Defect-Mediated Conductivity Enhancements in Na_{3-x} $Pn_{1-x}W_xS_4$ (Pn = P, Sb) using Aliovalent Substitutions

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Abstract

The sodium-ion conducting family of Na₃*Pn*S₄, with *Pn* = P, Sb, have gained interest for the use in solid-state batteries due to their high ionic conductivity. However, significant improvements to the conductivity have been hampered by the lack of aliovalent dopants that can introduce vacancies into the structure. Inspired by the need for vacancy introduction into Na₃*Pn*S₄, the solid solutions with WS₄²⁻ introduction are explored. The influence of the substitution with WS₄²⁻ for PS₄³⁻ and SbS₄³⁻, respectively, is monitored using a combination of X-ray diffraction, Raman and impedance spectroscopy. With increasing vacancy concentration improvements resulting in a very high ionic conductivity of $13 \pm 3 \text{ mS} \cdot \text{cm}^{-1}$ for Na_{2.9}P_{0.9}W_{0.1}S₄ can be observed. This work acts as a stepping-stone towards further engineering of ionic conductors using vacancy-injection via aliovalent substituents.

TOC Graphic



Lithium and sodium thiophosphates are currently of high interest for solid-state battery applications due to their high ionic conductivity and ductile nature. A multitude of structural classes such as $Li_{10}(Si/Ge/Sn)P_2S_{12}$,¹⁻⁹ $Li_6PS_5X^{10-15}$ in the case of Li^+ conductors and $Na_{11}Sn_2PS_{12}$,^{16–18} Na_3PS_4 and Na_3SbS_4 for Na^+ conductors have attracted quite some interest.^{19–24} However, while most of these compounds possess vacant sites for the mobile carrier within the structure, a prerequisite for fast ionic conduction,^{25,26} Na_3PS_4 and Na_3SbS_4 are intrinsically hindered by fully occupied Na^+ positions.

The structure of Na₃*Pn*S₄, with *Pn* = P, Sb, is displayed in Figure 1, crystallizing either in the cubic $I\bar{4}3m$ or tetragonal $P\bar{4}2_1c$ space groups. The cubic polymorph features a body centered cubic arrangement of the PnS_4^{3-} units, whereas the tetragonal polymorph exhibits a rotation of PnS_4^{3-} around the [111] axis that goes along with a minor tetragonal distortion along the *c*-axis. In the cubic structure, only one fully occupied Na⁺ position exists, while in the tetragonal structure, two Na⁺ positions can be found due to this tetragonal distortion resulting from the PnS_4^{3-} rotation.^{22,27} The resulting, fully occupied Na⁺ diffusion pathways are illustrated in Figure 1.



Figure 1. Unit cells of the (a) cubic and (b) tetragonal Na_3PnS_4 with Pn = P, Sb. A full Na^+ occupancy within the direct diffusion pathway can be found in both the (c) cubic and (d) tetragonal polymorphs.

Different isoelectronic substitutions, such as Se for S and Sb for P, have already been carried out toward improving our understanding of this material system.^{27–29} Substituting Se into the anion sublattice is possible over the whole stoichiometric range, leading to a stabilization of the

cubic polymorph and general expansion of the lattice.^{27,30} This Se substitution also leads to a lattice softening, which in turn decreases both the activation energy for ion migration and the Arrhenius prefactor.^{27,31} Interestingly, further improvement of the conductivity up to 0.74 mS·cm⁻¹ can be achieved by mixing Na₃PS₄ with small amounts of Na₄SiS₄, highlighting the potential of this material class.³² Furthermore, depending on the synthetic conditions, it is also possible to individually synthesize the cubic and tetragonal polymorphs. For example, while mechanical alloying leads to the cubic polymorph, the tetragonal structure can be obtained by classic high temperature synthesis.^{19,24,33} Compared to Na₃PS₄, Na₃SbS₄ exhibits an even higher ionic conductivity due to the overall lower activation barriers found in both polymorphs.³³ In line with the full occupancy of the Na⁺ positions, the transport properties seem to depend on the defect concentration and thus, the synthesis conditions, and less on the crystal structure itself.^{22,34} This assumption is further affirmed upon considering the exceptionally high conductivity of 3 mS·cm⁻¹ of the ball-milled t-Na₃SbS₄ (with 2.5 mol % Na vacancies).²⁹

