

Precise Control of π -electron Magnetism in Metal-free Porphyrins

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ABSTRACT

The porphyrin macrocycle can stabilize a set of magnetic metal ions, thus introducing localized net spins near the center. However, it remains illusive but most desirable to introduce delocalized spins in porphyrins with wide implications, for example, for building correlated quantum spins. Here, we demonstrate that metal-free porphyrins host delocalized π -electron magnetism as revealed by scanning probe microscopy and different level of theory calculations. Our results demonstrate that engineering of π -electron topologies introduces spin polarized singlet state and delocalized net spins in metal-free porphyrins. In addition, the π -electron magnetism can be switched on/off via STM manipulation by tuning the interfacial charge transfer. Our results provide an effective way to precisely control the π -electron magnetism in metal-free porphyrins, which can be further extended to design new magnetic functionalities of porphyrin-based architectures.

INTRODUCTION

Porphyrin molecules have a wide spectrum of physicochemical and biological properties with diversified applications in sensors, molecular electronics/spintronics and medical applications¹⁻³. The rich functionalities of porphyrin complexes originate from the center macrocycle, which can stabilize a variety of elements inside and thus gives rise to distinct structural, electronic, magnetic, optoelectronic, and transport properties⁴⁻¹³. Magnetic metalloporphyrins have been extensively studied by various techniques, exhibiting Kondo and spin excitation effects, Yu-Shiba bound states, and spin-state switching¹³⁻²⁷. Recently, hybridized graphene and metalloporphyrin nanostructures have been achieved, hosting combined intriguing electronic and magnetic properties²⁸⁻³³. The magnetism of metalloporphyrin architectures originates from the stabilized metal ions, and thus localizes at the center of the molecules. It remains elusive but most desirable to develop strategies to introduce delocalized spins in porphyrins for designer new magnetic functionalities in porphyrin architectures.

Recent advances of on-surface synthesis allow for tailoring π -electron topologies at the single chemical bond level, and thus provide the ability to precisely engineer π -electron magnetism in low-dimensional graphene nanostructures by introducing sublattice imbalance, topological frustration and topological defects³⁴⁻⁵⁰. Examples like triangulenes^{41,43,48,51}, bowtie-shape nanographenes⁴⁷, chiral graphene nanoribbons⁴⁵, polymers⁵² and nanographenes with topological defects^{46,49} have been recently synthesized and characterized on surfaces. Similar as nanographenes, the porphyrin macrocycle is aromatic and can be interpreted as a multiple-bridged aromatic diaza[18]annulene system^{53,54}. Thorough tailoring peripheral π -electron topology, it is in principle able to introduce tunable π -electron magnetism in metal-free porphyrin architectures, thus greatly enriching magnetic functionalities of porphyrin-based systems. Although substituted open-shell porphyrins have been realized using conventional solution synthesis methods, on-surface synthesis of unsubstituted porphyrins with π -electron magnetism remains to be explored^{55,56}.

Here, we demonstrate such ability to precisely control delocalized magnetism in metal-free porphyrins through tailoring their peripheral π -electron topologies. Using combined solution and on-surface synthesis, three metal-free porphyrins with different π -electron topologies have been synthesized. Their chemical structures have been resolved by using high resolution non-contact atomic force microscopy, permitting us to clearly distinguish their π -electron topologies at the single chemical bond level. Using low-temperature scanning tunneling spectroscopy, their π -electron magnetism has been confirmed by Kondo effects. The experimental observations can be nicely elucidated by clear empirical rule, mean-field Hubbard model and spin-polarized density functional theory calculations. In addition, the π -electron magnetism can be reversibly switched on/off via STM manipulation by

tuning the charge transfer from the underneath metal substrate. The realization of π -electron magnetism in metal-free porphyrin may be further extended to build correlated spins in porphyrin-based nanostructures, with implications for quantum information and spintronics applications.

