

Two are better than one: a design principle for ultralong persistent luminescence of pure organics

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Abstract

Because of their innate ability to store and then release energy, long persistent luminescence (LPL) materials have garnered strong research interest in a wide range of multidisciplinary fields, such as biomedical sciences¹, theranostics², and photonic devices³. Although many inorganic LPL systems with afterglow durations of up to hours and days have been reported⁴, organic LPL (OLPL) systems have had difficulties reaching similar timescales⁵. We propose in this work a design principle based on the successes of inorganic systems to produce an OLPL system through the use of a strong organic electron trap. The resulting system generated detectable afterglow for up to 7 hours, significantly longer than any other reported OLPL system. The design strategy demonstrates an easy methodology to develop organic long persistent phosphors, opening the door to new OLPL materials.

Main

Long persistent luminescent⁶ (LPL) materials are highly sought after in areas such as life sciences⁷, the biomedical field⁸, optoelectronics⁹, photovoltaics¹⁰, and energy storage¹¹ as they offer fascinating possibilities for their ability to store and slowly release excited state energy. The most successful LPL materials make use of transition metals and rare-earth metal ions^{12,13}. Although the metals grant exceptionally long afterglows, they are not without their inherent drawbacks. In addition to the high costs of these rare-earth metals, many inorganic LPLs require harsh synthetic procedures⁴, further increasing processing costs. LPL materials are also desirable for biomedical

applications¹⁴ for their ability to overcome autofluorescence but metal-containing long-lived phosphors cloud their safety for such uses. Organic LPL (OLPL) materials offer the promise of a multitude of benefits: easy synthesis and modification for targeted functionality, easy processing, and can be biologically inert. However, the development of OLPL materials have encountered many obstacles. To access long lived states in organic compounds, there have been many designs to exploit the excited triplet state. Though access to and from the triplet state is a forbidden process and once thought to be too inefficient for effective use at room temperature¹⁵, recent advances have vastly increased intersystem crossing efficiency by enhancing spin-orbit coupling (SOC) with the use of heteroatoms^{16,17}, the carbonyl functional group^{18,19}, heavy atom effects²⁰, and multimer-enhance ISC²¹. Equally important is protecting the long-lived triplet due to the fact that they are particularly sensitive to molecular vibrational quenching and atmospheric oxygen. In this regard, recent works have accomplished this through the use of crystals^{22,23}, MOFs²⁴, H-aggregation²⁵, and others^{26,27}. Although there have been many achievements in generating organic room-temperature phosphorescent systems²⁸⁻³¹, their lifetimes cannot match the durations of inorganic LPL materials.

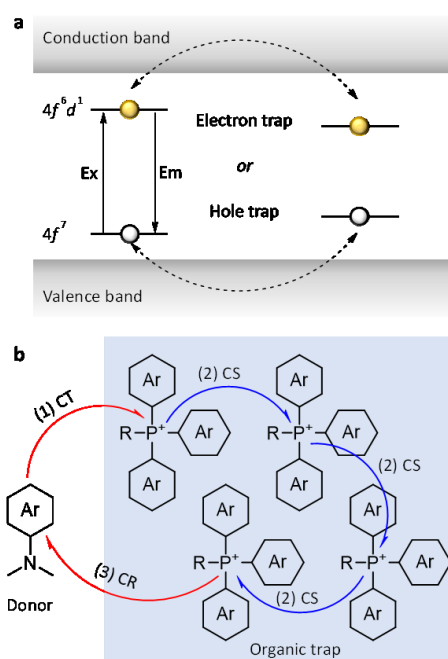


Figure 1 | Schematic representations of mechanisms behind inorganic and organic persistent luminescence.

a, Inorganic long persistent luminescence can be achieved through electron or hole traps mechanisms. In the electron trap, it is proposed after excitation (Ex), the excited electron travels through the conduction band to an electron accepting trap. In the hole trap, an electron fills the hole by traveling through the valence band. In both cases, relaxation becomes blocked, either because the excited electron has migrated away, or the hole has become filled. Thermal disturbances restore the electron or hole, producing afterglow emission (Em). **b**, A proposed mechanism for organic long persistent luminescence where a cationic quaternary phosphonium core acts as an organic trap and an aromatic amine as an electron donor molecule. First (1) photoinduced charge transfer (CT) occurs between the donor and acceptor molecules followed by (2) charge separation (CS). Multiple CS can occur before the final step (3), charge recombination (CR), resulting in OLPL.

