Structural causes of singlet/triplet preferences of Norrish Type II reactions in carbonyls

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Abstract

Photolysis thresholds are calculated for the Norrish Type II (NTII) intramolecular γ -hydrogen abstraction reaction in 22 structurally informative carbonyl species. The B2GP-PLYP excited state S_1 and T_1 thresholds agree well with triplet quenching experiments. However, many linear-response methods deliver poor S_1 energetics, which is explained by a S_1/S_0 conical intersection in close proximity to the S_1 transition state. Multiconfigurational CASSCF calculations confirm a conical intersection features across all carbonyl classes.

Structure–activity relationships are determined that could be used in atmospheric carbonyl photochemsitry modelling. This is exemplified for butanal, whose NTII quantum yields are too low when used as a 'surrogate' for larger carbonyls, since but anal lacks the γ -substitution that stabilises the 1,4biradical. Reaction on T_1 dominates only in species where the S_1 thresholds are high — typically ketones. The α, β unsaturated carbonyls cannot cleave the α - β bond, causing them to photoisomerise. A concerted S_0 NTII mechanism is calculated to be viable and may explain the recent detection of NTII photoproducts in the photolysis of pentan-2-one below the T_1 threshold.

Introduction

The Norrish Type II reaction mechanism

The Norrish Type II (NTII) photolysis reaction is available in larger carbonyl species that have a γ -hydrogen. As shown schematically in Figure 1, the NTII reaction involves intramolecular H-abstraction by a photoexcited C=O chromophore to form a biradical.^{1–3} The NTII reaction dominates the photolysis of carbonyls with main alkyl chains of four carbon atoms or longer.^{4,5} NTII reactions can involve both the excited electronic singlet (S_1) and triplet (T_1) states as shown in Figure 1 below. The propensity for singlet or triplet reactions depends on the particular carbonyl species,^{6,7} and is dictated by the S_1 and T_1 NTII photolysis threshold for each carbonyl. In principle a similar S_0 pathway could also exist, and is examined in this paper.

With reference to Figure 1, if the available photon energy is above the photolysis threshold for the S_1 cyclic transition state (TS) then the 1,5–H-shift can occur on S_1 , leading to a singlet biradical.⁸ However, if the photon energy is near or below the S_1 NTII threshold then reaction on S_1 will be slow and intersystem crossing (ISC) to T_1 can occur, resulting in the formation of a triplet biradical. The triplet biradical is long lived because it correlates with excited state products, that is, either the alkene or enol NTII product is in an electronic excited



Figure 1: Schematic illustrating the Norrish type II reaction pathways on S_1 (*light blue*) and T_1 (*orange*). Following photoexcitation to S_1 the carbonyl forms a cyclic 6-membered transition state and undergoes a 1,5– H-shift to form a 1,4-biradical. If the available energy is above the S_1 photolysis threshold a singlet biradical will be formed, otherwise intersystem crossing (ISC) to T_1 can occur, leading to a triplet biradical. A S_1/S_0 conical intersection is proposed to occur along the S_1 reaction pathway, allowing ultrafast crossing to S_0 .

state, and these products are not accessible at actinic UV energies (> 300 nm). The alkene and enol photoproducts observed in NTII reactions are ground state species, thus for NTII reaction to occur via T_1 there must be ISC to the S_0 state. A conical intersection (CI) between S_1 and S_0 is also proposed along the S_1 NTII reaction coordinate,⁹ lowering the lifetime of the singlet biradical by allowing ultrafast crossing to S_0 and hence rapid formation of NTII photoproducts. The different lifetimes of the singlet and triplet biradicals are also reflected in different NTII products observed for reaction on S_1 and T_1 , such as cyclisation or retention of stereochemistry.¹⁰

The 1,4-biradical can undergo different subsequent reactions: either cleavage of the α , β bond to form an enol and an alkene (Scheme 1), or Norrish Yang cyclisation (NYC) to form a substituted cyclobutanol^{11,12} (Scheme 2). Additionally, reaction from the biradical can be frustrated, as the abstracted H-atom can backtransfer to re-form the original carbonyl.



Scheme 1: Norrish Type II cleavage to an alkene and enol.



Scheme 2: Norrish Yang cyclisation to a substituted cyclobutanol.

The prevalence of the reactions shown in Schemes 1 and 2 depends on: the carbonyl substituents,^{10,13} chemical environment,^{14,15} and degree of energetic and conformational stabilisation of the biradical intermediate.^{1,3} Each of these factors influence the conformation of the *p*-type radical singly occupied molecular orbitals (SOMOs) of the biradical. The cleavage of the biradical to an alkene and an enol requires the SOMOs to overlap the sites of the newly forming C=C bonds (Scheme 1), and these gauche conformations are readily accessible in carbonyls without bulky α -substituents.¹⁰ Consequently, cleavage is the dominant NTII photolysis reaction in the gas phase for linear carbonyls.

Norrish Yang cyclisation occurs in many carbonyl species, ^{10,12,16} and requires a conformation with in-phase overlap of the biradical SOMOs to form a new C–C bond (Scheme 2). The NYC reaction is common for solvated carbonyls, ^{15,16} and a synthetic strategy to form 4-membered rings is to use carbonyl substituents that promote conformations with overlapping SOMOs for cyclisation.¹⁷

However, NYC also occurs in the gas phase, provided a cyclisation conformation can be reached before cleavage to an alkene and enol occurs. Indeed, several classes of carbonyl species exclusively undergo NYC upon photolysis due to conformational effects. These include α -diketones,^{18,19} α -fluoro ketones,²⁰ and β , γ unsaturated ketones.²¹ Additionally, intramolecular abstraction of hydrogens further than the γ -position can occur,^{22,23} but quantum yields (QYs) for these alternative abstractions are low whenever a γ -H is available²⁴ and these alternate abstractions must result in cyclisation rather than cleavage.²⁵ The canonical 6membered NTII TS to form a 1,4-biradical is the most energetically favourable intramolecular H-abstraction pathway due to its low degree of ring strain,^{17,26} and so is the only TS discussed in this paper.

This paper is concerned with the NTII reaction in gas-phase carbonyls that bear the simple hydrocarbon chains typical of the volatile organic compounds (VOCs) found in high concentrations in the troposphere.^{27,28} Under tropospheric conditions NTII predominantly leads to alkene and enol photoproducts.²⁹ Indeed, these are the only NTII photoproducts considered in the Master Chemical Mechanism.^{30,31} However, NYC is still observed as a minor channel in gas-phase photolysis of these 'typical' carbonyls,^{11,32} and the occurrence of a cyclisation product provides mechanistic insight.

In recent gas-phase FT-IR photolysis experiments on pentan-2-one, Kharazmi observed 1methylcyclobutanol as a photoproduct.³² For this to occur the biradical must be able to undergo internal bond rotation to access the conformation shown in Scheme 2, which does not happen with singlet biradicals as they retain stereochemistry in their photoproducts.⁸ Additionally, triplet quenching studies show that NYC photoproducts arise purely from triplet NTII reactions, since the T_1 biradical can undergo conformational rearrangement before cleavage occurs.^{12,33}

This T_1 biradical cyclisation mechanism is consistent with the lifetime of the biradical measured in femtosecond mass spectrometry experiments on a series of ketones: the 1,5–H-shift occurs on a 70–90 fs timescale, while subsequent cleavage from this biradical is an order of magnitude slower at 400–700 fs.³⁴

The photolysis data by Kharazmi on 1methylcyclobutanol formation from pentan-2one are also consistent with a triplet cyclisation mechanism. The QY of the NYC product of pentan-2-one is ~0.05 at 300 nm (above the expected S_1 NTII threshold) and increases for longer wavelengths, up to ~0.17 at 325 nm.³² This suggests that NYC becomes increasingly important at lower photolysis energies since reaction on S_1 is no longer rapid, allowing ISC and reaction on T_1 to occur. The QY of 1methylcyclobutanol then drops below 0.05 at 330 nm, indicating that the T_1 NTII photolysis threshold also becomes inaccessible at this lower energy.