RESULTS AND DISCUSSION

Using combined in-solution and on-surface synthesis, metal-free porphyrins with different π -electron topologies have been synthesized. As shown in Fig. 1a, the precursor 5,10,15,20-tetrakis(2,6-dimethylphenyl)porphyrin is synthesized in solution by using freshly distilled pyrrole and 2,6-dimethylbenzaldehyde (cf. detailed synthesis method in supplementary fig. 1). After in-solution synthesis, the precursor is thermally deposited on a Au(111) substrate held at room temperature following a subsequent annealing to 295 °C for 10 minutes. At elevated temperature, thermal-induced cyclodehydrogenation of methyl units occurs, giving rise to fully aromatic porphyrins^{39,57}. Except for the main cyclodehydrogenation, two common side reactions take place occasionally on Au(111): i) few methyl units detach from the precursors forming five-membered carbon rings after subsequent cyclodehydrogenation, and ii) partial dehydrogenation sometimes occurs and introduces SP_3 hybridized carbon sites at edges. As shown in Fig. 1b, three products have been frequently observed and characterized. The Product #1 contains two pentagon carbon-rings at opposite sides, and the Product #3 has a SP_3 hybridized carbon site passivated by two hydrogen atoms. Their chemical structures have been clearly resolved by high resolution nc-AFM imaging using a CO-functionalized tip (cf. Fig. 1c-e). The number of unpaired π -electrons (radicals) in the resulted porphyrins has been determined by using Clar's empirical rule (cf. all the possible π -electron networks in supplementary Fig. 3 and 4)^{58,59}. As shown in Fig. 1, the Product #1 is a close-shell molecule, while the Product #2 and #3 have two and one unpaired π -electron(s), respectively (as marked by red dots in Fig. 1). The presence of radicals suggests an open-shell configuration in the product #2 and #3, thus giving rise to delocalized π -electron magnetism.

The energy spectra and frontier molecular orbitals of the resulted porphyrins have been calculated by using mean-field Hubbard model (MFHM) and spin-polarized density functional theory (SP-DFT), which have been widely used to determine the magnetic structure of nanographenes and show nice agreements with experiments³⁷. As shown in Fig. 2a-c, the Product #1 has a close-shell ground state as confirmed by combined Clar empirical rule, MFHM and SP-DFT calculations. We notice that MFHM calculations qualitatively agree with SP-DFT, despite the exact energy position of each state varies a little bit from each other. Experimentally, differential conductance (dI/dV) spectroscopy has been used to detect the electronic structure of achieved porphyrins. Together with spatial mappings, dI/dV spectroscopy is able to resolve both the energy positions and spatial distribution of molecular

orbitals with ultra-high energy and spatial resolution. As shown in Fig. 2d, dI/dV spectra taken above the center of the Product #1 reveal two resonances at -1.4 V and 0.9 V, which can be assigned to the highest occupied molecular orbital (HOMO) and lowest unoccupied molecular orbital (LUMO), respectively. dI/dV mappings at these two energies resolve the spatial distribution of these two molecular states, which agree well with HOMO and LUMO orbital shapes as calculated by MFHM and SP-DFT. These results confirm that the Product #1 is a close-shell molecule without any unpaired π -electrons.

Despite its slightly different π -electron topology from that of the Product #1, the Product #2 has drastically different magnetic ground state. The Product #2 has in total 58 π -electrons, with 2 π -electrons more than those of the product #1. Although both of them host even number of π -electrons, the Product #2 has an open-shell electronic structure as opposed to the close-shell configuration of the Product #1 (cf. Fig. 3a). The reason is due to topological frustration, that is, it is impossible to concomitantly pair all p_z orbitals to form π bonds in the Product #2. This picture can be verified by both MFHM and SP-DFT calculations as shown in figure 3b-e. The calculated energy levels in Fig. 3b-c show that there are two degenerated states below fermi level, which are spatially separated with enhanced orbital density at opposite sides (cf. Fig. 3d and e), and singly occupied by two electrons with oppositely oriented spins. In other words, the ground electronic configuration breaks spin-spatial symmetry and exhibits antiferromagnetic ordering, which can be nicely seen in the calculated spin density maps in Fig. 3f and g. The mechanism behind is very similar to Clar Goblet, where topological frustration induces spin polarized singlet state in the bowtie-shape nanographene^{36,47}.

However, experimental dI/dV measurements reveal that the Product #2 hosts a net spin of $S=1/2$ by showing a sharp Kondo resonances instead of the expected singlet ground state with a net spin of $S=0$ (cf. Fig. 4g). As illustrated in Fig. 4a, we attribute this discrepancy to the interfacial charge transfer, where one electron is transferred from molecule to gold substrate making it positively charged. Similar charging effects have been previously observed in several graphene nanostructures, such as the zigzag termina of $N=7$ armchair graphene nanoribbons^{60,61}, and $N=6$ zigzag graphene nanoribbon³⁹. As shown in Fig. 4b-c, the electronic structure of the positively charged Product #2 is calculated by MFHM and SP-DFT, revealing a singly occupied molecular state below fermi level. This means the charged molecule is with a net magnetic moment of $S=1/2$. On metal substrate, this magnetic moment will be screened by Au(111) electron reservoir, which can be understood by single impurity Anderson Model^{17,62}. As shown in Fig. 4g, dI/dV spectra reveal such screening effect by showing a sharp Kondo resonance at Fermi level. In addition, the observed spin density distribution nicely agrees with calculated results, further supporting that the observed zero-bias peak originates from Kondo resonance (cf. Fig. 4e-f). As per

Anderson model, the Kondo peak width varies with temperature with full width at half maximum (FWHM) of $\Gamma = \sqrt{(\alpha k_B T)^2 + (2k_B T_K)^2}$, where T is the temperature, T_K is the Kondo temperature, and α is the slope of linear growth of the width at $T \gg T_K$. The experimentally obtained temperature dependent FWHM can be well fitted by the previous equation, giving a Kondo temperature of 22 K (cf. Fig 4h-i).