By applying concepts from the mechanistic understand of inorganic LPL systems to OLPLs, it can be possible to merge the diverse benefits of organic systems with the performance of inorganics. One of the most successful methods to achieve persistent luminescence is through the use of traps (Fig. 1a), which allow for energy storage and its subsequent slow release via thermal disturbances. Inspired by the role of inorganic trap species, we hypothesized a purely organic system made of two parts, a biontentic system of an emitter and trap (Fig. 1b). Just as the heavy metals store charge carriers and slowly release them back into the system for eventual emission, organic molecules could theoretically do the same thing.

To test our mechanistic principle, we first set out to find an ionic core system that could work as a strong electron acceptor. In our search, we discovered, synthesized, and characterized (Fig. S1-5) a series of pure organic quaternary phosphonium bromide salts, TPP-3CB, TPP-4CB, TPP-3C2B, and TPP-4C2B, which exhibit observable afterglow (Fig. 2 and Supplementary Video 1-4). The phosphonium core serves as a good acceptor and the bromide, as a donor. These four compounds were purified and recrystallized several times before characterization, with the ionic nature of the compounds aiding the formation of the crystal growth. Their crystal structures displayed a distorted tetrahedral geometry (108.5° to 111.2°) and the distance between the phosphonium and bromide ions were measured to be in the range of 4.52 Å to 4.90 Å (Fig. S6a). The electrostatic interactions as well as $\text{CH}\cdots\pi$ interactions (Fig. S6b), work together to suppress molecular motion in the crystal enhancing phosphorescence efficiency. The absorption spectra of TPP-3CB, TPP-4CB, TPP-3C2B, and TPP-4C2B were recorded in dilute acetonitrile solutions (1×10^{-5} M), and they showed similar absorption bands at 265 nm, 268 nm, and 275 nm which were attributed to π - π^* transitions from the phenyl moieties of the phosphonium core (Fig. S7). The crystals of all four phosphonium bromide salts showed an emission maximum of 480 nm (Fig. 2b), and surprisingly, their excitation spectra (Fig. S8) at 480 nm emission revealed a maximum excitation band at 310 nm which was attributed to a photoinduced charge transfer between the phosphonium core and its bromide counterion. Thus, subsequent PL and lifetime measurements of TPP-3CB, TPP-4CB, TPP-3C2B, and TPP-4C2B crystals were recorded by 310 nm excitation. Exceptionally, the phosphorescence quantum yield (Φ_p) of TPP-4C2B was found to be 17.2% which was ~ 4 times higher than the others. From the time-resolved PL decay curves (Fig. 2c), the average lifetimes of TPP-3CB, TPP-4CB, TPP-3C2B, and TPP-4C2B crystals at 480 nm measured at room temperature (298 K) were 157 ms, 200 ms, 159 ms, and 164 ms, respectively, indicating the ultralong after-glow nature of these materials and the average lifetimes of these phosphonium salts measured at 77 K in dilute solution were found to be 632 ms, 410 ms, 252 ms, and 717 ms, respectively (Fig. S9-12). Under 254 nm excitation, intense greenish blue emission could be observed from the crystals (Fig. 2d).