In order to understand the photolysis of larger

carbonyls, it is therefore crucial to have accurate estimates of S_1 and T_1 NTII thresholds. The cause of the short S_1 biradical lifetimes, and the high rate of H-backtransfer reversion to the starting carbonyl, was not well understood mechanistically when data was obtained in the 1970s.³⁵ To explain these phenomena, the presence of a S_1/S_0 conical intersection (CI) along the NTII reaction pathway was proposed in the 1990s.^{9,36} The possibility of a concerted S_0 NTII photolysis pathway has also been proposed,¹⁶ and recent wavelength-dependent FT-IR photolysis data shows the formation of NTII photoproducts in pentan-2-one at energies below the predicted T_1 NTII threshold.³²

Dataset of NTII-capable carbonyls

This paper aims to derive NTII threshold structure-activity relationships (SARs) generalisable to a large variety of carbonyl species. A significant proportion of biogenic volatile organic compounds (VOCs) are larger (5-6 carbon) molecules.^{28,37–39} This includes the atmospherically abundant isoprene-derived C5 carbonyls,⁴⁰ and the C6 carbonyls like 'leaf-aldehyde' (hex-3-enal).⁴¹

Carbonyls larger than C6 are not studied here because: the volatility and abundance of carbonyls decrease with chain length;²⁸ substitution patterns adjacent to incipient γ -C[•] and C[•]-O radical sites can still be studied with molecules of this size; and the computational expense when performing $\mathcal{O}(N^7)$ (EOM-)CCSD calculations is limited.

The 22 NTII-capable carbonyl species studied here are shown in Scheme 3, and their systematic names are listed below by carbonyl class:

- Saturated aldehydes: butanal, pentanal, 2-methylpentanal, 3-methylpentanal, 4-methylpentanal, 2-ethylbutanal, 3,3dimethylbutanal, hexanal, hex-3-enal;
- Saturated ketones: pentan-2-one, hexan-2-one; 3-methylpentan-2-one, 4methylpentan-2-one;
- *α*,*β***-Unsaturated aldehydes:** but-2enal, 2-methylpent-2-enal, 3-methylpent-



Scheme 3: Dataset of carbonyl bearing γ -hydrogens, for which Norrish Type II reaction thresholds are calculated.

2-enal, hex-2-enal;

- *α*,*β***-Unsaturated ketones:** hex-3-en-2one, hex-1-en-3-one, hex-4-en-3-one;
- Dicarbonyls: 2-oxobutanal.

Previous electronic structure calculations of the NTII reaction

This paper is only concerned with NTII reaction energetics to determine photolysis SARs. Theoretical work on other mechanistic aspects of the NTII reaction, such as the dynamics of the biradical, are available in other studies, but only for a few carbonyl species.^{42–44} An overview of the previous modelling work for the NTII reaction in carbonyls is presented below to give context to this current study.

Early electronic structure calculations Early MINDO/3 semi-empirical calculations in 1978 were able to provide the key mechanistic details of the S_1 and T_1 NTII reactions, including the importance of the biradical intermediate,⁴⁵ though the method mistakenly identifies the S_1/S_0 region of degeneracy on the path to the biradical formation as an avoided crossing, rather than a true conical intersection. Later modelling on T_1 using UHF and UMP2 methods made clear the preference for NTII γ -H abstraction via a planar 6-membered TS, and its similarity to the ground state retro-ene and McLafferty rearrangements,⁴⁶ which will be discussed later in the context of an S_0 NTII mechanism. Semi-empirical methods have also been

used to study other mechanistic steps involving the 1,4-biradical, including cyclisation and H-backtransfer pathways.⁴⁷

Electronic states of NTII in unsaturated carbonyls The photochemistry of butyrophenone has been modelled using multiconfigurational complete active space (CAS) methods.⁴⁸ Butyrophenone has a phenyl group conjugated with the C=O moiety and as an aromatic carbonyls it has similarities to the α , β -unsaturated carbonyls discussed here. This includes a T_2 state that is energetically accessible at tropospheric photon energies,⁴⁹ as well as a $S_1/T_2/T_1$ three-state intersection. This intersection is not at configurations accessed by the triplet biradical, and does not intersect with S_0 , and so does not decrease the triplet biradical lifetime. However, an S_1/S_0 CI was located. From this, it can be surmised that only the T_1 NTII reaction results in a biradical with significant lifetime, even in species where conjugation to the C=O moiety complicates the electronic states involved.

NTII/NTIa competition and substituent effects Both the Norrish Type I α -bond cleavage (NTIa) and NTII reactions in butanal have been modelled using the CASSCF, DFT, and MP2 levels of theory by Chen and Fang.⁵⁰ These calculations demonstrate that in the tropospheric energy range, NTIa reaction in butanal only occurs on T_1 after ISC; whereas NTII can occur on both the S_1 or T_1 states. Chen and Fang also located an S_1/S_0 CI along the S_1 NTII reaction coordinate, after the S_1 1,5–H-shift TS. A recent theoretical study highlights the role of 'singlet state only' NTII and NTIa mechanisms in 5-methylhexan-2-one derivatives.⁵¹ While this study also calculates reaction enthalpies on substituted reaction products, this is estimated with single point energy calculations on products and not by locating TSs, and so does not provide data on photolysis thresholds.

Dynamics of the biradical The dynamics of the triplet 1,4-biradical formed from pentanal have been simulated using semi-empirical multireference configuration interaction (MRCI) cal-

culations at two timescales: 10 ps and 100 ps. 43 In the 10 ps simulations both NTIa and the NTII γ -H transfer occurred. In the 100 ps simulations neither the biradical cyclisation nor cleavage were observed, inconsistent with the 400–700 fs biradical lifetime estimated by femtosecond mass-spectrometry experiments.³⁴ Nevertheless, these dynamics simulations indicate that the 1,5-H-shift and subsequent reactions occur on separate time-scales in the triplet biradical, allowing conformational rearrangements to occur. The MRCI dynamic simulations predicted a NTIa QY of 0.34 and NTII QY of 0.66 on the 100ps timescale, in qualitative but not quantitative agreement with observed NTII QYs of ~ 0.8 in pentanal. 4,52

Multiconfigurational calculations on prototypical aldehydes Kletskii *et al.* used the multiconfigurational XMCODPT2 method⁵³ to characterise Norrish reactions in butanal, pentanal, and but-2-enal.⁵⁴ XMCQPDT2 provides a correct representation of wavefunctions with multireference character, and regions with electronic state degeneracies such as conical intersections, both of which are incorrectly modelled with linear response (LR) time-dependent density functional theory (TD-DFT) or equationof-motion coupled cluster (EOM-CC) methods.⁵⁵ However, like CASPT2, the XMCQDPT2 method is more computationally demanding than single-reference methods,⁵⁶ and the energetics delivered for carbonyl photochemistry can often be worse than those from many DFT functionals.²⁶ While the Kletskii *et al.* paper provides the largest qualitatively correct survey of the electronic structure and energies along the NTII reaction for a few carbonyls (butanal, pentanal, pentan-2-one, but-2-enal), we wish here to expand the study of the NTII to many more more molecules.

Computational Methods

The computational protocols used in this paper to study S_1 , T_1 , and S_0 NTII photolysis thresholds are largely the same as reported by the authors in an earlier publication on the NTI reaction in carbonyls.⁴⁹ Photolysis thresholds are calculated by the addition of the intrinsic reaction barrier (relative to the relevant electronic state minimum) to the 0–0 excitation energy (from the S_0 minimum to the excited state minimum). The 0–0 excitation energies used here are calculated using protocol detailed by the authors in an earlier paper studying carbonyl excitation energies.⁵⁷ The methods are summarised below.

$S_0 \rightarrow S_1$ Excitation energies

Briefly, the chosen protocol for 0–0 excitation energies is: EOM-CCSD vertical excitation energies scaled by factors derived from CCSDR(3) results,⁵⁷ EOM-CCSD//TD-CAM-B3LYP S_1 relaxation energies, and (TD-)CAM-B3LYP zeropoint vibrational energies (ZPVEs). A moleculeindependent energy shift, equal to the mean signed error of the 0–0 excitation energy protocol, is applied to correct for the lack of triples in the calculated relaxation energy, as discussed in the earlier carbonyl excitation paper.⁵⁷

Optimisation of S_1 1,5–H-shift transition states

In our earlier study on the NTI reaction, photolysis thresholds and intrinsic reaction barriers determined using TD-B3LYP were found to reasonable agreement with respect to experimental data.⁴⁹ Here, however, the presence of an S_1/S_0 CI near S_1 biradicals and/or TS structures means that linear response theories, like TD-DFT, may not be appropriate and multiconfigurational methods must be used.⁵⁸ This, and other potential limitations of the theoretical methods used, are discussed below.

