that of the zero-field case. This behavior can be understood by consideration of the energetic contributions of the permanent dipole (i.e. first term in equation 1) and the induced dipole (i.e. second and higher order terms in equation 1) for the di-fluoro system in Figs. 2b and 2c respectively. Note that these data are also provided for the di-bromo and di-nitrile systems in S2.

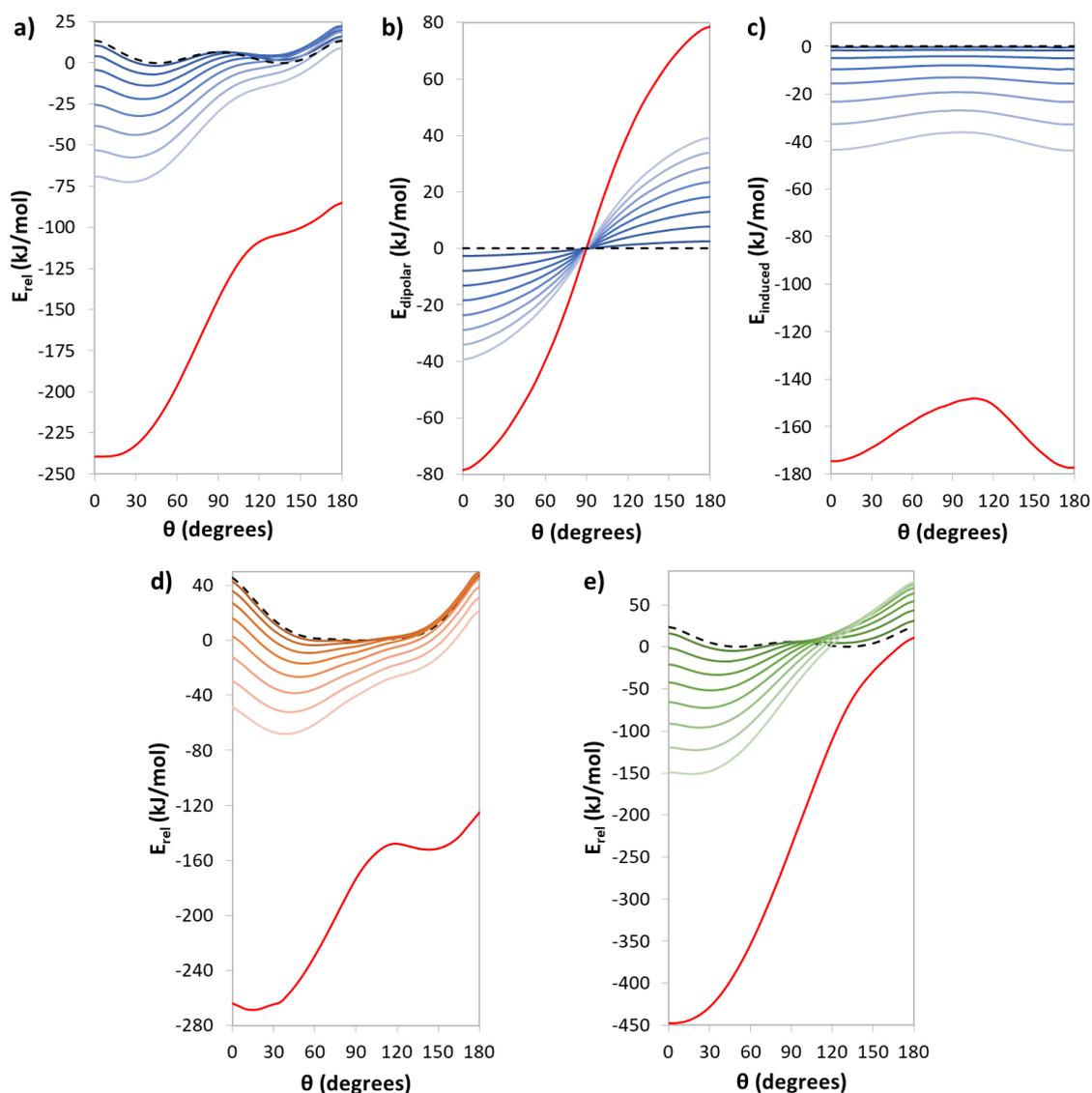


Figure 2. a) Relative total energies (E_{rel}) with respect to θ with relative to the lowest energy zero field (dashed line) conformation of the di-fluoro system. b) Energy contribution due to the permanent dipole (second term in equation 1) for the di-fluoro system. c) Energy contribution due to the induced dipole (third and higher order terms in equation 1) for the di-fluoro system. d) E_{rel} with respect to θ