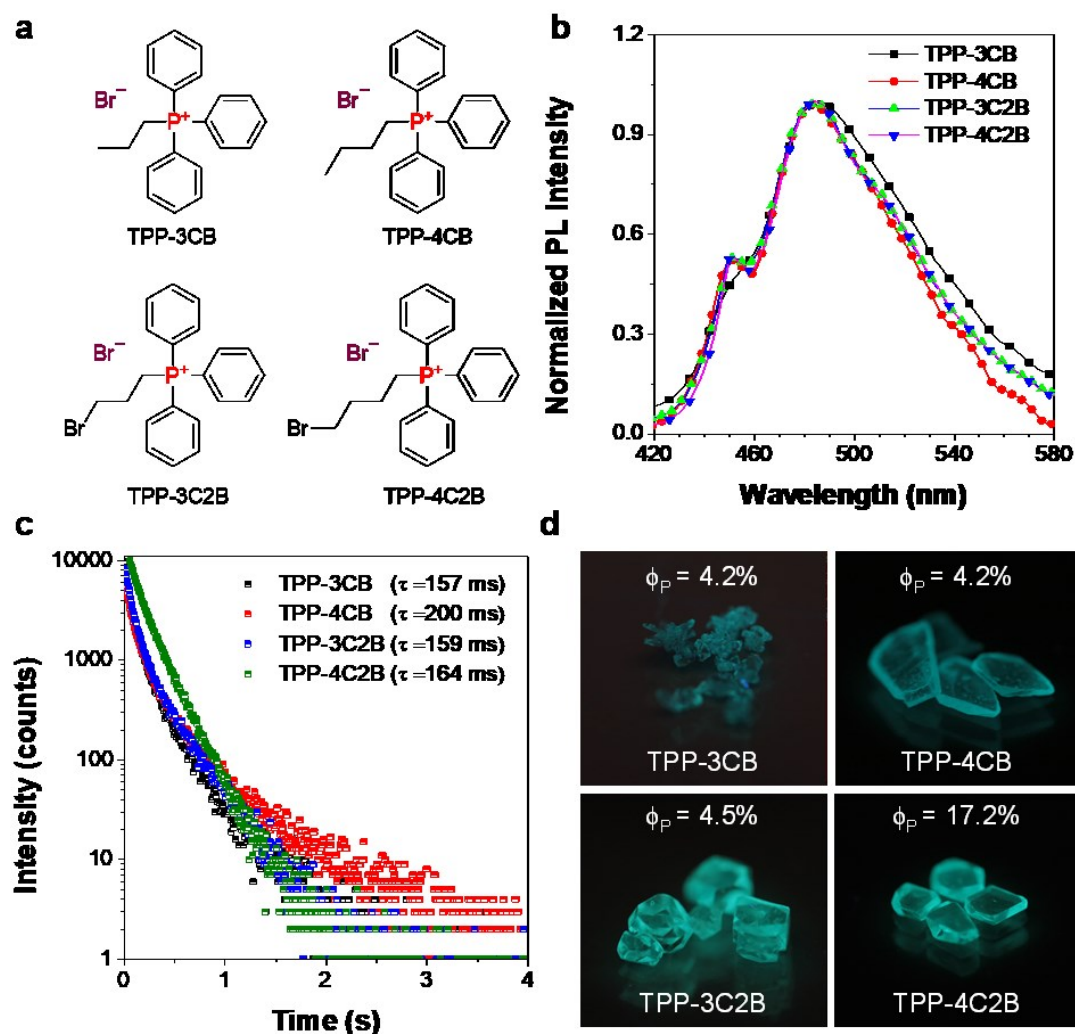


Figure 2 | Chemical structure and photophysical properties of phosphonium halides. a, Chemical structures, b, PL spectra and c time-resolved PL decay (at 480 nm) of crystalline powders of TPP-3CB, TPP-4CB, TPP-3C2B, and TPP-4C2B at 298 K. d, Photographs of TPP-3CB, TPP-4CB, TPP-3C2B, and TPP-4C2B under UV excitation at 254 nm as well as their phosphorescence quantum yields (Φ_P).

We further took the crystal of TPP-3CB as a target model and employed combined quantum mechanics and molecular mechanics method (QM/MM) to simulate the photophysical process in the crystal, optimizing the excited singlet (S_1) and triplet (T_1) states and the molecular orbitals furnished by the ONIOM method (Fig. 3a-c and S13). At the S_1 geometry (Fig. 3a), the molecular orbitals of the S_1 and T_1 of TPP-3CB (Fig. S14) show complete charge transfer from its bromide anion to the phosphonium cation, furnishing excited radical pairs. Initially, a large separation of the orbitals involved in the transition from S_0 to S_1 resulted in a small exchange energy from S_1 to T_1 state, producing an energy gap between S_1 and T_1 that is smaller than 0.1 eV facilitating intersystem crossing. Furthermore, the bromide anion endows the transition process with a spin-orbit coupling constant as large as 327.39 cm^{-1} , which further enhances the transition between the S_1 and T_1 state. As such, both the complete charge-transfer characteristic and the strong SOC provides a firm foundation for the spin-orbit charge transfer (SOCT) intersystem crossing process from S_1 to T_1 for the phosphonium bromide complex. Furthermore, the excited state radical pair is also susceptible to hyperfine coupling (HFC) effects which provide another channel