Hydrogen-transfer TSs can be difficult to locate with DFT methods, and usually require functionals with a high component of Fock exchange.^{59,60} Notably, B3LYP is poor at describing hydrogen transfer TSs^{61,62} and neither TD-B3LYP calculations nor TD-BH&HLYP calculations (with 50% HF) could locate the S_1 1,5– H-shift TSs here. Instead, TD-CAM-B3LYP is used to successfully optimise the required S_1 TSs, noting that, from earlier work, 49,57 the TD-B3LYP and TD-CAM-B3LYP S_1 results had comparable accuracy. Single point energies were determined on S_1 at the TD-B2GP-PLYP level of theory using TD-CAM-B3LYP geometries.

A zero-point vibrational energy (ZPVE) scaling factor of 0.976 was used for S_1 TD-CAM-B3LYP/def2-TZVP ZPVE corrections. This value was derived for CAM-B3LYP/ma-TZVP calculations, but it has been shown that the minimal augmentation has negligible effect on the scale factor.^{63,64} It is also assumed, in the absence of additional data, that the scale factor for S_0 ZPVEs is also applicable to S_1 . As discussed previously,⁴⁹ the few literature studies on excited state scaling factors^{65,66} indicate that this assumption will lead to differences of less than 1 kJ/mol in ZPVE corrected energies.

The 1,5–H-shift TS to formation of the 1,4biradical, shown in Figure 1, is also likely to have significant biradical character. The S_1/S_0 CIs are expected to be near the reaction coordinate for formation of the 1,4-biradical,^{9,54} and again will have biradical character. In principle all of these configurations should be described using multiconfigurational methods.^{67,68} For S_0 biradicals, DFT functionals are able to deliver reasonable energies.^{69,70} For carbonyls specifically, previous PhD work by Shaw on carbonyl epoxidation reactions was able to predict S_0 biradical TSs accurately with UM06-2X calculations.²⁶ This same work showed that carbonyl T_1 stationary points with significant biradical character calculated with UM06-2X level were within 10 kJ/mol benchmark results from higher levels of theory.²⁶

This work also studied carbonyl 1,5–H-shift TSs on S_1 using DFT and multiconfigurational methods, and while TD-M06-2X intrinsic barriers were ~20 kJ/mol higher than those calculated using the XMCQDPT2 method (see Table S1), the TD-DFT results appeared to be in better agreement with experiment. Thus, somewhat fortuitously, DFT-based methods appear capable of predicting S_1 NTII energies to somewhat acceptable accuracy — assuming TD-DFT formalism is valid, that is, the effects of the S_1/S_0 CI are negligible.

CASSCF calculations

Multiconfigurational CASSCF calculations were used to determine a qualitatively accurate picture of any S_1/S_0 minimum energy conical intersections (MECIs). Gaussian 09 or Gaussian 16 were used for all CASSCF calculations and MECI searches. TD-CAM-B3LYP S_1 NTII TS geometries were used as the starting points for all CASSCF MECI searches, since the S_1/S_0 MECIs were assumed to be in close proximity to the S_1 1,5–H-shift TSs.

An (8,7) active space was used for initial optimisations, matching previous literature.^{50,54,71} The four occupied orbitals in the active space are the: $\gamma C-H \sigma$, $C-O \sigma$, $C=O \pi$, and oxygen lone pair n orbitals; and the virtual orbitals are the: $\gamma C-H \sigma^*$, C-O σ^* , and C=O π^* orbitals. This was then expanded to a (10,8) active space by including the other oxygen-centred non-bonding orbital (see Figure S1). For unsaturated species the π and π^* orbitals at the point of unsaturation were also included in the active space, and in these species the C–O σ and σ^* orbitals were removed to keep a manageable (10,8) active space (see Figure S2). The MECI structures calculated here with the 6-31+G(d) basis set for butanal, pentanal, pentan-2-one, and but-2-enal were consistent with those calculated by Kletskii et al. with the larger 6-311++G(d,p) basis set.⁵⁴ Further details are provided in the supporting information.

Results

Norrish Type II reaction energies

Our best calculations for: S_1 , T_1 , and S_0 NTII photolysis thresholds relative to the S_0 minimum; intrinsic reaction barriers relative to the corresponding S_1 or T_1 minimum; and $S_0 \rightarrow T_1$ and $S_0 \rightarrow S_1$ 0–0 excitation energies; are reported in Table 1. All energies in Table 1 are ZPVE corrected.

To the author's knowledge, no direct experimental determination of photolysis threshold energies for the NTII reaction in any carbonyl has been published. An indication of the S_1 and T_1 photolysis thresholds can be inferred from

	0–0 Excitation Energy		Photolysis Threshold			Intrinsic Barrier	
	$S_0 \rightarrow S_1$	$S_0 \rightarrow T_1$	S_1	T_1	S_0	S_1	T_1
	Composite ^a	B2GP-PLYP	Composite	B2GI	P-PLYP	B2GP-PLYP// CAM-B3LYP	B2GP-PLYP
Aldehydes:							
Butanal	353	325	$382 \ (\sim 382)^b$	368	239	29	43
Pentanal	352	324	369	359	236	17	35
Hexanal	354	320	374	358	232	20	39
2-Ethylbutanal	362	324	376	359	225	14	35
2-Methylpentanal	348	323	359	355	227	11	32
3-Methylpentanal	350	321	367	358	225	17	37
4-Methylpentanal	352	324	362	351	233	10	27
3-3-Dimethylbutanal	346	318	372	364	224	26	46
Hex-3-enal Ketones:	349	320	387	376	280	38	56
Pentan-2-one	371	339	$399 \ (>382)^b$	381	242	28	42
Hexan-2-one	366	336	$381(<382)^{b}$	370	239	15	34
3-Methylpentan-2-one	358	333	379	374	231	21	42
4-Methylpentan-2-one	362	335	387	382	239	25	46
Hexan-3-one Enals:	368	338	394	383	243	26	45
But-2-enal	330	298	335	323		5	25
Hex-2-enal	327	298	318	314		-9	16
2-Methylpent- 2 -enal ^c	335	247^{c}	323	321		-12	74^c
3-Methylpent-2-enal Enones:	320	291	311	308		-9	17
Hex-3-en-2-one	315	296	319	316		4	20
Hex-1-en-3-one	313	288	341	332	229	28	44
Hex-4-en-3-one Dicarbonyls:	324	298	341	330		17	32
2-Oxobutanal	251	125	297	148	128	46	23

Table 1: 0-0 excitation energies, NTII photolysis thresholds, and NTII intrinsic reaction barriers, as indicated. (kJ/mol)

^a Scaled EOM-CCSD//B2GP-PLYP E^{vert} + EOM-CCSD//TD-CAM-B3LYP E^{relax} + (TD-)CAM-B3LYP ΔE^{ZPVE} + Δ HLC, as described in reference 57.

 b Experimental photolysis threshold energy bounds are shown in parentheses, estimated from triplet quenching studies (see text).