relative to the lowest energy zero field (dashed line) conformation of the di-bromo system. e) E_{rel} with respect to θ relative to the lowest energy zero field (dashed line) conformation of the di-nitrile system. In all plots the applied E-field strength varies from $5.1 \times 10^6 - 7.7 \times 10^7$ V/cm (dark to light line colour) and the red lines denotes the strongest field of 1.5×10^8 V/cm.

The energy due to the permanent dipole is an asymmetric function of θ with respect to $\theta = 90$ (i.e. $-\boldsymbol{\mu} \cdot \mathbf{F} = -\mu F \cos(\theta)$) which equally and oppositely affects the energy for dipolar alignment and anti-alignment, respectively. In Fig. 2c we can see that the contribution from the induced dipolar and higher order terms to the energy is found to vary much less with respect to θ , and is always negative (i.e. energy lowering). For higher field strengths there is a tendency for a more significant energetic effect of the induced dipole for conformations close to full alignment or full anti-alignment of the two aryl rings. For lower field strengths the magnitude of the energetic contribution of the permanent dipole is higher than that of the induced dipole. In this regime induced dipole leads to a slight energetic downshifting of the θ -dependent E_{rel} curves but the shapes of the curves are largely unchanged (see S3). However, for field strengths $\geq 6.7 \times 10^7$ V/cm the energy lowering from the induced dipole is greater in magnitude than the maximum energy variation due to the permanent dipole. In such cases, the assumption that the field induced energy variation of the system only depends on the orientation of the permanent dipole, will significantly underestimate the effect of the applied E-field. This is in line with the acknowledged importance of the induced dipole in tuning the catalytic effect of external fields in reactivity.^{28,52} For the highest strength E-field considered of 1.5×10^8 V/cm, the energy lowering from the induced dipole is more than double the maximum energy variation due to the permanent dipole. Overall, in all systems the energy lowering effect of the induced dipole acts to: i) increase the energy lowering of the permanent dipole when the molecular dipole is aligned with \mathbf{F}_x , and ii) largely cancel the energetically destabilising effect of the permanent dipole (or even reverse it at high field) when the molecular dipole is anti-aligned with \mathbf{F}_x . In the SI we provide plots of the energetic contributions due to the permanent dipole and the induced dipole interacting with the considered range of E-fields for the di-bromo and di-nitrile systems.