Using similar techniques, the magnetic properties of the Product #3 has been determined. Since one carbon site is passivated by two hydrogen atoms, the Product #3 has one π -electron less than the Product #2 with an odd number of 57 π -electrons in total. It is impossible to concomitantly pair odd number of p_z orbitals to form π bonds, thus generating one π radical as marked by a red dot in Fig. 5a. MFHM and SP-DFT calculations confirm the presence of a singly occupied state below Fermi level in the Product #2, that is, a net spin of $S=1/2$. As shown in Fig. 5d, the passivation of π -radical locally modifies its spin density distribution, which gets significantly quenched near the CH_2 site (cf. Clar non-kekulé structures in Supplementary Fig. 3). The emergence of net spins in porphyrins with odd number of π electrons is very similar to nanographenes with sublattice imbalance as proposed by E. Lieb in 1989³⁴. As shown in Fig. 5g, dI/dV spectroscopy measurements confirm the emergence of a net spin of $S=1/2$ in the product #3. The net spin couples with underneath electron reservoir, generating Kondo resonances by showing a sharp Fano-shape peak at Fermi level in dI/dV spectra with peak intensity proportional to spin density magnitude. In addition, the resolved spin density distribution fits very well with the calculated result (cf. Fig. 5e and f). The nice agreements between experiments and calculations in gas phase indicate that the Product #3 is neutral on Au(111).

Although the Product #3 shares very similar adsorption geometries with that of the Product #2, the former is neutral on Au(111) while the latter is positively charged. By carefully checking the AFM images in Fig. 1c-e, one can see that the presence of sp^3 carbon (CH_2) site slightly lifts the corner near the CH_2 site away from the surface, which is imaged as a local bright protrusion in nc-AFM imaging due to a stronger Pauli repulsion. This slight adsorption difference leads to a completely different charging behavior. Previous DFT calculations by using a set of van der Waals correction schemes reveal that despite very close adsorption heights of most graphene nanostructures on Au(111), the magnetism (charge transfer) is very sensitive to small perturbations (adsorption height variations), supporting our observations⁶³. To gain more evidences of the charging effect, we performed STM manipulation to locally modify the adsorption configurations. After each manipulation, we performed nc-AFM imaging, constant-height current imaging and dI/dV spectroscopy to trace the charging transfer. As shown in Fig. 6a, the porphyrin molecule is initially neutral by showing a Kondo resonance at Fermi level. After manipulation, the molecule rotates an angle of 185° around the out-of-plane axis (cf. Fig. 6b). Although nc-AFM and constant-height current image

indicate a similar adsorption geometry after rotation, dI/dV spectroscopy reveals that the molecule is positively charged by showing a broad peak above Fermi level instead of exhibiting a sharp Kondo resonance at Fermi level (cf. the No. 2 spectrum in Fig. 6d). In other words, the electron previously occupied the singly occupied state transfers to gold substrate, thus making the molecule positively charged and leading all orbitals doubly occupied. Further manipulation experiment was used to rotate the molecule again. As shown in Fig. 6c, the Kondo resonance appears after a successive rotation of 125° , suggesting the molecule becomes neutral again. In total, we performed successive 9 times STM manipulations, and found that the charging effect is extremely sensitive to adsorption variations (cf. supplementary Fig.8). Similar manipulation experiments have been performed on the Product #1 and #2 without obtaining this switching effect. These results indicate that slightly decoupling the molecule from surface can effectively quench the charging effect, such as by lifting the molecule slightly away from the substrate by functional groups.

CONCLUSION

In conclusion, we demonstrate an effective approach to engineer π -electron magnetism in metal-free porphyrins by using combined non-contact atomic force microscopy, scanning tunneling microscopy/spectroscopy as well as theory calculations at different levels. Three types of porphyrins with different π -electron topologies have been successively synthesized and characterized down to the single chemical bond level. The π -electron network in porphyrins with even number of π electrons can either host a close-shell configuration or spin polarized singlet ground state depending whether the network gets frustrated or not. Porphyrins with odd number of π electrons host net spin similar as sublattice imbalance induced spins in nanographenes. In addition, the magnetism can be switched on/off by tuning charge transfer from underneath Au(111). Our method reported herein provides ample opportunities for further designer correlated spins in porphyrin-based nanostructures to study their many-body effects.