^c 2-Methylpent-2-enal has a low calculated $S_0 \rightarrow T_1$ 0–0 energy, and high T_1 NTII barrier, because the S_0 minimum and T_1 NTII TS are calculated to be planar, while the alkyl group in the T_1 minimum is calculated to be twisted around the C=C bond.

triplet quenching studies of particular carbonyl species, or the formation of NTII photoproducts thought to arise only from T_1 NTII reactions (such as cyclobutanols).¹⁰ Experimental triplet quenching studies on butanal, pentan-2one, and hexan-2-one have been performed using a single photolysis wavelength of 313 nm (382 kJ/mol).^{6,7} For butanal both singlet and triplet NTII products were observed, indicating a S_1 NTII threshold of approximately 382 kJ/mol, since reaction on S_1 must be relatively slow for ISC to be competitive.⁷ Only triplet NTII products were observed in the 313 nm photolysis of pentan-2-one, indicating an S_1 NTII photolysis threshold above 382 kJ/mol, and a T_1 threshold below 382 kJ/mol.⁶ The NTII photoproducts for

hexan-2-one at 313 nm were all unquenchable by triplet quenchers, *i.e.* the photoproducts were all singlets, and so the S_1 NTII threshold should be below 382 kJ/mol.⁶

The NTII S_1 and T_1 photolysis thresholds predictions in Table 1 are consistent with these experiments: the S_1 NTII threshold predicted for butanal (382 kJ/mol) is close to 382 kJ/mol, while pentan-2-one is predicted to have a higher S_1 NTII threshold (399 kJ/mol), and hexan-2-one a slightly lower S_1 NTII threshold (381 kJ/mol). The predicted thresholds in Table 1 are also in quantitative agreement with TD-M06-2X/6-311++G(3df,2p) calculations by Shaw, which predict S_1 NTII photolysis thresholds of: 385 kJ/mol for butanal, 368 kJ/mol for pentanal, 398 kJ/mol for pentan-2-one, and 379 kJ/mol for hexan-2-one. 26

More broadly, there is a lack of NTII photolysis data on the other carbonyl species in Table 1. However the predicted energies in Table 1 are consistent with some general features of NTII photolysis:

- The relative proportion of S_1 and T_1 NTII reaction products of saturated carbonyls is wavelength-dependent within the actinic range.³² This provides upper and lower bounds for the NTII photolysis thresholds of ~300–330 nm (362.5–400 kJ/mol), which agrees with Table 1.
- Photolysis is competitive between the S_1 and T_1 states, implying that their photolysis thresholds must be similar in energy. Given that $S_0 \rightarrow S_1$ 0–0 excitation energies are higher than $S_0 \rightarrow T_1$ 0–0 excitation energies, the NTII intrinsic reaction barriers must be lower on S_1 . In contrast, NTI*a* intrinsic reaction barriers are 50–70 kJ/mol higher on S_1 than on T_1 .⁴⁹
- Predicted NTII thresholds for α, β unsaturated carbonyls are lower than those for saturated species. On both S_1 and T_1 , these thresholds are all lower than 340 kJ/mol. This is partly due to the lower 0–0 excitation energies of α, β -unsaturated carbonyls due to delocalisation of their frontier orbitals,^{49,57} but also low intrinsic barriers due to an allylic resonance stabilisation effect that will be discussed later.

The TD-B2GP-PLYP//TD-CAM-B3LYP S_1 NTII intrinsic reaction barriers in Table 1 are small, mostly under 30 kJ/mol. For some α, β unsaturated species they are negative, indicating a barrierless S_1 NTII reaction pathway. Below, these S_1 intrinsic reaction barriers are calculated using a range of single point energy methods, and issues in the calculation of S_1 NTII TS energies are identified and diagnosed.

Benchmarking S_1 NTII intrinsic reaction barriers

Three of the S_1 NTII intrinsic reaction barriers in Table 1 are negative, those for: hex-2-enal, 2-methylpent-2-enal, and 3-methylpent-2-enal. These reactions may indeed be barrierless, or the TD-B2GP-PLYP NTII TSs on S_1 may have a different geometry to those optimised with TD-CAM-B3LYP. A third option is the presence of an S_1/S_0 CI may degrade the quality of the TD-DFT energies. There is therefore a need to benchmark the S_1 NTII intrinsic reaction barriers and characterise any CIs.

Unfortunately, there are no reliable experimental benchmarks for the S_1 NTII intrinsic reaction barriers; the ranges extracted from experiment are not sufficiently precise. There are a few previous literature calculations of these S_1 intrinsic reaction barriers, ^{5,26,54} and these are included in Table S1, where they are compared to calculations with four linear response methods (TD-CAM-B3LYP, TD-M06-2X, TD-B2GP-PLYP and EOM-CCSD) on all carbonyl species in this paper. In short, large energetic discrepancies are seen between calculations using these different methods.

The multireference XMCQDPT2 calculations of Kletskii *et al.*, which are based on CASSCF geometries, predicted lower barriers than the TD-M06-2X calculations of Shaw, as does the CIS(D) barrier predicted for butanal by Tadic *et al.* The TD-M06-2X/6-311++G(3df,2p) calculations by Shaw predicted S_1 NTII intrinsic reaction barriers approximately 10–20 kJ/mol higher than CASSCF and CASPT2 values,²⁶ however it was shown that these higher intrinsic TD-M06-2X barriers are more consistent with experimental photolysis thresholds. Importantly, from the comparisons in Table S1 predicted S_1 intrinsic reaction barriers are sensitive to the level of theory used.

Moreover, Kletskii *et al.* find that, at the XM-CQDPT2 level of theory, the S_1 NTII reaction in but-2-enal is barrierless, that is, a negative energy is reported in Table S1. A number of S_1 NTII reactions, for enals and enones, are also found to be barrierless at the levels of theory used here. Literature calculations on carbonyl S_1 NTII intrinsic reaction barriers using wavefunction methods have predicted low (<15 kJ/mol) barriers, ^{5,54} or even barrierless processes, for α, β -unsaturated species like but-2-enal.⁵⁴

A possible reason for the large energetic discrepancies between methods could be the breakdown of the linear response approximation at either, or both, the S_1 minimum energy geometry and the S_1 NTII TS. Single point energies at S_1 minimum energy geometries are reported in Table S2. These are broadly consistent for all four levels of theory, and moreover are consistent with experimental 0–0 excitation energies. Therefore, the discrepancies arise from calculations on the S_1 NTII TS, and this is shown by anomalously low $S_0 \rightarrow S_1$ single point energies at S_1 NTII TS geometries reported in Table S3.

Comparison between S_1 and T_1 B2GP-PLYP energies

These large variations and/or oscillations in energy are diagnostic of the presence of a MECI near the TS.^{55,72–74} As an additional 'flag' for the presence of an S_1/S_0 CI, the S_1 NTII intrinsic reaction barriers were compared to the corresponding T_1 energies. For both S_1 and T_1 , the NTII reaction takes place via a species with biradical character where the radical centres are separated by two carbon atoms. As a result, the electronic character of the NTII TSs should be similar on both electronic states, and should display similar energetic trends with respect to substitution for the saturated carbonyls considered. This is not necessarily true for the enals and enones, however, as the double-bond lies between the radical centres in the α , β -unsaturated species in this dataset. Furthermore, unlike the S_1 energies, the optimised T_1 TS energies are expected to be considerably more reliable because the T_2 state is significantly higher in energy, and a T_2/T_1 CI is not expected. The T_1 UB2GP-PLYP calculations are expected to yield accurate and reliable energies as they have been validated against experiment.⁴⁹ Our calculated NTII intrinsic reaction barriers on S_1 (open symbols) and T_1 (closed symbols) are compared in Figure 2 below.

The S_1 TD-B2GP-PLYP//TD-CAM-B3LYP

energies in Figure 2 follow, for the most part, the same energetic trends as the T_1 UB2GP-PLYP energies. They are systematically 15–30 kJ/mol lower than the corresponding T_1 energies. However, the trends in NTII intrinsic reaction barriers with substitution patterns are not calculated to be in quantitative agreement for the ketones on S_1 and T_1 . The intrinsic T_1 NTII barrier is much higher than the S_1 barrier because, as in the footnote of Table 1, the T_1 minimum energy geometry of this molecule is calculated to be twisted around the C=C bond, whereas the S_0 and S_1 minima are planar. Additionally, the S_1 NTII intrinsic reaction barrier for 2-oxobutanal is higher than the T_1 energy, but this may be due to the different electronic nature of α -dicarbonyls.^{75,76} Thus, comparison of the S_1 TD-B2GP-PLYP and T_1 UB2GP-PLYP energies in Figure 2 does not indicate, in itself, a failure of the TD-DFT method.

In contrast, the EOM-CCSD//TD-CAM-B3LYP S_1 NTII intrinsic reaction barriers shown in Figure 2 (open markers with internal crosses) are inconsistent with the T_1 energies. The butanal, 3-methylpentanal, 3,3-dimethylbutanal, hexan-2-one, hexan-3-one, hex-1-en-3-one, and 2-oxobutanal EOM-CCSD//TD-CAM-B3LYP S_1 intrinsic reaction barriers are higher than, or close to, the corresponding T_1 barrier. The EOM-CCSD S_1 energy for 2-oxobutanal, however, is similar to the TD-B2GP-PLYP S_1 energy. For the other six carbonyls, the discrepancy may be indicative of a breakdown of the linear response theory in the vicinity of a CI. The only caveat is that EOM-CC energies are extremely sensitive to geometry.⁷⁷ To investigate this further, S_1/S_0 MECI searches were performed on select carbonyl species, and a range of diagnostics determined.