The existence of the fully occupied positions of Na⁺ are the reason that has made it difficult to optimize the ionic conductivity. The typical aliovalent dopants within the PS4³⁻ and SbS4³⁻ units are often tetravalent such as Si⁴⁺, Ge⁴⁺, Sn⁴⁺.³⁵ However, while these dopants work well in systems with a high number of vacant sites, they introduce more Na⁺ as mobile carries and hence cannot be incorporated in the structures of Na₃PS₄ and Na₃SbS₄. While vacancy-injection via substitution of S²⁻ with Cl⁻ and Na⁺ with Ca²⁺ has been proposed,^{36,37} no substitution on the MS4³⁻ tetrahedral center was shown until recently.³⁸ Thus, inspired by the need to create more vacancies in Na₃PS₄ and Na₃SbS₄, this work explores aliovalent substitutions within Na_{3-x}*Pn*₁₋ *xWx*S₄. Incorporation of W⁶⁺ on the tetrahedral site, forming WS4²⁻ tetrahedra, leads to the introduction of vacancies via charge compensation. It should be noted that, in addition to the substitution with W⁶⁺, the substitution of Na₃*Pn*S₄ with Mo⁶⁺ was attempted but synthetically unsuccessful.



Figure 2. Obtained X-ray diffraction patterns of (a) $Na_{3-x}P_{1-x}W_xS_4$ and (b) $Na_{3-x}Sb_{1-x}W_xS_4$ with changing W content. The grey bars highlight the reflections of WS₂. Refined lattice parameters of (c) $Na_{3-x}P_{1-x}W_xSb_4$ and (d) $Na_{3-x}Sb_{1-x}W_xS_4$, showing the structural changes during substitution. In both solid solutions, changes in the lattice parameters can be found during substitution, leading to a decrease in the tetragonal distortion (c/a ratio) and ultimately a transition from the tetragonal into the cubic crystal structure.

All experimental details on the solid-state syntheses and characterization can be found in the Supporting Information. Figures 2 a,b show the X-ray diffraction patterns of Na_{3-x}*Pn*_{1-x}W_xS₄ from $0 \le x \le 0.1$. Rietveld refinements of the laboratory X-ray diffraction data were performed, but deemed unreliable due to unphysically large thermal displacement parameters of Na⁺ of > 10 Å². While these large displacement parameters serve as an indicator for fast ionic conduction, an accurate refinement of atomic positions and occupancies becomes impossible within the whole solid solution, underlining the need for high-resolution scattering experiments. Nevertheless, the lattice parameters obtained with Pawley refinements of the solid solutions Na_{3-x}P_{1-x}W_xS₄ and Na_{3-x}Sb_{1-x}W_xS₄ can be found in Figure 2c and 2d, respectively. All Pawley fits to the diffraction data, as well as a table containing the refined lattice parameters, can be

found in the Supporting Information. With increasing W content, interesting changes in the patterns and the lattice of Na_{3-x}P_{1-x}W_xS₄ and Na_{3-x}Sb_{1-x}W_xS₄ can be observed:

(1) A certain phase fraction of the precursor WS₂ can be found in all samples of W substituted Na_{3-x}P_{1-x}W_xS₄ and above x > 0.075 minor reflections of WS₂ can also be seen in Na_{3-x}Sb_{1-x}W_xS₄. As WS₂ is isotypic with the MoS₂ layered structure, all of these reflections are likely influenced by texturing of a minor impurity precursor phase fraction or suggest a limited solubility in the structure of Na_{3-x}P_{1-x}W_xS₄. It is therefore difficult to assess the real solubility limit in the structure, based on these phase fractions.