ASSOCIATED CONTENT

Supporting Information. Detailed descriptions of experimental and theoretical procedures, Clar non-kekulé structures of the product #2 and 3, Magnetic field dependence of Kondo resonance, dI/dV spectra and mappings, and simulated local density of state maps. This material is available free of charge via the Internet at <http://pubs.acs.org>.

AUTHOR INFORMATION

Yan Zhao and Kaiyue Jiang contributed equally to this work.

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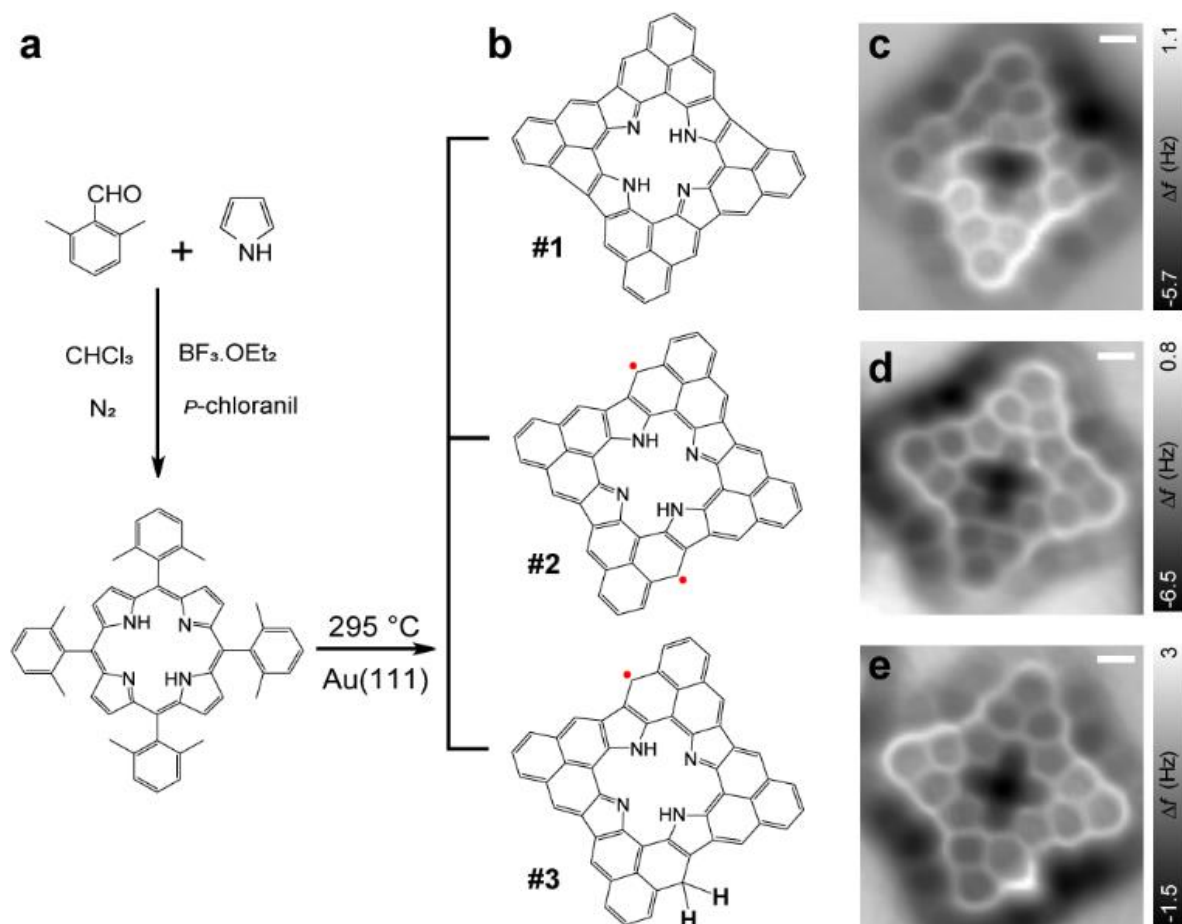


Figure 1. Synthesis of porphyrins with different π -electron topologies. (a), In-solution synthesis of molecular precursor 5,10,15,20-tetrakis(2,6-dimethylphenyl)porphyrin. (b), Chemical structures of three On-surface synthesized porphyrins with red dots indicating the unpaired π -electrons. (c-e), Nc-AFM frequency shift images (Resonant frequency: 26 KHz, Oscillation amplitude: 80 pm, Scale bars: 0.4 nm) of the porphyrins in (b).