Further diagnostics for S_1 calculations

Diagnostic calculations were performed on the TD-CAM-B3LYP S_1 NTII TS geometries, as well as the S_1 minimum energy geometries. This includes the \mathcal{T}_1 CCSD diagnostic for multireference character,^{78,79} as well as $S_1 - S_0$ and $S_2 - S_1$ electronic state energy gaps. These diagnostic calculations are reported in Table S4.



Figure 2: Intrinsic Norrish Type II reaction barriers, relative to relevant electronic state minimum for each molecule. Filled markers (\bullet) are T_1 intrinsic barriers, calculated at the UB2GP-PLYP/def2-TZVP level of theory. Open markers are S_1 barriers. Completely unfilled (\bigcirc) markers show intrinsic S_1 barriers calculated at the TD-B2GP-PLYP/def2-TZVP//TD-CAM-B3LYP/def2-TZVP level of theory, while open markers with internal crosses (\oplus) show intrinsic S_1 barriers calculated at the EOM-CCSD/def2-TZVP//TD-CAM-B3LYP/def2-TZVP level of theory. All energies are ZPVE corrected. Symbol shape and colour indicate carbonyl classes: linear aldehydes (\bullet), branched aldehydes (\bullet), linear ketones (\bullet), enones (\blacktriangledown), and dicarbonyls (\blacklozenge).

The \mathcal{T}_1 diagnostics in Table S4 for the TD-CAM-B3LYP S_1 minimum energy structures are all below the 0.02 threshold considered acceptable for closed-shell singlet calculations.⁷⁸ For S_1 NTII TS geometries several \mathcal{T}_1 diagnostic values were at or above the 0.02 threshold: pentanal (0.055), hex-3-enal (0.041), 4-methylpentan-2one (0.027), hex-4-en-3-one (0.022), and but-2enal (0.020). In all of these species the $S_1 - S_0$ energy gaps are < 100 kJ/mol, as determined by 10-root TD-CAM-B3LYP single point energy calculations. The $S_1 - S_0$ separations for but-2-enal, hex-3-enal, and hex-4-en-3-one are < 30 kJ/mol. These energy gaps are lower than observed for species with acceptable \mathcal{T}_1 diagnostics, with the exception of pentan-2-one which has a low (65 kJ/mol) energy gap without a \mathcal{T}_1 diagnostic above 0.02 for the S_1 NTII TS.

None of the S_1 NTII intrinsic reaction barriers for these species with poor \mathcal{T}_1 values were identified, from Figure 2, as being potentially problematic. Indeed, there is also no indication, on the \mathcal{T}_1 diagnostics presented in Table S4, that the S_1 calculations for butanal, 3-methylpentanal, hexan-2-one, hexan-3-one, hex-1-en-3-one, and 2-oxobutanal are problematic, despite their inconsistent EOM-CCSD energies in Figure 2. It appears a poor \mathcal{T}_1 diagnostic is not necessarily an indication that the S_1 NTII energy calculations by TD-DFT or EOM-CCSD may be unreliable. The converse also does not hold, an acceptable \mathcal{T}_1 diagnostic cannot be taken as assurance that the S_1 energy is reliable. All linear response S_1 calculations for the NTII reaction should therefore be treated with caution.

S_1/S_0 conical intersections

MECI searches were performed on species with large \mathcal{T}_1 diagnostic values, as well as for butanal, hexanal, and 2-oxobutanal. This ensured that all classes of carbonyl were considered, to determine whether an S_1/S_0 MECI in the vicinity of the S_1 NTII reaction coordinate is a feature common across all carbonyl classes. While the energetics of CASSCF will be poor due to lack of dynamic electron correlation, the geometries should be reliable as CASSCF bond lengths are found to differ by $\sim 0.01-0.02$ Å com-

pared to CASPT2 geometries for small organic molecules, including saturated and unsaturated carbonyls.⁸⁰

The optimised S_1/S_0 MECIs for 4methylpentan-2-one and 2-oxobutanal, are overlaid on their corresponding S_1 NTII TS structures in Figure 3. These two examples are representative of all species, and similar figures for all other MECIs are available in Figure S3 of the supporting information. The striking similarity of the S_1/S_0 MECIs for all carbonyls considered suggests similar S_1/S_0 MECIs will be present for all carbonyls that bear a γ -hydrogen. This mechanism is therefore likely to be important in all tropospheric NTII photolysis reactions.



(a) 4-Methylpentan-2-one

Figure 3: Structures of the S_1/S_0 MECIs (coloured) overlaid on the S_1 NTII TS structures (*black silhouette*). MECIs are calculated at the CAS(10,8)/6-31G(d) level of theory, while S_1 NTII TS structures are from CAM-B3LYP/def2-TZVP calculations.

For all carbonyls considered, the S_1/S_0 MECI has a very similar structure to the S_1 NTII TS, and is on the 1,4-biradical product side of the TS. This is the reaction coordinate for S_1 NTII reactions. These calculations therefore support the generic S_1 NTII mechanism shown in Figure 1, where crossing to S_0 occurs before the biradical intermediate is formed, implying short S_1 lifetimes.

Despite the presence of S_1/S_0 MECIs in the vicinity of S_1 TS, and the observation of problematic \mathcal{T}_1 diagnostics and small $S_1 - S_0$ energy gaps, no clear approach is apparent for determining whether the calculated S_1 NTII photolysis thresholds are reliable. Perhaps the best available indication is the similarity of calculated TD-B2GP-PLYP S_1 energies to the trends observed for the T_1 intrinsic reaction barriers, as was shown in Figure 2. In terms of literature precedent, there are suggestions that time-dependent

double-hybrid DFT methods are more reliable in the vicinity of CIs for acetophenone,⁸¹ due to the treatment of double excitations,^{58,82} and this is broadly supported in Figure 2. Therefore TD-B2GP-PLYP S_1 NTII intrinsic reaction barriers were reported in Table 1, and are used below in the discussion of the NTII reaction on different electronic states.

Discussion

The S_1 NTII reaction

The S_1 NTII photolysis threshold energies from Table 1 are shown in Figure 4, alongside a dashed blue line denoting the 313 nm photolysis energy used in NTII triplet quenching studies of butanal, pentan-2-one, and hexan-2-one.^{6,7} A solid purple line at 400 kJ/mol is used to indicate an approximate maximum photon energy in the troposphere in the figures below. As discussed previously, the S_1 photolysis thresholds in Figure 4 are consistent with experimental singlet/triplet ratios.



Figure 4: Photolysis thresholds for the S_1 NTII reaction, calculated at the TD-B2GP-PLYP//TD-CAM-B3LYP + CC/DFT'' level of theory (see text).

The calculated S_1 NTII thresholds of ketones are ~10 kJ/mol higher than the equivalent aldehyde species, reflective of higher $S_0 \rightarrow S_1 \ 0-0$ excitation energies since the intrinsic barriers are similar. The α, β -unsaturated carbonyls are predicted to have the lowest S_1 1,5–H-shift photolysis thresholds, due to their low 0–0 excitation energies. The enals are calculated to have slightly higher S_1 NTII thresholds than the enones, again due to differences in $S_0 \rightarrow S_1$ 0–0 excitation energies, since the S_1 intrinsic reaction barriers calculated for enals are low, or even barrierless. Because hex-3-enal is β , γ unsaturated there is no conjugation with the C=O chromophore in this molecule, and hence its excitation energies and photolysis thresholds resemble those of the saturated carbonyls. 2-Oxobutanal is predicted to have a low (294 kJ/mol) S_1 NTII photolysis threshold due to low excitation energies in α -dicarbonyls.