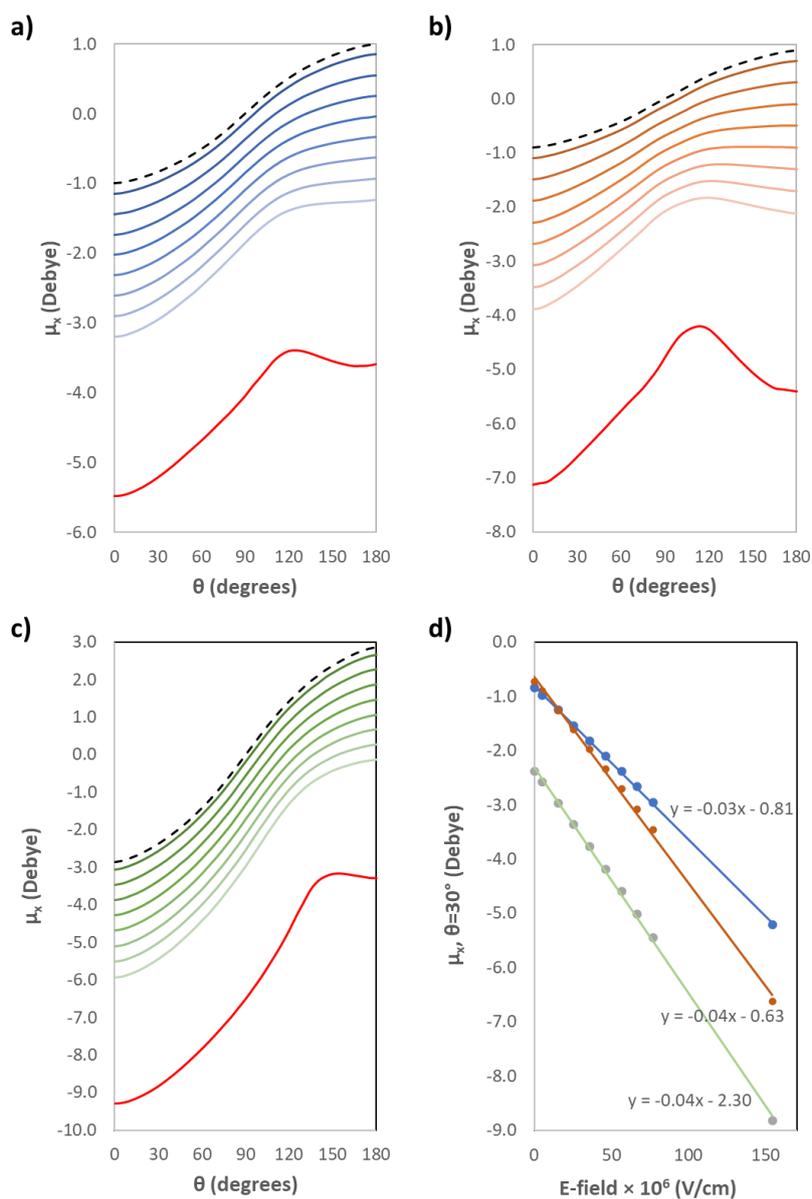


Figure 3. Variation in the x-component of the total dipole moment (μ_x) with respect to F_x and θ for: a) the di-fluoro system, b) the di-bromo-system, and c) the di-nitrile system. In d) we show linear fits to the variation of total dipole moment at $\theta = 30^\circ$ with respect to increasing E-field strength for each system.

The effect of an applied E-field can also be assessed by examining how it affects the overall the dipole moment. In Figure 3a-c we show how the x-component of the total dipole moment varies with respect to increasing field strength and θ for each system. In terms of total dipole moment the di-fluoro and di-bromo systems show very similar behaviour at zero field and the lowest field strengths. However, with increasing field strength the total dipole in the di-bromo system quickly becomes more negative than the corresponding dipole for the di-fluoro system. This is mainly due to

the higher polarisability (α) of the larger Br atoms as compared to the more compact less polarisable F substituents. This difference can be more clearly seen in Fig. 3d where the increase in the magnitude of x-component of the total dipole moment for each of the three considered systems is plotted with increasing magnitudes of F for a fixed angle of $\theta = 30^\circ$. To first order the total dipole can be approximately expressed as the sum of the permanent dipole (μ_0) and the product between α and applied E-field (i.e. $\mu = \mu_0 + \alpha F$). The gradient of a linear fit to the data in Fig. 3d thus provides us with an estimate of α for each system. Here, we see that although the di-fluoro and di-bromo systems have a similar zero field permanent dipole (i.e. a similar intercept) the steeper linear fit for the latter shows that this system has a significantly higher polarisability. Performing the same analysis for the di-nitrile system shows that the permanent dipole is significantly higher in magnitude than for the di-halogen based systems. Also, the polarisability of the di-nitrile system is almost as high as for the di-bromo system. This combination results in this system having the largest magnitude total dipole for all applied field strengths in Fig. 3d.