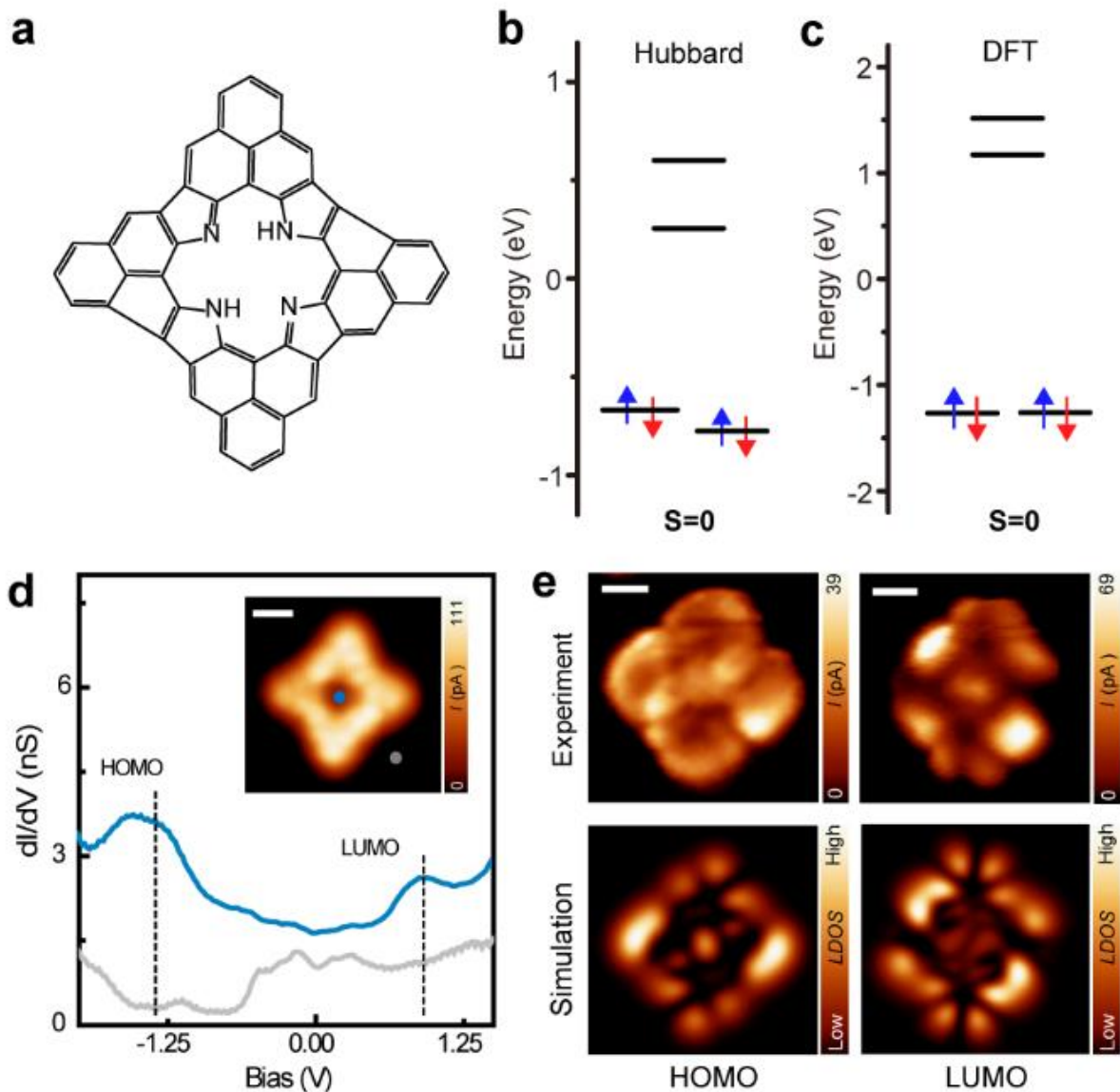


Figure 2. Electronic structure of a close-shell porphyrin with two embedded pentagon rings. (a), Chemical structure of the product #1. (b-c), Mean-field Hubbard model and spin-polarized density functional theory calculated energy spectrum of the porphyrin in (a). (d), dI/dV spectra taken at the two locations marked on the inset current image (scale bar: 0.25 nm). Two resonances have been observed and assigned to the highest occupied molecular orbitals (HOMO) and lowest unoccupied molecular orbitals (LUMO). (e), Up: spatial resolved local density of states maps taken at the biases marked by dashed lines in (d); Down: Simulated density of states maps. Scale bars: 0.5 nm.