for intersystem crossing³²; as the radicals are located on two separate locations, they experience different local magnetic fields leading to spin-mixing where rephasing (Fig. 3e) and spin-flipping transitions facilitate intersystem crossing. In this way, upon excitation the molecule initially becomes an excited singlet radical pair (¹RP), but due to SOCT and HFC effects, not only does the ¹RP easily intersystem cross to become an excited triplet radical pair (³RP), the ³RP can also easily intersystem cross back into the ¹RP (Fig. 3d-e). According to the calculations, the RP state changes when the molecule relaxes to its T₁ geometry, that is, at the optimized T₁ geometry, charge-transfer/radical pair nature is lost. The molecular orbitals indicate at the T₁ minimum, the exciton is localized on the phosphonium core (Fig. 3c) indicative of a second charge transfer. The T₁ becomes energetically separated from the S₁ with an energy difference of 0.957 eV thus trapping the excited state for phosphorescence.

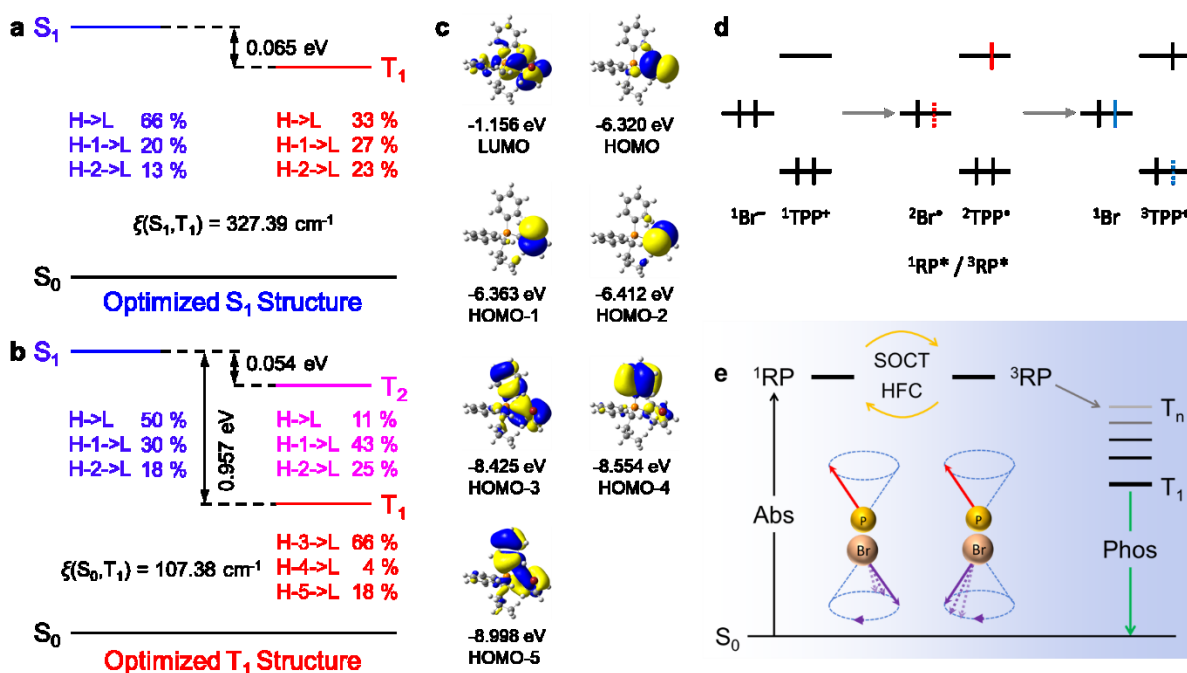


Figure 3 | QM/MM calculations at TD-DFT level of the singlet and triplet excited states. The TPP-3CB crystal structure was used to obtain an optimized geometry for the **a**, S_1 and **b**, T_1 excited state surrounded by 43 molecules. **c**, Molecular orbitals calculated from the optimized T_1 structure. **d**, Schematic representation of electrons in TPP-3CB starting from the ground state, after the first charge transfer excitation (colored in red) into the radical pair (RP), and the second charge transfer (colored in blue) into the T_1 state. **e**, Modified Jablonski diagram showing the formation of the singlet radical pair (¹RP) after absorption (Abs). Spin-orbit charge transfer (SOCT) and hyperfine coupling (HFC) can aid in the interconversion between the ¹RP and the triplet radical pair (³RP) as illustrated by the rephasing of the radical pair electrons in separate locations. The RP then relaxes into the T_1 state before phosphorescence is observed.

Considering that it has been reported that excited-state charge separation can be beneficial for OLPLs³³ and, taking cues from inorganic LPL systems, if we introduce moieties to “trap” charge carriers, it may be possible to further extend the afterglow duration. To verify our hypothesis, we attempted to develop a biontentic OLPL system using a phosphonium bromide salt as a strong electron acceptor and N,N-dimethylaniline (DMA) as an electron donor.