The above considerations can help inform us as to which type of system would form the basis for a responsive and efficient molecular switch. Firstly, for each system we identify the minimum applied field which induces a single energy minimum in the E_{tot} versus θ curve. We note that although the di-bromo system has a single energy minimum at zero applied field, a small field of 1.5×10^7 V/cm induces a weak double well energy profile which then disappears at slightly higher field strengths. For the di-nitrile system 1.5×10^7 V/cm is already enough to induce a single energy well, whereas a larger field strength of 3.6×10^7 V/cm is required for the di-fluoro system (see Fig. 4). The reason for the di-nitrile system being more efficient with respect to this measure (i.e. requires the smallest E-field) is most likely due to the higher permanent dipole in this system. We also note that the E-fields needed to induce a conformational change in our dipolar biphenyl systems are all smaller than those required (ca. 8.0×10^7 V/cm) to rotate the central phenyl ring of terphenyl derivative without a permanent dipole moment.¹³ This highlights the importance of having large permanent dipole moments when it comes to inducing conformational changes by means of modest E-fields.

Using the angular dependence of E_{tot} at a modest E-field of 3.6×10^7 V/cm we can also estimate the order of magnitude of speed at which these systems could function as switches. At this E-field strength the majority (~66%) of the field-induced drop in the total energy is determined by the $\mu \cdot F$ term in equation 1 (i.e. the dipolar approximation is reasonable). Solving equation 3 for the interval between the approximate angle of minimum energy (θ_0) and (θ_0+25 degrees) we obtain times of 0.21 ps, 0.22 ps, and 0.31 ps for the di-fluoro, di-bromo and di-nitrile systems respectively.

Although there is some variation in these times all are within a similar narrow range of 0.26 ± 0.05 ps for this 25° rotation. Assuming this is a rotation speed in these systems we can estimate that a switching cycle from a zero field minimum (taking the $\sim 135^\circ$ case for the di-fluoro and di-nitrile systems) to the E-field induced minimum ($\sim 35^\circ$), and back again, would take approximately 2 ps. We note that the fastest commercial CPUs operate with clock speeds (an indication of operations per second) well below 10 GHz.⁵³ Considering that in a fast 5 GHz CPU the time period for a full switching cycle is 200 ps, our results indicate that E-field gated molecular conformational change has the potential to provide ultrafast switching speeds.

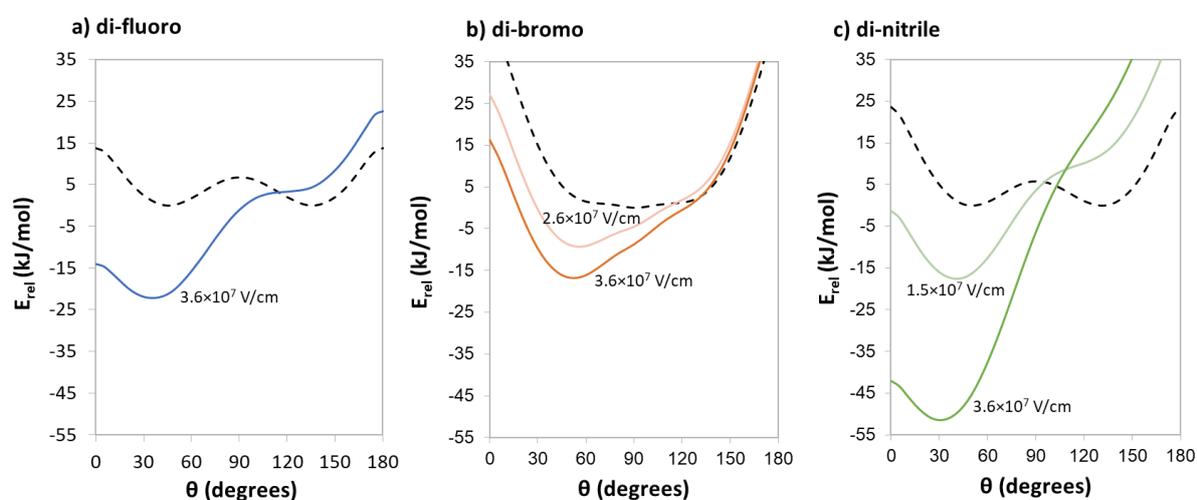


Figure 4. E_{rel} versus θ relation for low field strengths for the: a) di-fluoro, b) di-bromo, and c) di-nitrile systems. The lowest field strength shown in each plot denotes the minimum field required to induce a single energy minimum. In all cases the dashed line denotes the zero-field case. For comparison the E_{rel} behaviour for a field strength of 3.6×10^7 V/cm is also shown in each case