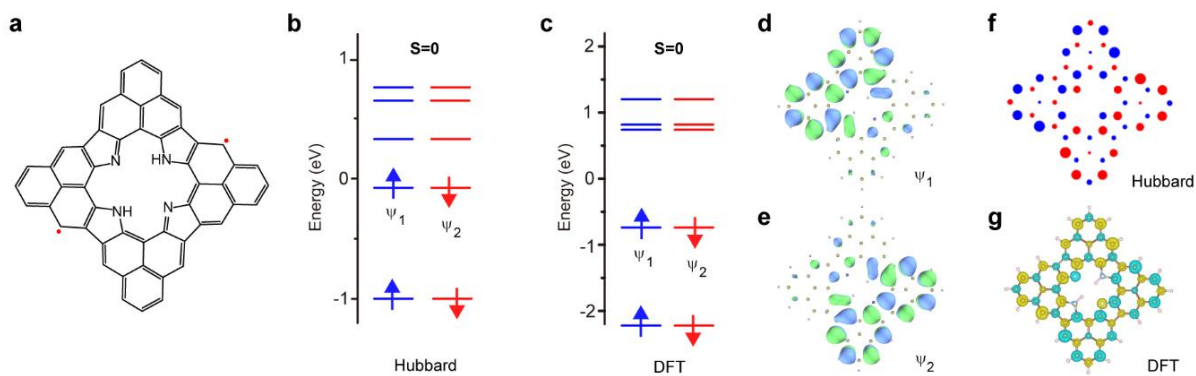


Figure 3. Spin polarized singlet ground state in the Product #2. (a), Chemical structure of the Product #2 with two unpaired π electrons as marked by red dots. (b-c), Mean-field Hubbard model and spin-polarized density functional theory calculated energy spectrum of the porphyrin in (a). (d-e), Calculated frontier orbitals of the two singly occupied states as marked in (b) and (c). (f-g), Mean-field Hubbard model and spin-polarized density functional theory calculated spin density distributions. Red/green: spin up, blue/yellow: spin down.