From Figure 4, the S_1 NTII photolysis thresholds of the saturated ketones are within ~ 20 kJ/mol of the maximum photon energy available in the troposphere ($\sim 400 \text{ kJ/mol}$). Since the intensity of solar radiation increases with wavelength, this means that S_1 NTII reactions of saturated ketones in the troposphere are likely to be slow, with small QYs. In saturated carbonyls, ISC to T_1 and reaction on T_1 will be competitive, if not the dominant photolysis pathway. Exceptions occur for saturated carbonyls whose S_1 NTII photolysis thresholds are lowered by substitution. The SARs dictating the lowering of photolysis thresholds with substitution patterns will be elucidated after discussion of NTII photolysis thresholds calculated on each electronic state.

The T_1 NTII reaction



Figure 5: Photolysis thresholds of the T_1 NTII reaction, calculated at the UB2GP-PLYP/def2-TZVP level of theory.

The calculated T_1 NTII photolysis thresholds are shown in Figure 5. Comparison of Figures 4 and 5 shows the same general relationships within and between carbonyl classes for both S_1 and T_1 NTII thresholds. However, all T_1 NTII thresholds are predicted to be lower than 380 kJ/mol (315 nm), meaning that the T_1 NTII reaction pathways should be accessible in the previous 313 nm or 315 nm photolysis experiments on butanal, pentanal, pentan-2-one, and hexan-2-one.^{6,7,32}

In principle, all T_1 NTII reactions are accessible under tropospheric conditions for all carbonyls considered here. Moreover, even at the lowest actinic energies at which saturated carbonyls absorb (for aldehydes this is ~ 360 kJ/mol,⁸³ the T_1 NTII photolysis channel is open for many carbonyls. This will lead to relatively large T_1 NTII QYs within the actinic range for most large carbonyls. This will also make experimental measurement of T_1 NTII photolysis thresholds in aldehydes difficult, as the UV absorption energies that can be accessed in an experiment will be above the T_1 NTII threshold. Similarly, the predicted T_1 NTII photolysis threshold of 2-oxobutanal, at 148 kJ/mol, will prevent a precise T_1 NTII photolysis threshold measurement.

The ketones are predicted to have T_1 NTII photolysis thresholds ~10 kJ/mol higher than equivalent aldehydes, and so their T_1 thresholds are above 360 kJ/mol (330 nm). Therefore, photolysis experiments on ketones with lower energy wavelengths, ~330 nm or longer, should show the photolysis behaviour of these species when the T_1 NTII reaction is inaccessible.

The enals and enones are predicted to have low T_1 NTII photolysis thresholds, at or below 330 kJ/mol. The α , β -unsaturated carbonyls have broader absorption cross-sections than saturated species — the longest wavelength at which typical enals and enones absorb is ~400 nm (~300 kJ/mol).⁸³ Although the T_1 NTII photolysis thresholds are energetically accessible in α , β -unsaturated carbonyls, as will be discussed below, the second step of the NTII reaction is not possible in these species.

The NTII reaction in α , β -unsaturated carbonyls

Of the enals, the photolysis of but-2-enal (crotonaldehyde) is the most experimentally studied.^{84–89} None of these studies, however, have detected the enol and alkene photoproducts which would be expected from the NTII reaction, despite the low predicted S_1 and T_1 NTII photolysis thresholds shown in Figures 4 and 5. This can be rationalised by the inability of α, β unsaturated species to cleave their α, β -bond after formation of the 1,4-biradical. The photolysis threshold energies shown in Figures 4 and 5 are for the initial 1,5–H-shift, but, as shown in Scheme 4, the α, β -double-bond prevents cleavage from occurring following electron migration to stabilise the γ -C[•] and [•]C–O radical centres.



Scheme 4: α, β -unsaturation prevents Norrish Type II bond cleavage to an alkene and an enol.

Instead, photoisomerisation dominates the photolysis of but-2-enal.^{85,88,89} Based on the inability of α , β -unsaturated species to undergo NTII reaction, and the high T_1 NTI*a* photolysis thresholds calculated for unsaturated species in earlier work,⁴⁹ ground state reactions and isomerisations are expected to dominate the photolysis of all α , β -unsaturated species.

The S_0 concerted NTII reaction

There have been proposals of an alternative S_0 concerted NTII reaction where the 6-membered cyclic TS leads directly to alkene and enol formation in one step,¹⁶ as shown in Figure 6a. Analogous reaction mechanisms are well documented in organic chemistry in the class of 'retro-ene' pericyclic signatropic reactions^{90,91} (Figure 6c), and the McLafferty rearrangement in mass spectrometry⁹² (Figure 6b).

Complete retention of stereochemistry should be seen in concerted S_0 NTII reaction products, by analogy to suprafacial migration in [1,5]sigmatropic rearrangements.⁹⁰ However, stereochemistry alone is not sufficient to disambiguate

$ \overset{R}{\longrightarrow} \overset{H}{\underset{R'}{\overset{O^*}{\longrightarrow}}} \overset{R}{\longrightarrow} \overset{H}{\underset{P'}{\overset{H}{\longrightarrow}}} \overset{O}{\underset{R'}{\overset{H}{\longrightarrow}}} $	$\begin{bmatrix} R \bigcup_{i=1}^{H} O \\ \vdots \bigcup_{R'} \end{bmatrix}^{t} \xrightarrow{R} \mathbb{I} + \begin{bmatrix} H O \\ \vdots \bigcup_{R'} \end{bmatrix}^{t}$	$ \overset{R}{\underset{R'}} \overset{H}{\underset{R'}} \overset{R}{\underset{R'}} \overset{R}{\underset{R'}} \overset{H}{\underset{R'}} \overset{R}{\underset{R'}} \overset{H}{\underset{R'}} \overset{R}{\underset{R'}} \overset{H}{\underset{R'}} \overset{R}{\underset{R'}} \overset{R}{\underset{R'}} \overset{R}{\underset{R'}} \overset{R}{\underset{R'}} \overset{R}{\underset{R'}} \overset{R}{\underset{R'}} \overset{R}{\underset{R'}} \overset{R}{\underset{R'}} \overset{R}{\underset{R''}} \overset{R}{\underset{R'''}} \overset{R}{R'''''''''''''''''''''''''''''''''''$
(a) Concerted	(b) McLafferty	(c) Retro-ene

Figure 6: The reaction schemes for the proposed concerted S_0 NTII reaction (a), alongside the McLafferty rearrangement (b), and the retro-ene reaction (c).

between S_1 or S_0 NTII reactions, since the S_1/S_0 conical intersection promotes cleavage following an S_1 1,5–H-shift before stereochemical scrambling can occur. Evidence for S_0 NTII reactions therefore relies on energetic arguments, reinforcing the importance of accurate prediction of S_1 , T_1 , and S_0 NTII photolysis threshold energies.

We calculate the reaction thresholds for the S_0 NTII reaction, which is a concerted reaction where the intramolecular hydrogen transferred is accompanied by α , β -bond breaking in the same step. These S_0 reaction thresholds are shown in Figure 7, and have been labelled as 'photolysis thresholds' to emphasise that, although these energies are S_0 intrinsic reaction barriers, S_0 NTII reactions under tropospheric conditions would occur by absorption of actinic photons. Concerted S_0 NTII TSs could not be located for most enals and enones: the double-bonds at the α , β -position in these species prevent α , β -bond cleavage from accompanying the 1,5–H-shift.



Figure 7: Photolysis thresholds for the concerted S_0 NTII reaction, calculated at the B2GP-PLYP/def2-TZVP level of theory.

The S_0 NTII photolysis thresholds shown in Figure 7 are all well below the longest wavelength range where saturated carbonyls absorb (~330 nm / 362.5 kJ/mol).⁸³ If formation of NTII photoproducts in pentan-2-one in recent

330 nm FT-IR experiments by Kharazmi is evidence for the existence of a S_0 NTII pathway,³² then the thresholds shown in Figure 7 indicate that energetically accessible S_0 NTII reaction pathways should be available for all saturated carbonyl species. Hex-3-enal is predicted to have the highest concerted S_0 NTII photolysis threshold, since the γ -H transferred is adjacent to the β , γ -double-bond and so the energy of the closed-shell TS to transfer the γ -H is raised. Apart from hex-3-enal, the concerted S_0 NTII threshold energies of all other carbonyls are similar, within an $\sim 225-245$ kJ/mol range.