We were able to dope DMA into TPP-3C2B by first mixing DMA with a dichloromethane (DCM) solution of TPP-3C2B and then allowing crystals to grow. These crystals exhibited a strong emission at 500 nm upon 365 nm excitation (Fig. 4a). For reference, the absorption spectrum of DMA in acetonitrile has two bands at 253 nm and 300 nm (Fig. S15) and has an emission band at 350 nm. Since neither DMA nor TPP-3C2B have absorption bands at 365 nm (Fig. S16), this suggests the excitation of the TPP-3C2B:DMA crystals has photoinduced charge-transfer characteristics. Furthermore, the emission band of TPP-3C2B:DMA at 500 nm is ~20 nm red-shifted from TPP-3C2B, suggesting an exciplex emission (Fig. S17). In addition, both HPLC (Fig. S18a) and HRMS (Fig. S19) characterizations confirm the presence of DMA in the TPP-3C2B:DMA samples. After 60 seconds of 365 nm UV excitation, long-lived charge-separated states could be generated, and emission could be detected for up to 7 h (Fig. 4b and Supplementary Video 5) by a CMOS camera at room temperature and ambient conditions. It should be noted that near the 7 h mark, though weak, faint emission from the crystal could still be captured by the camera (Fig. S20). Due to the persistent nature of the emission, we were unable to measure the lifetime of its decay curve. Instead of using a spectrophotometer to measure the lifetime, the average greyscale values of the photographs were calculated using a fluorescence microscope to represent the overall brightness of the fluorescent image. The corresponding data points give an exponential decay with calculated average lifetime of 22 minutes from the luminescence of TPP-3C2B:DMA (Fig. 4c). The TPP-3C2B:DMA crystals have an emission quantum efficiency (ϕ) of 14.9%.

The crystalline nature of TPP-3C2B:DMA protects the photo-generated radicals from atmospheric oxygen and suppresses nonradiative deactivation pathways allowing for such a long afterglow. These crystals exhibit good stability showing no visible changes to its afterglow duration even after being kept for more than 45 days under dark conditions (Fig. S21 and Supplementary Video 6).

The reproducibility of this OLPL system was repeated several times by simply mixing of DMA/TPP-3C2B (10:1 mol ratio) in DCM/EA (1:1 v/v) with varying maximum LPL durations ranging from 1 to 5 h. We believe a key factor of the OLPL duration is the amount of DMA trapped inside the crystals. At this point in time, we are investigating how to control this parameter.