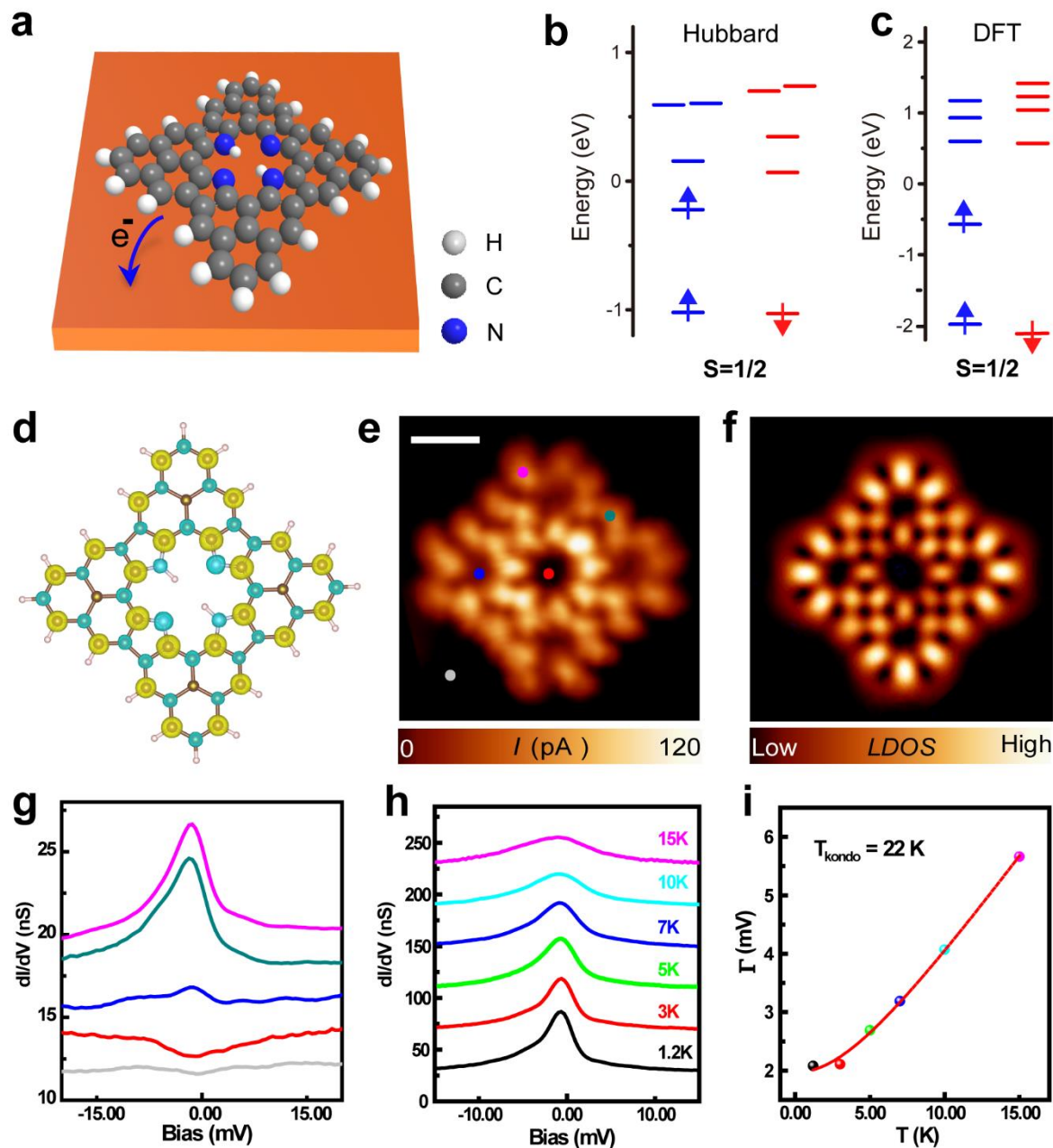


Figure 4. Electronic structure of the positively charged Product #2. (a), Illustrative showing charge transfer from the Product #2 to underneath Au(111), thus making it positively charged. (b-c), Mean-field Hubbard model and spin-polarized density functional theory calculated energy spectrum of the charged Product #2. (d), Spin-polarized density functional theory calculated spin density distributions. Green: spin up, Yellow: spin down. (e-f), Constant-height current image (Bias voltage: 1 mV, Scale bars: 0.5 nm), and simulated DOS map of the charged Product #2. (g), dI/dV spectra spectra taken on the locations marked in (e). (h), Temperature dependent dI/dV spectra showing the broadening of Kondo resonance peak as increase the temperature. (i), The full-width at half maximum of Kondo resonance peak as a function of temperature. The solid red line is fitted curve giving a Kondo temperature of 22K.

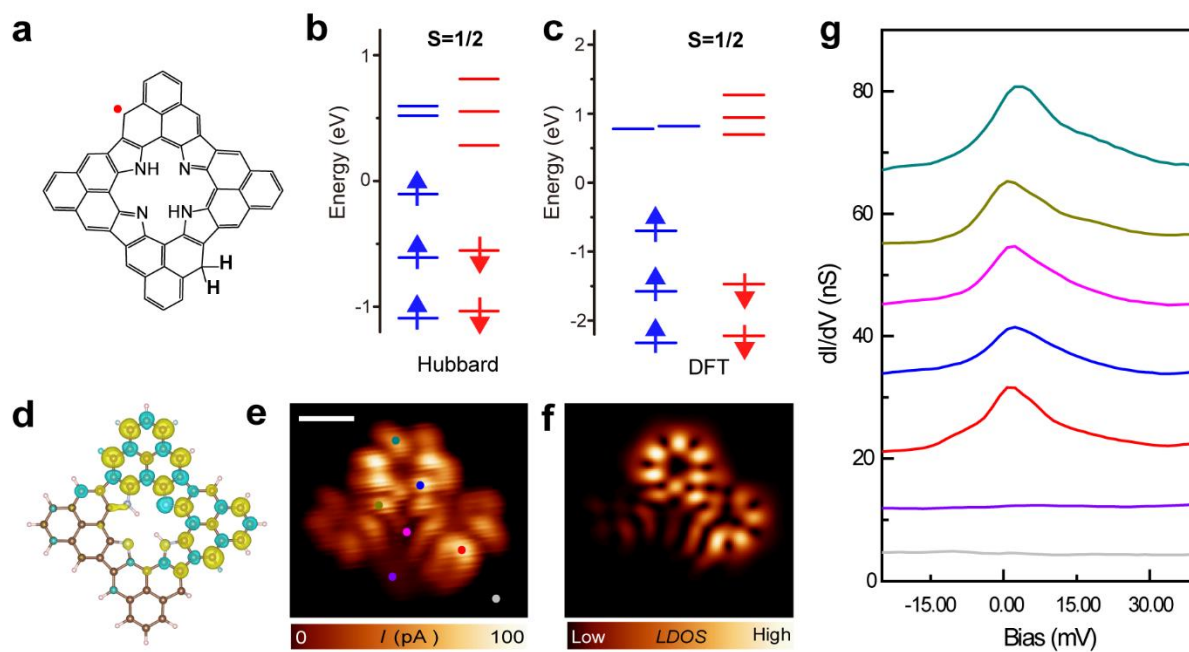


Figure 5. Electronic structure of the Product #3. (a), Chemical structure of the Product #3 with one unpaired π electron as marked by a red dot. (b-c), Mean-field Hubbard model and spin-polarized density functional theory calculated energy spectrum of the Product #3. (d), Spin-polarized density functional theory calculated spin density distributions. Green: spin up, Yellow: spin down. (e-f), Constant-height current image (Bias voltage: 1 mV, Scale bars: 0.5 nm), and simulated DOS map of the charged product #2. (g), dI/dV spectra spectra taken at the locations marked in (e).