A master equation scheme for butanal photolysis proposed by Shaw²⁶ included the concerted S_0 NTII pathway, but this pathway was found to be inconsequential at 290–330 nm photolysis wavelengths. So, while the S_0 NTII pathway should be energetically accessible in butanal, it is predicted to be kinetically unimportant where S_1 and T_1 channels are open. In contrast, recent 330 nm pentan-2-one photolysis data by Kharazmi indicates that an S_0 NTII reaction is occurring in pentan-2-one.³² As seen in Figure 7, S_0 NTII photolysis threshold energies do not explain the different photochemical behaviour of butanal and pentan-2-one, since the thresholds for both molecules are significantly lower than the photolysis energies used.

Alternatively, the predicted T_1 NTII threshold for pentan-2-one could be wrong, leading to miss-assignment of the electronic state involved, though a 15 kJ/mol overestimate by the UB2GP-PLYP/def2-TZVP calculations seems outside of the accuracy demonstrated when benchmarking to experimental T_1 thresholds,⁴⁹ and is in good agreement with other estimates of the T_1 thresholds.³² If another photolysis pathway apart from the NTII reaction could form ethene, it could explain the QYs at 330 nm, however our own calculations exploring other possible ground state reactions do not present a viable alternative pathway.

The master equation calculations of Shaw on butanal indicate that, from 310–330 nm, NTI*a* reactions on either T_1 or S_0 dominate. It is possible that the S_0 NTII reaction is observed for pentan-2-one and not butanal because the T_1 NTI*a* and NTII photolysis thresholds are higher for pentan-2-one than for butanal,⁴⁹ and so reaction on T_1 is closed at longer wavelengths for pentan-2-one, allowing significant ISC to S_0 and an appreciable amount of S_0 NTII reactions.

SARs for Excited State NTII Reactions

The SARs elucidated in the following subsections are derived from T_1 NTII photolysis energies. Since both the S_1 and T_1 states arise from $n \to \pi^*$ excitations, and the S_1 and T_1 carbonyl biradicals have similar structures and electron distributions, these SARs should also apply to S_1 NTII photolysis. They are not, however, applicable to concerted S_0 NTII reactions. Since the S_0 NTII photolysis thresholds are predicted to be below actinic energies, a simple SAR for S_0 NTII pathways is to assume that they are accessible in the troposphere for all but α, β -unsaturated carbonyl species.

Substitution at the γ -position

Substitution at the γ -position is the clearest SAR for excited state NTII processes, leading to a significant decrease in excited state photolysis thresholds. For instance, upon moving from but anal with a primary γ -carbon, to pentanal with a secondary γ -carbon, the T_1 NTII threshold drops by ~10 kJ/mol. This γ -substitution effect in carbonyl NTII photolysis is consistent with triplet quenching experiments: at 313 nm pentan-2-one NTII photolysis is a triplet process, whereas it is a singlet process in hexan-2one⁶ since the extra γ -substituent lowers the S_1 NTII threshold to become energetically accessible. This stabilisation is due to hyperconjugation of the incipient 1,4-biradical at the position of γ -H abstraction, as illustrated in Figure 8. Similar to the relatively *late* transition states in the NTIa and NTIb reactions,⁴⁹ NTII photolysis thresholds are also stabilised through hyperconjugation.

Further chain extension, *e.g.* moving from pentanal to hexanal, does not further decrease the T_1 NTII threshold. However additional substitution, such as the tertiary γ -carbon centre in



(a) Hyperconjugation at (b) 4-methylpentanal T_1 the γ -position. NTII TS.

Figure 8: Substitution at the γ -position directly stabilises the incipient 1,4-biradical in the NTII reaction. This is illustrated in the Lewis structure in (a), and shown in (b) through the unpaired electron density of the T_1 NTII TS calculated for 4-methylpentanal.

4-methylpentanal, drops the threshold by ${\sim}20$ kJ/mol compared to the primary $\gamma\text{-carbon}$ in butanal.

A simple SAR for excited state NTII reactions is that the photolysis threshold will decrease by ~10 kJ/mol per extra degree of substitution at the γ -position. This effect is seen across the other carbonyl classes in Figure 5, including the ketones (*cf.* pentan-2-one and hexan-2-one), the enals (*cf.* but-2-enal and hex-2-enal), and the enones (*cf.* hex-4-en-3-one and hex-3-en-2-one).

Substitution at carbonyl, $\alpha,$ and β positions

The incipient radical position in the $^{\circ}C-O$ moiety of the forming biradical could also be stabilised by hyperconjugation. Alkyl substitution at the carbonyl-carbon should provide adjacent $C-H \sigma$ -bonds for hyperconjugation. However, the calculated NTII intrinsic reaction barriers for ketones in Table 1 are typically raised compared to equivalent aldehydes, rather than lowered. This indicates that any potential hyperconjugative stabilisation gained by the ketone alkyl group is offset by destabilisation of the excited C=O chromophore by the alkyl auxochrome.

For substitution at the α -position, T_1 thresholds in Figure 5 do not indicate a clear effect. The T_1 NTII photolysis threshold of 2-methylpentanal is predicted to be lowered by 4 kJ/mol compared to α -unsubstituted pentanal. Similarly, the T_1 threshold for 3-methylpentan-2-one is lowered 7–9 kJ/mol compared to α -unsubstituted ketones. However,

the lack of NTII threshold energy lowering in the T_1 intrinsic NTII barrier of doubly- α substituted 3,3-dimethylbutanal (though threshold lowering is seen on S_1) indicates that any α -substitution effect is not as straightforward as the γ -substitution effect

Substitution at the α -position should deliver no hyperconjugative stabilisation of the C[•] center if it was strictly localised as in a simple Lewis structure (Figure 9a). However the α -carbon does show a small amount of spin density delocalisation at a low isosurface value of 0.005 (Figure 9b), which could confer some of this T_1 NTII barrier lowering predicted for α -substituted species.



(a) Lewis structure with lo- (b) 2-methylpentanal T_1 calised radical centres. NTII TS.

Figure 9: Substitution at the α -position may indirectly stabilise unpaired $^{\circ}C-O$ electron density in the incipient 1,4-biradical in the NTII reaction. This is not shown by the Lewis structure of 2-methylpentanal in (a), but a small indication can be seen in the unpaired spin density of 2-methylpentanal (b) at an isosurface value of 0.005.

Energy differences of ~4 kJ/mol are however at the limit of chemical accuracy, and so using such differences to infer SARs is tentative. Additionally, given the limited number of β substituted species shown in Figure 5, there is insufficient data to draw a SAR based on substitution at the β -position. It would be expected that β -substitution affects NTII thresholds similarly to α -substitution, since both place an alkyl chain two carbons away from a biradical centre.

These SARs hold for alkyl substituents, and not substituents which are strongly electron withdrawing or donating.¹⁶ While substituents at the α or β position do not strongly alter energetic thresholds, they do affect the preferred conformation of the resulting biradical and hence influence the ratio of cyclisation vs. cleavage products of the NTII reaction.^{16,93,94} For aryl carbonyls, substitution at α and β positions of aromatic rings can alter NTII reactivity, as was observed in photolysis studies of substituted valerophenones.⁹³ Since aryl carbonyls are not common in the troposphere, they are excluded from this study and SARs pertaining to aromatic carbonyls are not determined here.

α, β -Unsaturation and allylic radical resonance.

 α, β -Unsaturated species display lower S_1 and T_1 intrinsic reaction barriers in Table 1 than the saturated carbonyls. This can be rationalised by considering allylic radical resonance stabilisation.⁹⁵ From the perspective of Lewis structures, unsaturation near a radical site allows multiple resonance contributors, as shown in Figure 10a. This allylic radical stabilisation manifests as delocalisation of the unpaired spin density, shown for the example of the hex-4-en-3-one T_1 NTII TS in Figure 10b. In Figure 10b, the spin density has delocalised to the β -position; this is not seen in saturated carbonyls.



(a) Allylic radical resonance stabilisation.

(b) Hex-4-en-3-one T_1 NTII TS.

Figure 10: Unsaturated near incipient radical sites formed in the NTII reaction will lower reaction barriers through allylic resonance stabilisation. Lewis structure resonance contributors are illustrated in (a), whereas the spin density of the hex-4-en-3-one T_1 NTII TS shown in (b) demonstrates that the unpaired electron density delocalises over to the β -position.