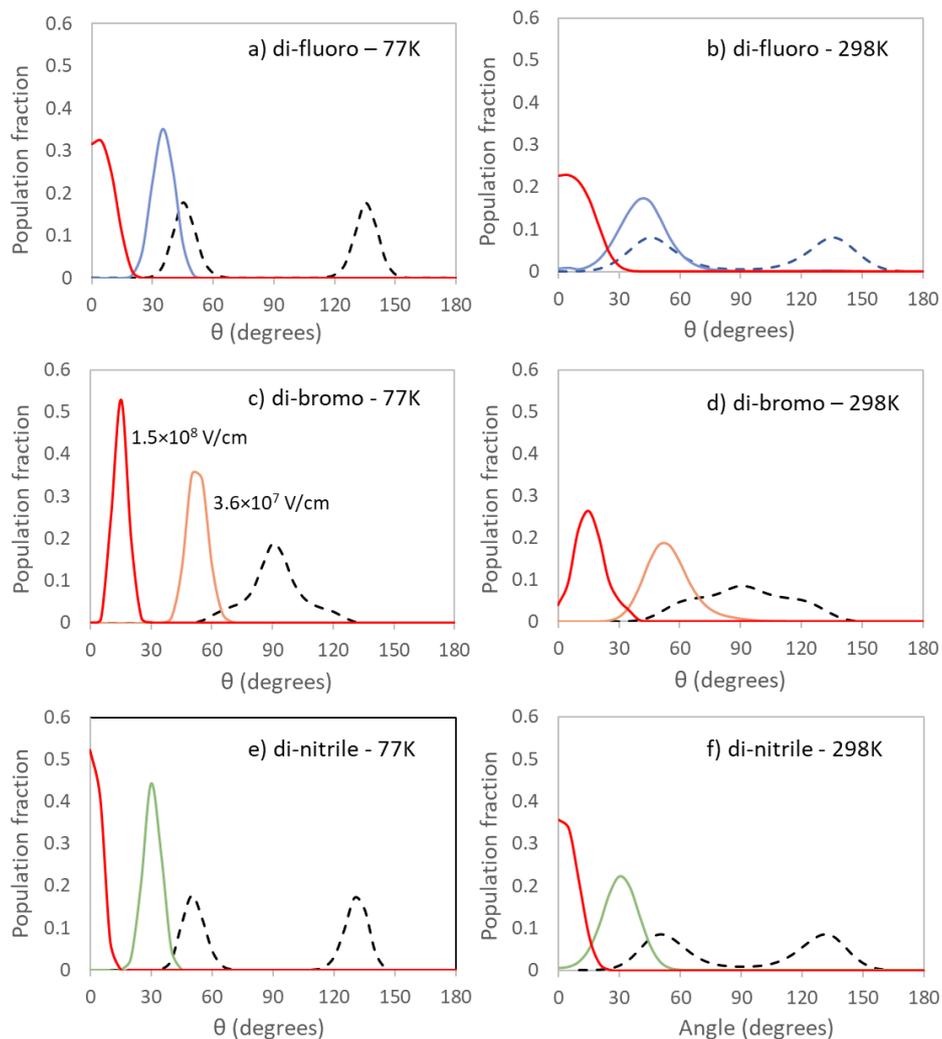


Figure 5. Boltzmann populations with respect to θ for: a) the di-fluoro system at 77K and b) at 298K, c) the di-bromo system at 77K, and d) at 298K, and e) the di-nitrile system at 77K, and f) at 298K. For each plot zero field is shown by a dashed line, the high field of 1.5×10^8 V/cm is indicated by a red line and a low field of 3.6×10^7 V/cm is shown by the remaining coloured line.