(2) However, a systematic shift of the lattice parameters can be found during the WS4²⁻ incorporation. With increasing W-content in Na_{3-x}*Pn*_{1-x}W_xS₄ the lattice parameter *a* increases while the lattice parameter *c* decreases and an overall decrease of the tetragonal distortion can be found. The changes in the lattice parameters and lattice volume (see Supporting Information Figure S3) are small for Na_{3-x}P_{1-x}W_xS₄, indicating a rather limited solubility when compared to Na_{3-x}Sb_{1-x}W_xS₄, as a much stronger trend can be found for Na_{3-x}Sb_{1-x}W_xS₄. The solubility limit seems to have been reached above $x \ge 0.1$ within Na_{3-x}Sb_{1-x}W_xS₄. This is further supported by the changing unit cell volume of Na_{3-x}Sb_{1-x}W_xS₄, which remains constant for x = 0.1 and 0.125, as shown in the Supporting Information Figure S5.

At this point, the differences in the changing lattice between the solid solutions $Na_{3-x}P_{1-x}W_xS_4$ and $Na_{3-x}Sb_{1-x}W_xS_4$ remain unclear as the tetrahedrally-coordinated ionic radius difference between Sb^{5+} vs. W^{6+} is low, whereas P^{5+} exhibits a much smaller ionic radius.³⁹ The differences are likely due to different solubility limits of WS_2 in the phases. In addition, competing effects during the introduction of Na^+ vacancies may lead to an overall lattice volume contraction (see Supporting Information Figure S5) and changing tetrahedral distortion. (3) In both solid solutions, the decreasing tetragonal distortion leads to a structural change from the tetragonal to the cubic polymorph around x = 0.1 and x = 0.075 in $Na_{3-x}P_{1-x}W_xS_4$ and $Na_{3-x}Sb_{1-x}W_xS_4$, respectively.

(4) Ultimately, while low-temperature X-ray diffraction will be needed to assess an accurate vacancy concentration and the possibility of altered diffusion pathways and behavior, a refinement of the W-occupancy in Na_{2.9}Sb_{0.9}W_{0.1}S₄ was possible using synchrotron X-ray diffraction. A Rietveld refinement of the sample with the highest W-content against the higher resolution data can be found in the Supporting Information, Figure S6. The refinement provides a W-occupancy of 11.0(4) % and with it a high Na⁺ vacancy concentration.

These structural data suggest that the introduction of W^{6+} in Na_{3-x}*Pn*_{1-x}W_xS₄ is possible. In addition, Raman spectroscopy is quite reliable in detecting the local polyhedral species in thiophosphates,^{31,40} and can be used to monitor the generation of WS₄²⁻ tetrahedra within the structure. The Raman spectra of Na_{3-x}P_{1-x}W_xS₄ and Na_{3-x}Sb_{1-x}W_xS₄, are shown in Figure 3a and 3b, respectively. For comparison, representative Raman spectra of solid WS₂ and the WS₄²⁻ ion in solution are shown.^{41–43}

In the case of Na_{3-x}P_{1-x}W_xS₄, all Raman modes of the PS₄³⁻ can be found, along with the main symmetric A_1 stretching mode at 413 cm^{-1,31} In the Raman spectra of Na_{3-x}Sb_{1-x}W_xS₄, three Raman signals can be noticed in the range of 340 cm⁻¹ to around 410 cm⁻¹ corresponding to the different modes of the SbS₄³⁻ tetrahedra in Na₃SbS₄.^{44,45} With increasing W-content, the intensity of the main modes decrease and additional Raman active modes can be found. Additionally, a different signal around 470 cm⁻¹ begins to evolve, corresponding to the WS₄²⁻ tetrahedra (shown in grey). A minor intensity corresponding to the WS₂ precursor can be found in the solid solution Na_{3-x}P_{1-x}W_xS₄, indicating a different and lower solubility limit of WS₄²⁻ in Na_{3-x}P_{1-x}W_xS₄ compared to Na_{3-x}Sb_{1-x}W_xS₄. The assignment of all Raman modes for PS₄³⁻, SbS₄³⁻, WS₂ and WS₄²⁻ can be found in the Supporting Information, Table S3 as well as all Raman spectra including x = 0.025 and x = 0.075 in Figure S6. While minor shifts and differences in intensity are expected due to the reference spectrum stemming from thiometallocomplexes in solution, comparison of the reference spectra with the data corroborate the incorporation of WS₄²⁻ tetrahedra in Na_{3-x}Sb_{1-x}W_xS₄, as well as Na_{3-x}P_{1-x}W_xS₄ to a minor extent.