The lowering of excited state NTII intrinsic reaction barriers by allylic radical resonance stabilisation in α , β -unsaturated carbonyls is significant, approximately 15–20 kJ/mol lower compared to corresponding saturated carbonyls. For instance, the intrinsic T_1 NTII barrier for but-2-enal is 18 kJ/mol lower than for butanal, in hex-3-en-2-one the barrier is 14 kJ/mol lower compared to hexan-2-one, and in hex-2-enal the barrier is 23 kJ/mol lower than in hexanal. A simple SAR to account for this radical allylic resonance is to posit that α , β -unsaturation leads to 15–20 kJ/mol lower intrinsic excited state NTII barriers, compared to the corresponding saturated carbonyl. Combined with the SAR of an approximately 30–50 kJ/mol 0–0 excitation energy decrease in α , β -unsaturated carbonyls, this leads to excited state NTII 1,5–Hshift photolysis thresholds in α , β -unsaturated carbonyls ~50 kJ/mol lower than the equivalent saturated species.

Barrierless, or small barrier, excited state NTII reactions

The intrinsic reaction barriers for the excited state NTII 1,5–H-shift photolysis thresholds in Table 1 are mostly under 30 kJ/mol on S_1 and under 50 kJ/mol on T_1 . If hex-3-enal is removed, because its β , γ -unsaturation raises abstraction barriers, the least substituted saturated carbonyls, such as butanal, 3,3-dimethylbutanal, and hexan-3-one, have the highest intrinsic reaction barriers. As above, α , β -unsaturation is predicted to lower NTII intrinsic reaction barriers by $\sim 15-20$ kJ/mol, and barriers should decrease by $\sim 10 \text{ kJ/mol}$ for each extra degree of γ -substitution. Since the S_1 NTII intrinsic barriers for butanal and pentan-2-one are predicted to be ~ 30 kJ/mol, lowering of the photolysis energy due to the SARs above leads to predictions of small barriers, or even barrierless processes, for carbonyl species that are both α, β -unsaturated and substituted at the γ -position. For example, the butanal S_1 NTII intrinsic reaction barrier is predicted to be 29 kJ/mol, but the but-2-enal intrinsic reaction barrier is significantly smaller, $\sim 5 \text{ kJ/mol}$, due to the α , β -unsaturation. When combined with extra γ -substitution, the S_1 NTII intrinsic reaction barriers in enals become negative, *i.e.* are submerged barriers.⁹⁶ This includes: hex-2-enal, 2-methylpent-2-enal, and 3-methylpent-2-enal in Table 1. This means that as long as photoexcitation to the S_1 minimum is possible, the S_1 NTII 1,5–H-shift should always be energetically accessible in these species.

The XMCQDPT2//CASSCF multiconfigurational calculations by Kletskii *et al.* predict a negative S_1 intrinsic reaction barrier for but-2enal, while in Table 1 this barrier is predicted to be slightly positive. However, in larger enals, substitution at the γ -position should lower the intrinsic reaction barrier further and thus be predicted to be negative. Therefore, these calculations present a model that γ -substituted enals may have barrierless S_1 NTII 1,5–H-shifts, with small barriers for enones.

It would be informative to study the entire S_1 NTII reaction pathway of these species with accurate multiconfigurational methods, which has recently been made possible with the development of analytic CASPT2 nuclear gradients.⁹⁷ However, as noted before, these α, β -unsaturated species are unlikely to undergo NTII cleavage, but instead will react through ground state isomerisation or dissociation pathways.

Conclusions

The calculations in this paper predict the energetics of S_1 , T_1 , and S_0 NTII reaction pathways in 22 carbonyl species and confirm the general schematic introduced in Figure 1: when S_1 photolysis thresholds are energetically accessible due to γ -substitution then formation of a S_1 biradical is rapid, the biradical species then crosses to S_0 via an S_1/S_0 conical intersection and cleaves to form an S_0 alkene and an enol. When the S_1 NTII threshold is inaccessible, then ISC to T_1 and formation of a long-lived triplet biradical occurs. The triplet biradical can then cleave or cyclise. While α, β -unsaturated carbonyls are all predicted to have very low NTII thresholds for the 1,5–H-shift reaction because of allylic radical stabilisation, these species cannot cleave the α, β -bond. Instead, it is likely they cross to S_0 as internally 'hot' molecules. On S_0 , the major unimolecular reaction pathways of these species are expected to be ground state isomerisation and dissociation reactions. Concerted S_0 pathways to cleavage to an alkene and an enol are calculated for most saturated species, and could explain the observation of ethene production from pentan-2-one at 330 nm.³² This concerted S_0 NTII pathway is also predicted

to be energetically accessible for the other carbonyls considered here. The reason this pathway has not been experimentally observed in other carbonyl species may be because reaction on S_1 and T_1 dominates at the experimental photolysis energies considered.

The NTII reaction is important in atmospheric models since 5–6 carbon species are common biogenic VOCs,^{40,98} which are oxidised in the atmosphere into carbonyls that can undergo the NTII reaction. However, the treatment of the NTII reaction in atmospheric chemistry models is inconsistent. The NYC reaction is not included in the photolysis schemes of major atmospheric chemistry models. NYC products are observed as a minor channel of T_1 NTII photolysis of gas-phase carbonyls. More pressingly, the NTII QY estimated for carbonyls is different between the two major atmospheric models (the MCM^{31} and GEOS-Chem⁹⁹), and both models are inconsistent with the experimental data and the photolysis thresholds calculated here.

For example, the MCM uses a NTII rate measured for butanal as the 'surrogate' rate for the wider array of carbonyls which undergo NTII reaction.^{31,100} GEOS-Chem is less detailed: the NTII reaction is excluded entirely from the generic photolysis scheme of aldehydes containing more than three carbon atoms.⁹⁹

The results of this paper show that using a generic NTII rate constant from butanal is inappropriate. Butanal should set a lower limit for NTII reaction QYs in aldehydes, as the SARs determined here predict that more substituted carbonyls should have lower NTII photolysis thresholds. Indeed, in 290–330 nm photolysis experiments (observing HCO[•] yields) the NTII QY was inferred to range from 0.52–0.56 for butanal, but was 0.68–0.91 for pentanal.¹⁰¹

The calculations in this paper characterised the NTII photolysis thresholds on all relevant electronic states, and located S_1/S_0 MECIs, for carbonyl species representative of each carbonyl class. There was no clear diagnostic for the failure of linear response methods near the vicinity of a MECI, however S_1 energies computed with time-dependent double-hybrid methods were consistent with the substituent effects on T_1 energies. These calculations provide comprehensive insight into the energetics of the NTII reaction. However, rationalising experimental data on the prevalence of cyclisation, cleavage, or H-backtransfer from the 1,4-biradical requires dynamic modelling.^{8,12,35,94,102} More theoretical studies of the NTII biradical⁴² are required, and *ab initio* dynamics has only been performed for the NTII reaction in pentanal.^{43,44} Furthermore, in order to assess the competition between NTII and other photolysis reactions on S_1 , T_1 , and S_0 kinetic modelling is required. The reaction energies, TS structures, and S_1/S_0 MECIs reported in this paper can be used as the basis for master equation calculations for NTII reactions across many carbonyl species.

Acknowledgement This work was supported by the Australian Research Council (grant DP160101792). It was also supported by grants of computer time under the Merit Allocation Scheme on the NCI National Facility at the Australian National University, as well as computer time on the computational cluster Katana supported by the Faculty of Science, UNSW Australia, and the computational cluster Artemis supported by the Sydney Informatics Hub at the University of Sydney. KNR recognises an Australian Government Research Training Program (RTP) scholarship.

Supporting Information Available

The Supporting Information contains: comparison between linear response S_1 NTII intrinsic reaction barriers and S_1 NTII intrinsic reaction barriers published in the literature; details of the active spaces used in CASSCF calculations; $S_0 \rightarrow S_1$ excitation energies calculated on S_1 minima and transition states using a range of linear response methods; visual overlays of S_1/S_0 minimum energy conical intersections and S_1 NTII transition states; diagnostic \mathcal{T}_1 values and $S_2 - S_1$ excited state energy separations. (PDF)

Cartesian co-ordinates for all optimised structures (ZIP)

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