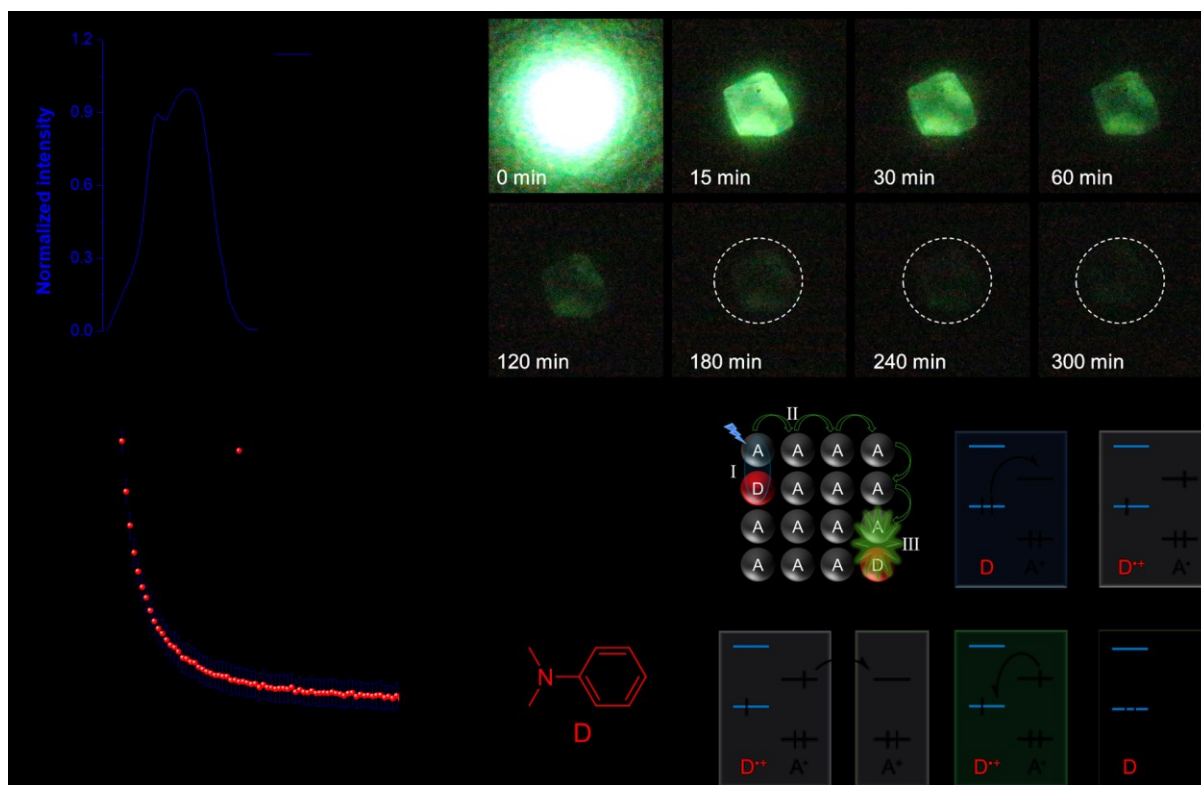


Figure 4 | Chemical structure and photophysical properties of TPP-3C2B:DMA. **a**, Excitation and emission spectra of TPP-3C2B:DMA crystals. Excitation spectrum was measured for the 500 nm emission band, and the emission spectrum was measured by 365 nm excitation. **b**, Photographs of TPP-3C2B:DMA crystals taken under 365 nm UV excitation and its subsequent afterglow. **c**, Intensity values measured from the photographs of the afterglow plotted over time. The plotted intensity values in red were averaged with the error bars in blue indicating the standard deviation. **d**, Left: chemical structures of TPP-3C2B labelled (A) for acceptor, and DMA labelled (D) for donor. Right: scheme representing the formation of the radical pair through charge transfer, subsequent charge separation and finally charge recombination to achieve long persistent luminescence.

Through careful analysis of the TPP-3C2B:DMA, it becomes possible to draw mechanistic insights to help design OLPL systems with even longer durations. The success of TPP-3C2B gives insight into factors that endow the system with such long persistent luminescence. One of the most important factors to the extraordinary duration of TPP-3C2B:DMA lies with the organic trap. Protection of the excited state is of utmost importance to ensuring high efficiency. In this case, the excited radical species, which are inherently reactive, can be easily lost. The phenyl rings decorating the phosphonium core stabilize the radical through delocalization and protect it by acting as bulky barriers that hinder unwanted reactions. In addition, the ionic nature of the material also aids in the crystallization, rigidifying the system to reduce non-radiative vibrational quenching via phonon emission. Tightly packed crystals also protect the system from oxygen quenching maximizing emission efficiency. Another important factor that bestows such long durations is the trap's ability to hold onto the electron. The cationic nature of the phosphonium core is an ideal target to accept the electron of the photoinduced charge transfer process. Furthermore, being surrounded by other cationic phosphoniums, the excited radical can migrate to multiple cores (Fig. 4d), before finally recombining with a DMA radical resulting in the observed emission. Stronger, and better designed traps can likely even further increase the time the radical is stored before it jumps to another location.