Figure 3. Raman spectra of (a) $Na_{3-x}P_{1-x}W_xSb_4$ and (b) $Na_{3-x}Sb_{1-x}W_xS_4$ for representative compositions, along with digitized Raman spectra of WS_4^{2-} polyhedra in solution,^{41,42} and single

crystalline WS_2 .⁴³ In both cases, the Raman modes of the WS_4^{2-} polyhedra are visible and increase with increasing degree of substitution. In addition, WS_2 can be found in the compositional series of $Na_{3-x}P_{1-x}W_xSb_4$.

In order to gain insight into how the increasing defect concentration by Na⁺ vacancies affect the ionic transport properties, AC impedance spectroscopy was performed on Na_{3-x}*Pn*_{1-x}W_xS₄. Representative Nyquist-plots of the impedance response at -20 °C are shown for all compositions in the Supporting Information, Figure S7. In the more resistive samples of Na₃PS₄ and Na₃SbS₄, two semicircles can be found and the bulk and grain boundary contributions can be fully distinguished. However, in all other samples, only one process is visible followed by the blocking electrode. While the capacitances of the resolvable spectra are in the range of 38– 64 pF·cm⁻² for the high-frequency process, corroborating that predominantly bulk transport is observed,⁴⁶ the bulk and grain processes of the highly conducting compositions cannot be distinguished. Therefore, all ionic conductivities and activation barriers presented here need to be considered as total conductivities.

The Arrhenius plots of all compositions can be found in Figure S7. It should be mentioned that for Na_{3-x}Sb_{1-x}W_xS₄ with x = 0.05 and x = 0.075, a change in the Arrhenius behavior can be observed at temperatures slightly above room temperature (see Supporting Information, Figure S8). Typically, the cubic structure is the high-temperature polymorph. In Na₃SbS₄, one would expect the phase transition from the tetragonal to the cubic structure at 170 °C.³³ Considering the Arrhenius behavior during heating, as well as the structural changes in Figure 2, it seems that the W substitution not only lowers the tetragonal distortion in the structure, but also the phase transition temperature, similar to Na₃PS_{4-x}Se_x.²⁷ Therefore, with increasing W content, the phase transition temperature seems to decrease. For x < 0.05, the phase transition is likely located above room temperature, whereas at $x \ge 0.075$, the cubic phase seems to be the stable polymorph at room temperature.

Figure 4 shows the ionic conductivity and activation barriers for $Na_{3-x}Pn_{1-x}W_xS_4$. With increasing W content, and with-it increasing vacancy concentration, the activation barriers for ionic transport decrease in both series of compositions. While the structural data suggests a low solubility of W⁶⁺ in $Na_{3-x}P_{1-x}W_xS_4$, the Raman data show the incorporation of WS_4^{2-} and the transport properties are affected. The incorporation of W^{4+} can be ruled out as a d² electronic system would not exist in a tetrahedral symmetry as seen in the Raman and it would further increase the electronic conductivity. In order to assess if the incorporation of WS_4^{2-} increases the electronic conductivity, DC polarization was exemplarily performed on $Na_{2.9}Sb_{0.9}W_{0.1}S_4$ and the results are shown in Figure S9. An electronic conductivity of $3.4 \pm 0.4 \cdot 10^{-6} \text{ mS} \cdot \text{cm}^{-1}$ shows a negligible electronic conductivity that is seven orders of magnitude lower than the ionic conductivity.