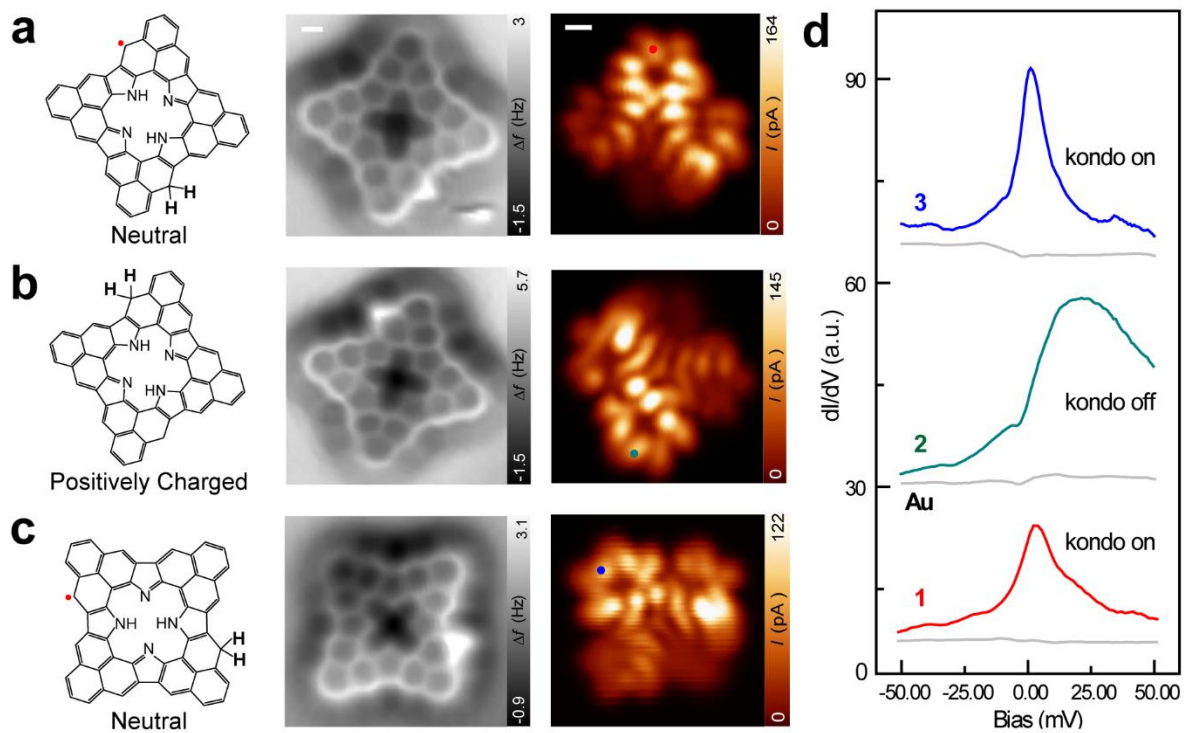
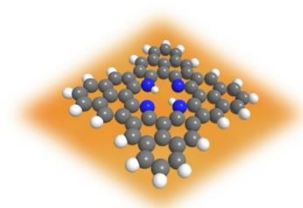
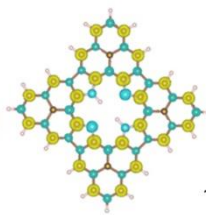


Figure 6. Switching magnetic state by STM manipulation. (a-c), From left to right: Chemical structure, nc-AFM image (Resonant frequency: 26 KHz, Oscillation amplitude: 80 pm, Scale bars: 0.2 nm), Constant-height current image (Bias voltage: 1 mV, Scale bars: 0.25 nm) of the Product #3 after two manipulation steps. (d), dI/dV spectra taken at the marked positions in (a-c).

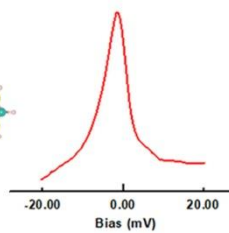
TOC Figure



Metal-free porphyrin



Delocalized π radical



Kondo resonance