More pertinent for actual devices is the stability of a particular conformational state induced by an applied field at finite temperatures. Here, for two temperatures ($T = 77$ K and $T = 298$ K) we compute the Boltzmann statistical probabilities of finding each system at each angle for three selected E-fields (see Fig. 5). Just below the boiling point of nitrogen (77.36 K), the Boltzmann populations for 77K in Figures 5a, 5c and 5e have peaks corresponding to ranges of θ with the lowest values of E_{rel} in Figures 2 and 4. For an applied field of 3.6×10^7 V/cm, where in all cases only one E_{tot} minima remains (see Fig. 4), all systems have a corresponding single peak in their Boltzmann population fraction. For this applied field strength, the di-nitrile system shows a relatively sharp ($\sim 5^\circ$ full width at half maximum - FWHM) and reasonably high (~ 0.45 fraction) peak centred at 30° , in line with the corresponding deep

sharp E_{tot} minimum (see Fig. 4c). In comparison, the di-fluoro and di-bromo systems have peaks with slightly lower heights (<0.4 fraction), slightly larger FWHM values (15° and 10° respectively) and positions 5° below and 10° above the di-nitrile peak, respectively. This indicates that the higher permanent dipole of the di-nitrile system leads to a greater degree of switching selectivity.

Progressively higher E-field strengths cause an increased planarity of the dipolar aryl ring and thus the peaks tend to be shifted towards 0° . For a field of 1.5×10^8 V/cm this shift is most pronounced for the di-fluoro and di-nitrile systems where the peak is centered at 0° . For the more sterically hindered di-bromo system the high field peak is centered around 15° . For a higher ambient temperature of 298 K there is a corresponding slight peak broadening (by 5 - 10°) and lowering of peak heights ($\sim 50\%$). This effect reflects the occupation of θ conformations with higher E_{rel} values. Generally, however, the increase in temperature does not lead to any significant change in the overall peaks positions or number of peaks, indicating that the conformational effects of an applied E-field would be observable under ambient conditions for all systems.