As the overall lattice volume within the solid solution is decreasing, the decreasing activation barriers are likely related to changing defect formation enthalpies due to the incorporation of vacancies, rather than changing widths of the diffusion pathways. Theoretical work also suggests that incorporation of vacancies should lead to lower activation barriers.³⁴ With these changes in the activation barriers and the changing carrier densities from the vacancy-injection, the ionic conductivity of both series increases significantly. While the W-substituted samples immediately exhibit a sharp increase of the conductivity and then a steadily rising conductivity with increasing W-content, the solid solution Na_{3-x}P_{1-x}W_xS₄ shows a much steeper jump at the tetragonal to cubic transition. As previously mentioned, the values presented here represent total conductivities and, especially in the pristine compositions, a large grain boundary contribution can be observed. The larger grain resistance is likely due to the differences in the synthetic procedure when compared to literature. Nevertheless, in the case of $Na_{3-x}P_{1-x}W_xS_4$, the unsubstituted compound exhibits a total conductivity of $0.04 \pm 0.01 \text{ mS} \cdot \text{cm}^{-1}$ and achieves $13 \pm 3 \text{ mS} \cdot \text{cm}^{-1}$ for x = 0.1, representing an increase of the conductivity exceeding three orders of magnitude. In the case of $Na_{3-x}Sb_{1-x}W_xS_4$, the pristine compound exhibits a conductivity of $0.01 \pm 0.002 \text{ mS} \cdot \text{cm}^{-1}$ and achieves $41 \pm 8 \text{ mS} \cdot \text{cm}^{-1}$ for x = 0.1. To the best of our knowledge, this is the highest ionic conductivity measured in the class of thiophosphate ionic conductors to date, showing the drastic influence of vacancy injection in $Na_{3-x}Pn_{1-x}W_xS_4$.



Figure 4. (a) Activation barriers of $Na_{3-x}Pn_{1-x}W_xS_4$ with Pn = P, Sb. and (b) room-temperature conductivities With increasing vacancy content, the conductivity changes drastically alongside with a continuous decrease of the activation barrier. Dashed lines represent guides-to-the-eye.

In conclusion, this work reports the aliovalent substitution in the solid solution series $Na_{3-x}Pn_{1-x}W_xS_4$ with Pn = P and Sb. X-ray diffraction and Raman spectroscopy corroborate that solid solutions can be formed and that WS_4^{2-} polyhedra are incorporated into the structure. With increasing W content, the structures become less distorted leading to a structural transition from the tetragonal to the cubic phase. The increasing vacancy concentration leads to orders of magnitude improvements in the ionic conductivity, exhibiting a very high ionic conductivity of $13 \pm 3 \text{ mS} \cdot \text{cm}^{-1}$ for $Na_{2.9}P_{0.9}W_{0.1}S_4$ and $41 \pm 8 \text{ mS} \cdot \text{cm}^{-1}$ for $Na_{2.9}Sb_{0.9}W_{0.1}S_4$.

Future work needs to further explore the structural changes and occupancies during the substitution using low-temperature, high-resolution diffraction methods, as well as elucidating the real defect concentration and changing phase transition temperatures. Further, an assessment of the influence of the W substitution on the electrochemical stability window and the applicability in a solid-state battery is needed. Overall, this work shows that aliovalent substitutions toward increasing vacancy content is possible by introducing WS_4^{2-} in the thiophosphates and thioantimonates, thereby opening up a new and very promising route for optimization of these classes of superionic conductors, beyond Na₃PS₄ and Na₃SbS₄.

Supporting Information.

All experimental details on synthesis, X-ray diffraction and impedance spectroscopy can be found here. All structural data as obtained from Pawley refinements can be found here. In addition, assignments of the Raman modes and all Raman spectra are given. Further, the impedance data and Arrhenius plot, as well as information on the Arrhenius behavior and phase transitions. Further, a Rietveld refinement against synchrotron data and DC polarization measurements of Na_{2.9}Sb_{0.9}W_{0.1}S₄ can be found.

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Notes

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