The principles exemplified by the success of TPP-3C2B can be applied to future OLPL systems producing even longer-lived emission and of varied colors and wavelengths. In addition, the OLPL system can be further tuned by modifying the donor molecule. Our group is currently working to develop OLPL systems based on these new insights.

In conclusion, we demonstrate a guiding design principle using strong organic traps to produce OLPL with previously unseen durations. The organic trap stabilizes and protects the excited radical, allowing the system to slowly recombine generating persistent luminescence. In TPP-3C2B:DMA, the cationic triphenylphosphonium core serves as a perfect trap, by first acting as a strong electron acceptor of the photoinduced charge transfer event, and then serving as multiple protective traps before the radical finally migrates back to a DMA for recombination. The insights demonstrated by our investigation will allow the future generation of highly persistent organic luminescent systems that may further propel the development and applications in the fields of medical science and optoelectronic devices.

Methods

All reagents were purchased from J&K Scientific. HPLC grade solvents were purchased from Merck. All the molecules synthesized were purified by column chromatography and recrystallized using double layer solution diffusion with dichloromethane/hexane or dichloromethane/ethyl acetate for three times, and fully characterized by ^1H NMR, ^{13}C NMR, ^{31}P NMR, high-resolution mass spectroscopy (HRMS) and elemental analysis. ^1H , ^{13}C , and ^{31}P NMR spectra were recorded on a Bruker AV 400 Spectrometer at 400, 100 and 162 MHz in CDCl_3 , respectively. High-resolution mass spectra were recorded on a GCT premier CAB048 mass spectrometer operating in MALDI-TOF mode. The elemental analysis was performed on a Thermo Finnigan Flash EA1112. Gel filtration chromatography was performed using a ZORBAX SB-C18 column (Agilent) conjugated to an Agilent 1260 Infinite HPLC system. Before running, each sample was purified via a 0.22 μm filter to remove any aggregates. The flow rate was fixed at 1.0 mL/min, the injection volume was 100 μL and each sample was run for 10 min. The absorption wavelength used was set at 254 nm. Acetonitrile/water (95:5 v/v) was used as the running solvent system. The photoluminescence spectra were measured on a PerkinElmer LS 55 spectrophotometer and Horiba Fluoromax 4 spectrofluorometer. Fluorescence quantum yields of the solids were recorded on Horiba Fluoromax 4 at room temperature with a calibrated integrating sphere system. The lifetime was measured on an Edinburgh FLSP 920 fluorescence spectrophotometer equipped with a Xenon arc lamp (Xe900) and a microsecond flash-lamp (uF900). Single crystal data were collected on a Bruker Smart APEXII CCD diffractometer using graphite monochromated Cu K α radiation ($\lambda = 1.54178 \text{ \AA}$). The photos and videos were recorded by a Canon EOS 7D Mark II. The TD-DFT calculations were performed with the Gaussian 16 program.

Data availability

The X-ray crystallographic data for the structures of TPP-3CB, TPP-4CB, TPP-3C2B, and TPP-4C2B have been deposited at the Cambridge Crystallographic Data Centre (CCDC), under deposition numbers 1919500-1919503. These data can be obtained free of charge from The Cambridge Crystallographic Data Centre via www.ccdc.cam.ac.uk/data_request/cif. Other data are available from the authors upon reasonable request.

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Contributions

P.A., N.L.C.L., J.L., X.Z., Z.H., R.T.K.K., J.W.Y.L., and B.Z.T. conceived the experiments. P.A., N.L.C.L., and B.Z.T. prepared the paper. P.A. and N.L.C.L. were primarily responsible for the experiments. H.H.Y.S. and I.D.W. carried out the single crystal X-ray diffraction measurements and analyses. J.L. and Q.P. contributed to the QM/MM TD-DFT calculations. All authors contributed to the data analyses.

Competing interests

The authors declare no competing interests.

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