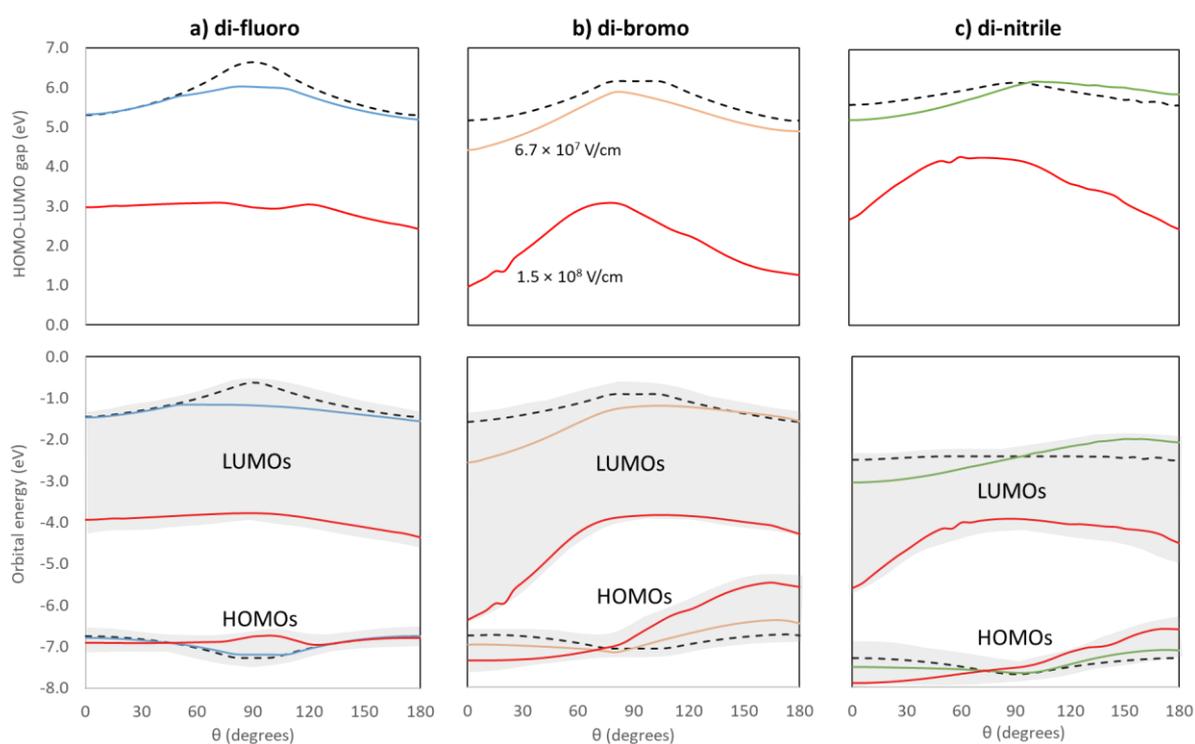


Figure 6. HOMO-LUMO gap energies (upper) and energies of individual HOMO and LUMO orbital energies with respect to θ for: a) the di-fluoro system, b) the di-bromo system, and c) the di-nitrile. All plots show data for zero field (dashed line), a field of 1.5×10^8 V/cm (red line) and a moderate field of 6.7×10^7 V/cm (remaining coloured line).

The effect of both varying θ and the applied E-field strength on the electronic structure of each system is shown in Figure 6. For zero field the average gap over the full range of θ between the energy of the highest occupied molecular orbital (HOMO) and that of the lowest unoccupied molecular orbital (LUMO) is similar for the di-halogen systems (5.83 eV for di-fluoro cases and 6.18 eV for di-bromo), but is somewhat lower for the di-nitrile system (5.01 eV). The maximum energy gap variation with respect to changes in θ is also less for the di-nitrile system (0.5 eV) than for the di-fluoro (1.3 eV) and di-bromo (1.0 eV) system. The most conspicuous trend in the zero field gap values with respect to θ is that all systems show a tendency to have a peak gap value close to $\theta = 90^\circ$. This behavior is typically linked to a corresponding maximum and minimum in the θ -dependent values of the respective LUMO and HOMO values. We note that $\theta = 90^\circ$ is where π -conjugation in all systems is minimized. The destabilization of the LUMO for this highly twisted conformation strongly suggests that the LUMO tends to have a stable delocalized character through the biphenyl bridge in each system when planar. In Figure 7 we show the LUMO orbital for the case of the di-fluoro system for $\theta = 0^\circ$ and $\theta = 90^\circ$ to confirm this. We note that electron conduction through the biphenyl system is known experimentally to decrease with decreasing π -conjugation (i.e. θ values closer to 90°)³⁰ indicating that the LUMO is likely to be one of the main channels for electron conduction in these systems. The general θ -dependent electronic response observed for zero field is quite robust with respect to applied E-fields and only starts to vary significantly for fields approximately $\geq 6 \times 10^7$ V/cm (see Figure 6). This confirms that the E-field gated conformational changes realized in these systems would yield associated π -conjugation changes, and thus conductive switching, up to relatively high field strengths, without modulating the electronic structure.

For relatively high fields, the energy of the LUMO is drastically lowered for all three biphenyl systems, whereas the energy of the HOMO is less affected. As a result, a large decrease of the HOMO-LUMO gap is observed for high fields. The large stabilization of the LUMOs in going from zero field to high fields originates in a drastic orbital shape change. As shown in Figure 7, the LUMO for the di-fluoro system for both $\theta = 0^\circ$ and $\theta = 90^\circ$ transforms into a Rydberg orbital under high electric fields. The virtual molecular orbitals (MOs) obtained at $F_x = 1.5 \times 10^8$ V/cm lying higher in energy than the LUMO also include a MO that is similar to the LUMO at zero field, as shown in Figure S4. The energy of this MO remains largely unchanged upon application of the electric field. The virtual MOs of the molecule at zero field, in turn, also include Rydberg states but they lie high in energy (see Figure S4). A comparison between the energies of the MOs displayed in Figure 7 with those of Figure S4 allows us to conclude transformation of the LUMO results from a large stabilization of one of the Rydberg orbitals upon application of large E-fields. Such E-field-induced stabilization of Rydberg orbitals is in line with the results recently found for ethylene in intense laser fields.⁵⁴ In broader terms, the strong

decrease of the HOMO-LUMO gap found for biphenyl derivatives under large fields is also in line with the results reported in Ref. 55 for diatomic molecules, where a large decrease in the gap was observed for E-fields larger than 5×10^7 V/cm.

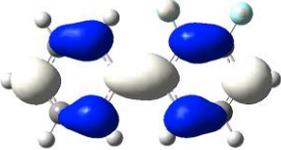
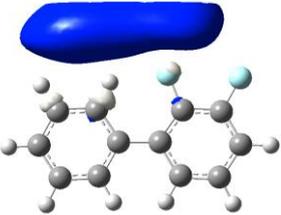
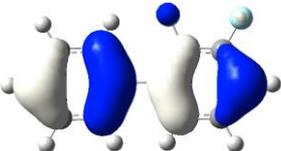
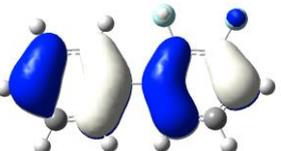
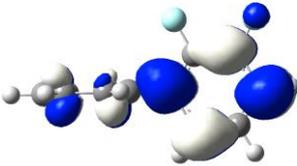
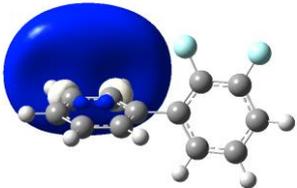
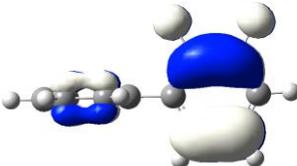
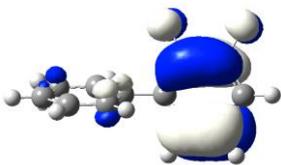
$\theta = 0^\circ$	$F_x = 0$ V/cm	$F_x = 1.5 \times 10^8$ V/cm
LUMO	 $E = -1.45$ eV	 $E = 3.94$ eV
HOMO	 $E = -6.75$ eV	 $E = -6.91$ eV
$\theta = 90^\circ$		
LUMO	 $E = -0.63$ eV	 $E = -3.78$ eV
HOMO	 $E = -7.28$ eV	 $E = -6.75$ eV

Figure 7. HOMO and LUMO orbitals and energies of the di-fluoro-biphenyl molecule for $\theta = 0^\circ$ and $\theta = 90^\circ$, as obtained at zero field and at the highest E-field employed in this study.

The E-field induced reshuffling of MOs and the concomitant change of the energy gap demonstrates that in the high-field regime, the physical properties of biphenyl systems will be governed both by the

induced conformation and the E-field induced changes in the electronic structure. The presence of E-field-stabilised Rydberg orbitals will likely cause significant changes in both the conduction and optical properties of biphenyl moieties.

Conclusions

We model the E-field response of the biphenyl system with three choices of dipolar substituents of one of the two aryl rings. In all cases, the E-field strengths required to induce conformational rearrangements in the twistable functionalized ring are similar to the largest E-field strengths used in silicon-based MOSFETS, and are well within those commonly found in biomolecular systems. At smaller E-field strengths the main variation in the total energy is due to the interaction of the permanent dipole with the applied field. However, the proportion of the energy change associated with the induced dipole and higher order polarization effects is always significant and rapidly becomes the dominant energy term with increasing E-field strength. This result clearly shows that E-field induced effects cannot be neglected when attempting to understand and model E-field responsive molecular switching. In line with this finding, we find that substituents that confer both a large permanent dipole moment and a modest polarizability (e.g. CN) entail that smaller E-fields are needed to trigger barrierless conformational rearrangements. For such systems the conformational switching selectivity is also likely to be enhanced. Practically, this suggests that such substituents could help design more energy efficient E-field responsive molecular switches. Looking further to eventual deployment of such switches as components in real devices we have also established that: i) the switching capability of dipolar aryl rings mediated by E-fields is robust at room temperature, ii) the electronic structure of these biphenyl systems is little affected at moderate fields - but can be strongly affected with high E-fields (which can energetically stabilise Rydberg LUMO orbitals that can drastically modify the physical properties of the system), iii) the approximate switching cycle times of twistable dipolar aryl rings estimated to be two orders of magnitude faster than the time period for the fastest commercial CPUs. These features suggest that dipolar conformational molecular switches based on substituted biphenyls or similar molecular systems could be a promising future basis for stable high performance computing in real world devices.

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